

Single Eddy Mixing in Cahn-Hilliard Flows

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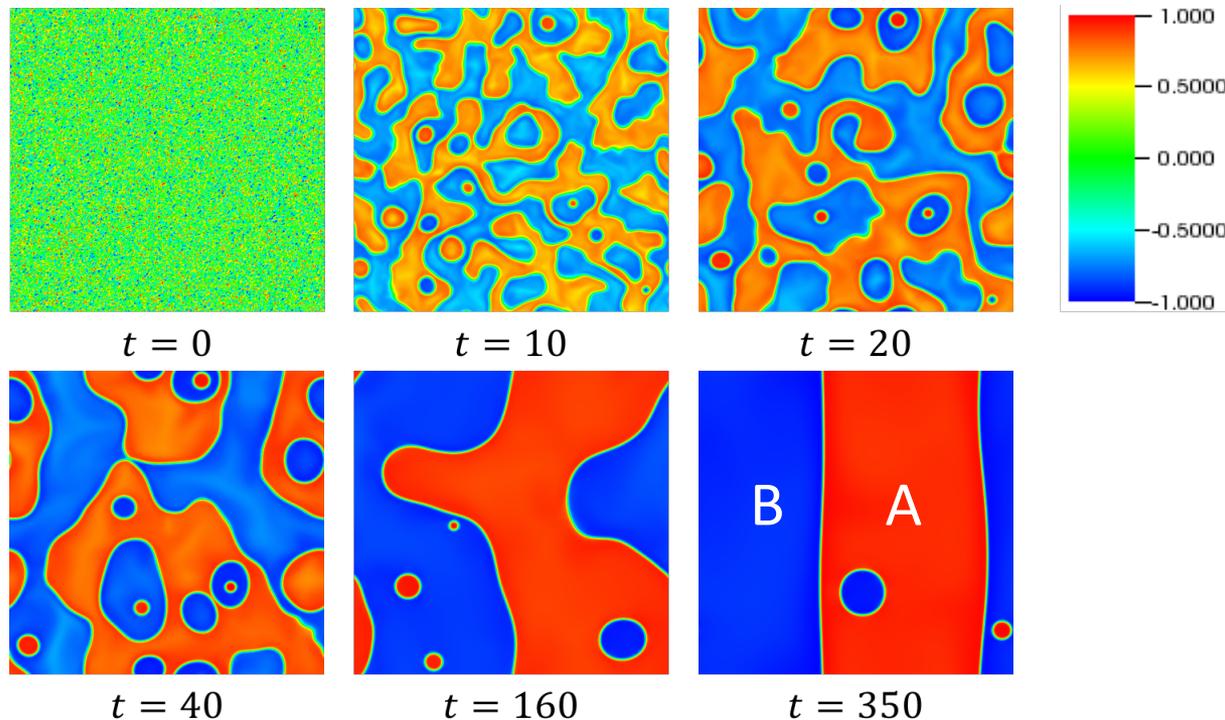
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What Is the CHNS System?

- Elastic media – Fluid with internal DoFs → “springiness”
- The Cahn-Hilliard Navier-Stokes (CHNS) system describes ***phase separation*** for binary fluid (i.e. ***Spinodal Decomposition***)
- Miscible phase -> Immiscible phase



What Is the CHNS System?

- How to describe the system: the concentration field
- $\psi(\vec{r}, t) \stackrel{\text{def}}{=} [\rho_A(\vec{r}, t) - \rho_B(\vec{r}, t)]/\rho$: scalar field
- $\psi \in [-1, 1]$
- CHNS equations:

$$\partial_t \psi + \vec{v} \cdot \nabla \psi = D \nabla^2 (-\psi + \psi^3 - \xi^2 \nabla^2 \psi)$$

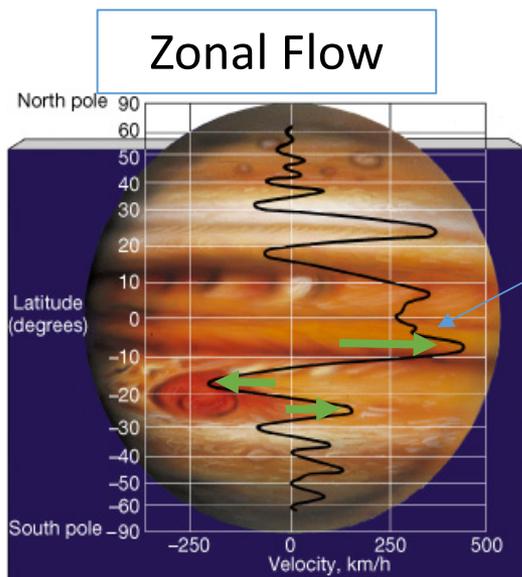
$$\partial_t \omega + \vec{v} \cdot \nabla \omega = \frac{\xi^2}{\rho} \vec{B}_\psi \cdot \nabla \nabla^2 \psi + \nu \nabla^2 \omega$$

Why Care?

- Useful to examine familiar themes in plasma turbulence from new vantage point
- Some key issues in plasma turbulence:
 1. Electromagnetics Turbulence
 - CHNS vs 2D MHD: analogous, with interesting differences.
 - Both CHNS and 2D MHD are *elastic* systems
 - Most systems = 2D/Reduced MHD + many linear effects
 - Physics of dual cascades and constrained relaxation → relative importance, selective decay?
 - Physics of wave-eddy interaction effects on nonlinear transfer (i.e. Alfvén effect \leftrightarrow Kraichnan)

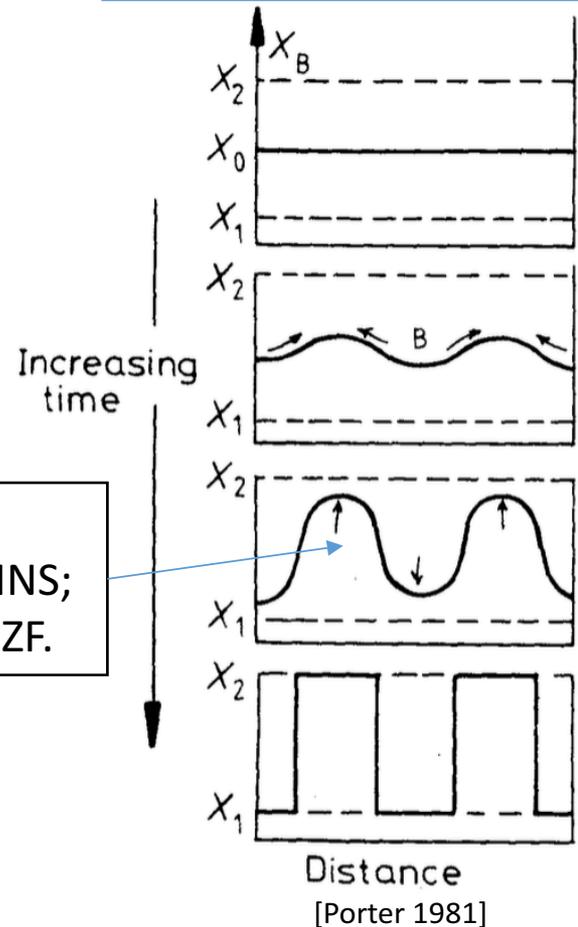
Why Care?

2. Zonal flow formation → negative viscosity phenomena → phase separation process
- ZF can be viewed as a “spinodal decomposition” of momentum.
 - What determines scale?



Arrows:
 ψ for CHNS;
flow for ZF.

Spinodal Decomposition



Why Care?

3. “Blobby Turbulence” → how to understand blob coalescence and relation to cascades?
 - CHNS is a naturally blobby system of turbulence.
 - What is the role of structure in interaction?
 - How to understand multiple cascades of blobs and energy?

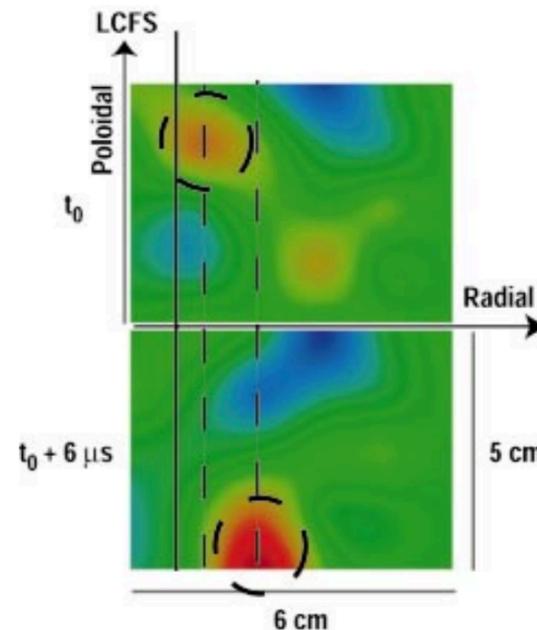


FIG. 4. (Color) Two frames from BES showing 2-D density plots. There is a time difference of $6 \mu\text{s}$ between frames. Red indicates high density and blue low density. A structure, marked with a dashed circle and shown in both frames, features poloidal and radial motion.

[J. A. Boedo et.al. 2003]

- CHNS exhibits all of the above, with many new twists

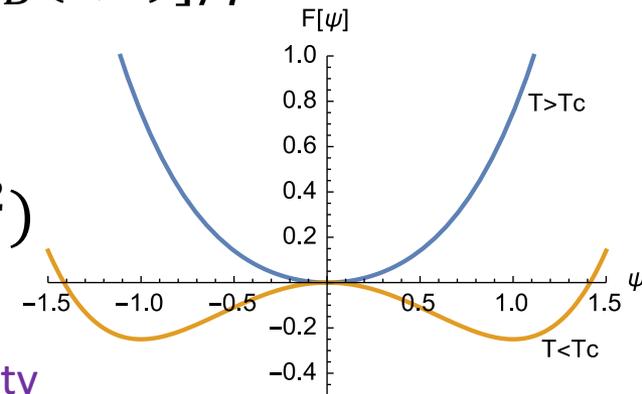
Outline

- A Brief Derivation of the CHNS Model
- 2D CHNS and 2D MHD
- Linear Wave
- Single Eddy Mixing in 2D MHD: Expulsion
- Single Eddy Mixing in 2D Cahn-Hilliard Flow
- 3 Stages
- Band Mergers
- Time Scales
- Conclusions

A Brief Derivation of the CHNS Model

- Second order phase transition \rightarrow Landau Theory.
- Order parameter: $\psi(\vec{r}, t) \stackrel{\text{def}}{=} [\rho_A(\vec{r}, t) - \rho_B(\vec{r}, t)]/\rho$
- Free energy:

$$F(\psi) = \int d\vec{r} \left(\underbrace{\frac{1}{2} C_1 \psi^2 + \frac{1}{4} C_2 \psi^4}_{\text{Phase Transition}} + \underbrace{\frac{\xi^2}{2} |\nabla \psi|^2}_{\text{Gradient Penalty}} \right)$$



- $C_1(T), C_2(T)$.
- Isothermal $T < T_C$. Set $C_2 = -C_1 = 1$:

$$F(\psi) = \int d\vec{r} \left(-\frac{1}{2} \psi^2 + \frac{1}{4} \psi^4 + \frac{\xi^2}{2} |\nabla \psi|^2 \right)$$

A Brief Derivation of the CHNS Model

- Continuity equation: $\frac{d\psi}{dt} + \nabla \cdot \vec{J} = 0$.
- Fick's Law: $\vec{J} = -D\nabla\mu$.
- Chemical potential: $\mu = \frac{\delta F(\psi)}{\delta\psi} = -\psi + \psi^3 - \xi^2\nabla^2\psi$.
- Combining \rightarrow Cahn Hilliard equation:

$$\frac{d\psi}{dt} = D\nabla^2\mu = D\nabla^2(-\psi + \psi^3 - \xi^2\nabla^2\psi)$$

- $d_t = \partial_t + \vec{v} \cdot \nabla$.
- Surface tension: force in Navier-Stokes equation:

$$\partial_t \vec{v} + \vec{v} \cdot \nabla \vec{v} = -\frac{\nabla p}{\rho} - \psi \nabla \mu + \nu \nabla^2 \vec{v}$$

- For incompressible fluid, $\nabla \cdot \vec{v} = 0$.

2D CHNS and 2D MHD

	2D MHD	2D CHNS
Magnetic Potential	A	ψ
Magnetic Field	\mathbf{B}	\mathbf{B}_ψ
Current	j	j_ψ
Diffusivity	η	D
Interaction strength	$\frac{1}{\mu_0}$	ξ^2

- 2D CHNS Equations:

$$\partial_t \psi + \vec{v} \cdot \nabla \psi = D \nabla^2 (-\psi + \psi^3 - \xi^2 \nabla^2 \psi)$$

$$\partial_t \omega + \vec{v} \cdot \nabla \omega = \frac{\xi^2}{\rho} \vec{B}_\psi \cdot \nabla \nabla^2 \psi + \nu \nabla^2 \omega$$

$-\psi$: Negative diffusion term

ψ^3 : Self nonlinear term

$-\xi^2 \nabla^2 \psi$: Hyper-diffusion term

With $\vec{v} = \hat{z} \times \nabla \phi$, $\omega = \nabla^2 \phi$, $\vec{B}_\psi = \hat{z} \times \nabla \psi$, $j_\psi = \xi^2 \nabla^2 \psi$. $\psi \in [-1, 1]$.

- 2D MHD Equations:

$$\partial_t A + \vec{v} \cdot \nabla A = \eta \nabla^2 A$$

$$\partial_t \omega + \vec{v} \cdot \nabla \omega = \frac{1}{\mu_0 \rho} \vec{B} \cdot \nabla \nabla^2 A + \nu \nabla^2 \omega$$

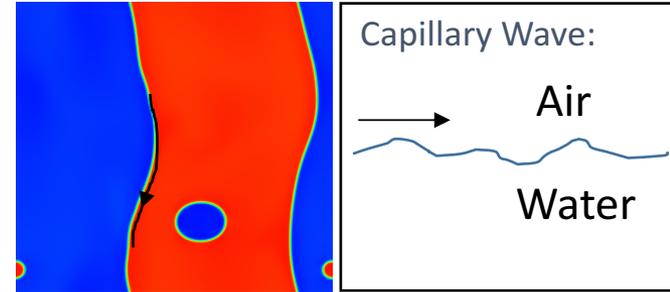
A : Simple diffusion term

With $\vec{v} = \hat{z} \times \nabla \phi$, $\omega = \nabla^2 \phi$, $\vec{B} = \hat{z} \times \nabla A$, $j = \frac{1}{\mu_0} \nabla^2 A$

Linear Wave

- CHNS supports linear “elastic” wave:

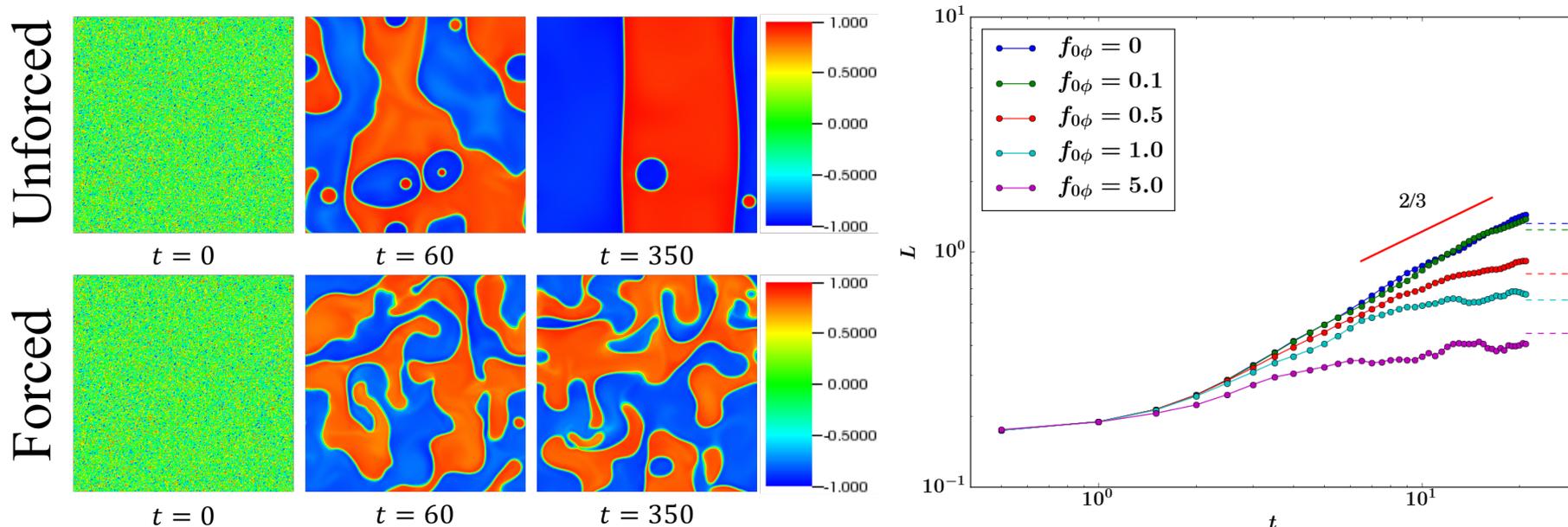
$$\omega(k) = \pm \sqrt{\frac{\xi^2}{\rho}} |\vec{k} \times \vec{B}_{\psi_0}| - \frac{1}{2} i(CD + \nu)k^2$$



Where $C \equiv [-1 - 6\psi_0 \nabla^2 \psi_0 / k^2 - 6(\nabla \psi_0)^2 / k^2 - 6\psi_0 \nabla \psi_0 \cdot i\mathbf{k} / k^2 + 3\psi_0^2 + \xi^2 k^2]$

- Akin to capillary wave at phase interface.
- Propagates ***only*** along the interface of the two fluids, where $|\vec{B}_{\psi}| = |\nabla \psi| \neq 0$.
- Analogue of Alfvén wave.
- Important differences:
 - \vec{B}_{ψ} in CHNS is large only in the interfacial regions.
 - Elastic wave activity does not fill space.

Results on CHNS Turbulence Study

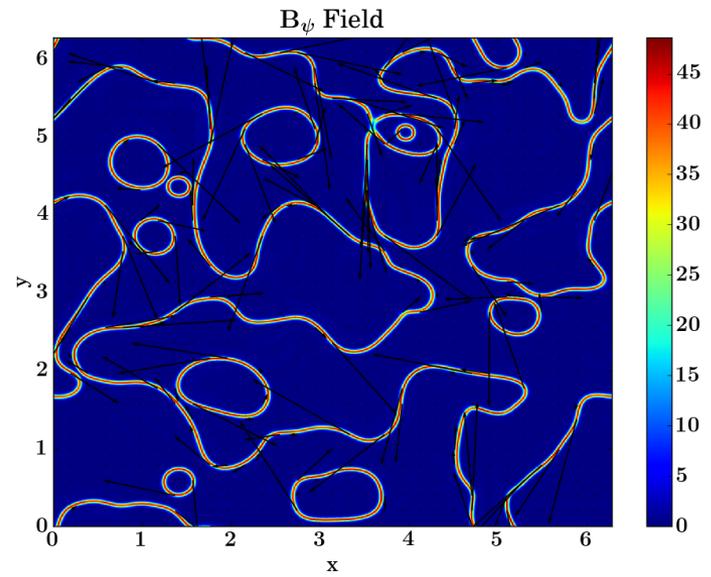
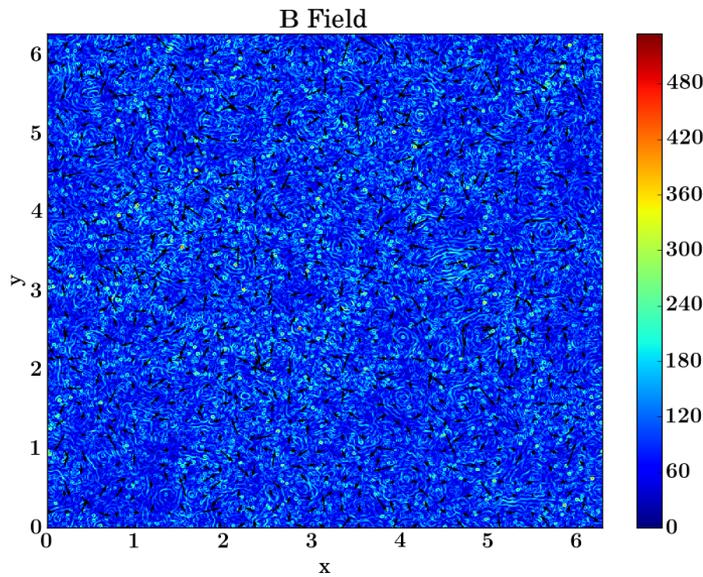


- Fluid forcing \rightarrow Fluid straining vs Blob coalescence
- Scale where turbulent straining \sim elastic restoring force (due surface tension): Hinze Scale

$$L_H \sim \left(\frac{\rho}{\xi}\right)^{-1/3} \epsilon_{\Omega}^{-2/9}$$

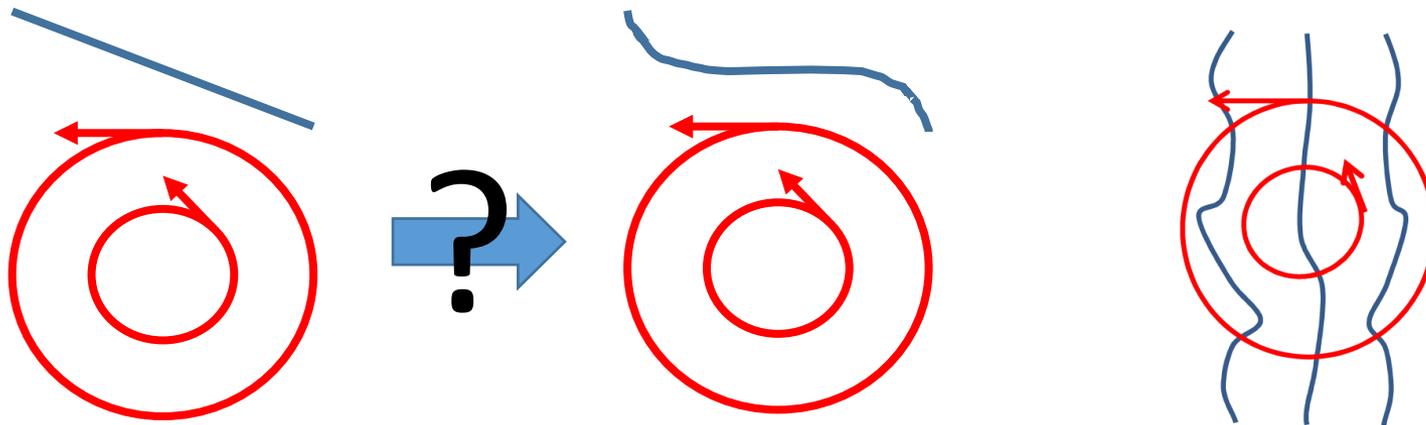
Results on CHNS Turbulence Study

- Dual cascade:
 - Inverse cascade of $\langle \psi^2 \rangle$
 - Forward cascade of E
- $k^{-7/3}$ for $\langle \psi^2 \rangle$ spectrum
- k^{-3} for kinetic energy spectrum (not $-3/2!$)
- Interface packing matters.



Why Single Eddy Mixing?

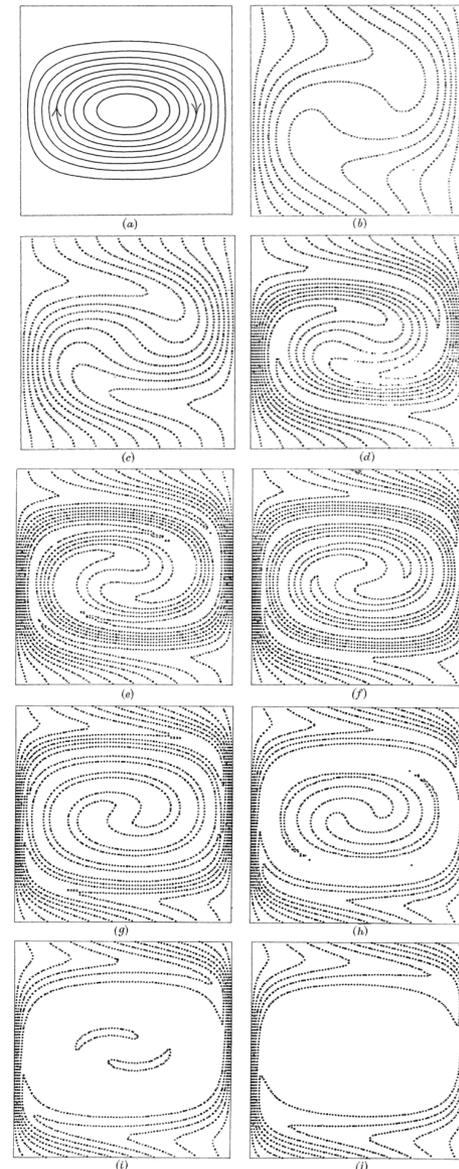
- Structures are the key \rightarrow need understand how a single eddy interacts with ψ field
- Mixing of $\nabla\psi$ by a single eddy \rightarrow characteristic time scales?
- Evolution of structure?
- Analogous to flux expulsion in MHD (Weiss, '66)



Single Eddy Mixing in 2D MHD: Expulsion

- When a convection eddy is imposed in a weak magnetic field, the magnetic field is expelled and amplified outside the eddy.
- This is called flux expulsion.
- The equation (kinematic, i.e. back reaction is ignored):

$$\partial_t A + \vec{v} \cdot \nabla A = \eta \nabla^2 A$$
- Also relevant to PV homogenization
→ ZF



Weiss 1966

Single Eddy Mixing in 2D MHD: Expulsion

- Main results of Weiss 1966 on Expulsion:
 - The final value of $\langle B^2 \rangle$ can be estimated by $\langle B^2 \rangle \sim Rm^{1/2} B_0^2$
 - The time for $\langle B^2 \rangle$ to reach a steady state is $\tau \sim Rm^{1/3} \tau_0$

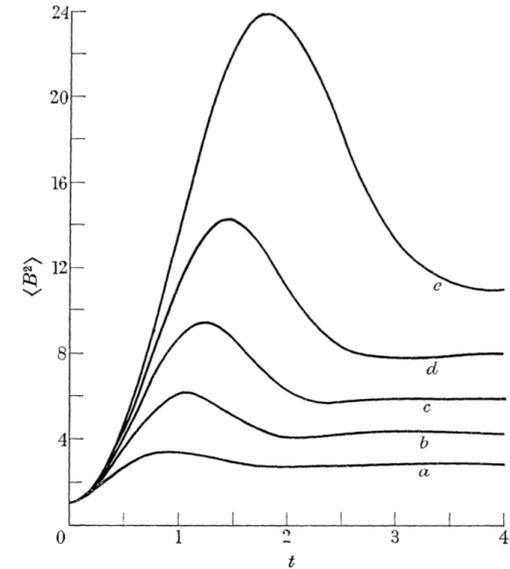


FIGURE 5. Magnetic energy as a function of time. Curves labelled *a, b, c, d, e* have $R_m = 40, 100, 200, 400, 1000$ respectively.

Weiss 1966

- Main results of Rhines and

Young 1983 on PV Homogenization:

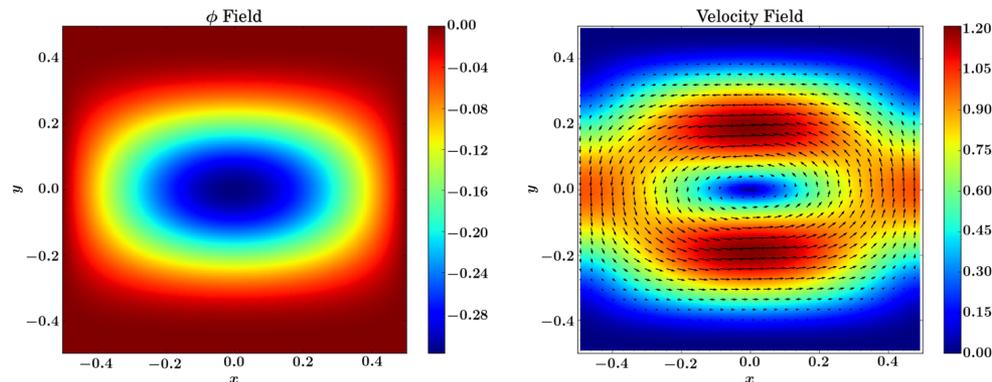
- Two stages: rapid and slow
- Rapid stage: dominated by shear-augmented diffusion, with time scale $\tau_{mix} \sim Rm^{1/3} \tau_0$
- Slow stage: usual diffusion, with time scale $\tau_{slow} \sim Rm \tau_0$

Single Eddy Mixing in 2D Cahn-Hilliard Flow

- The equation (kinematic, i.e. back reaction is ignored):

$$\partial_t \psi + \vec{v} \cdot \nabla \psi = D \nabla^2 (-\psi + \psi^3 - \xi^2 \nabla^2 \psi)$$

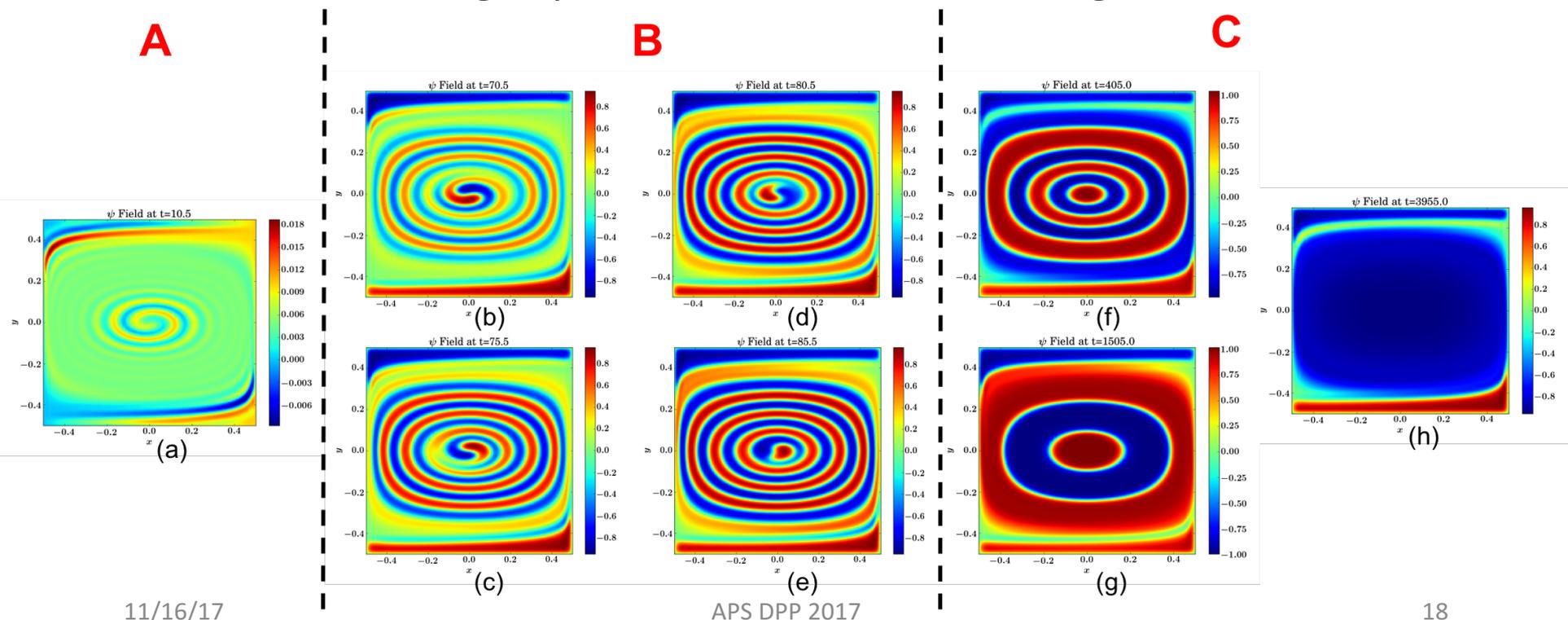
- Initial condition (analogous to Weiss 1966): \vec{B}_0 + eddy:



- Boundary condition: Dirichlet boundaries in both directions: $\psi = \psi_0$ and $\nabla^2 \psi = \nabla^2 \psi_0$ at boundaries.
- Dimensionless numbers:
 - $Pe = L_0 v_0 / D$ ($Pe \leftrightarrow Rm$)
 - $Ch = \xi L_0$

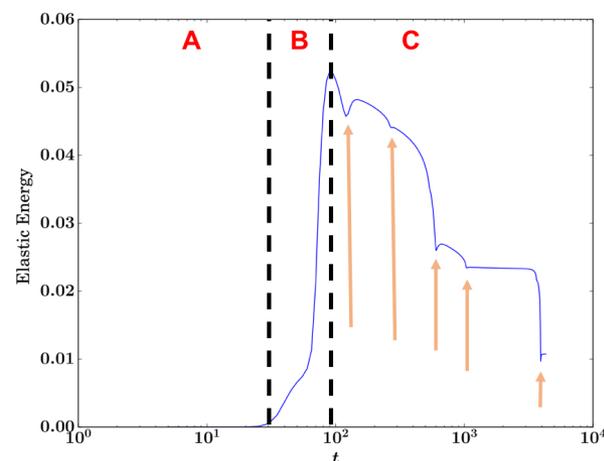
3 Stages

- 3 stages: (A) the “jelly roll” stage, (B) the *topological evolution* stage, and (C) the *target pattern* stage.
- ψ ultimately homogenized in slow time scale, but metastable target patterns formed and merge.

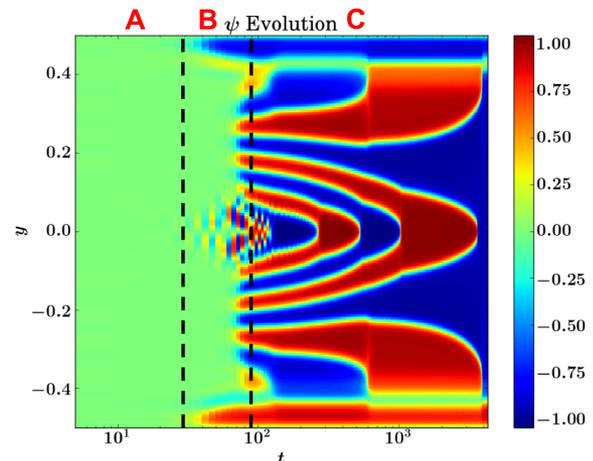


Band Merger

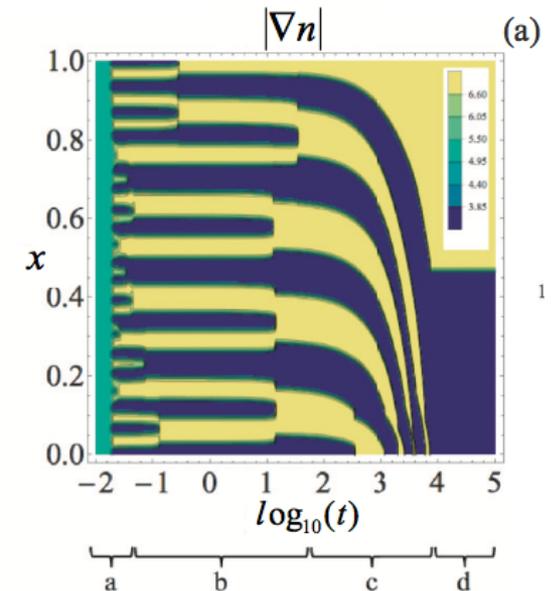
- The bands merge on a time scale **exponentially long** relative to eddy turnover time.
- The 3 stages are reflected in the elastic energy plot.
- The target bands mergers are related to the dips in the target pattern stage.
- The band merger process is similar to the step merger in drift-ZF staircases.



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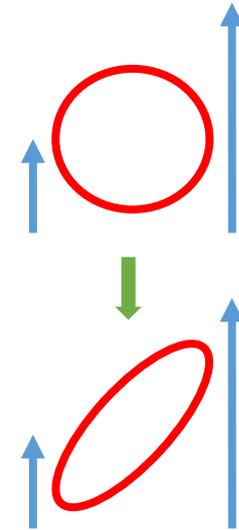


Ashourvan et al. 2016

Time Scales

- Analogous to the $Rm^{1/3}$ time scale in MHD, the mixing time scale of the shear + dissipation hybrid case is $\tau_{mix} \sim Pe^{1/5} Ch^{-2/5} t_0$.
- Brief derivation:
 - CH equation $\rightarrow \langle \delta r^4 \rangle \sim D \xi^2 t$
 - Relate δr and δy according to shear s :

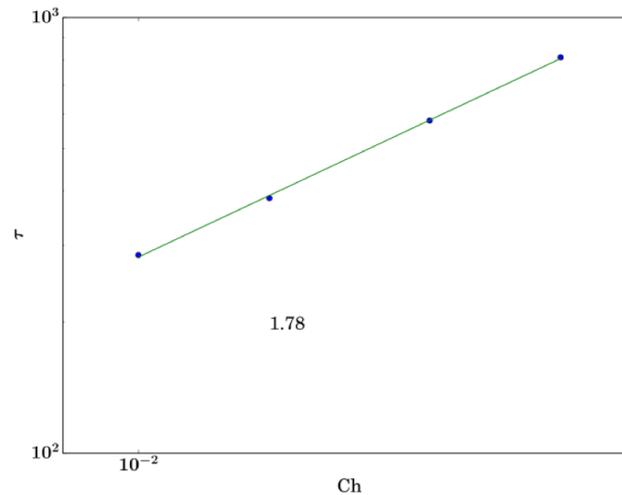
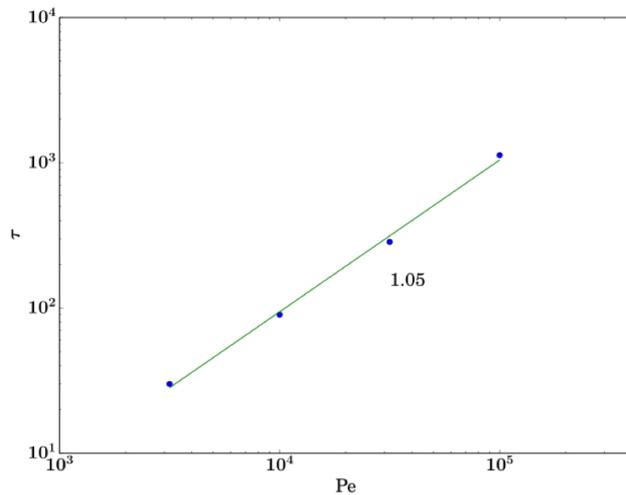
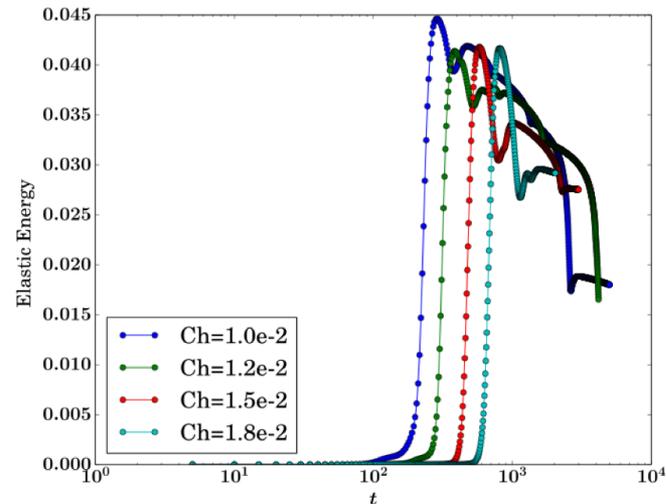
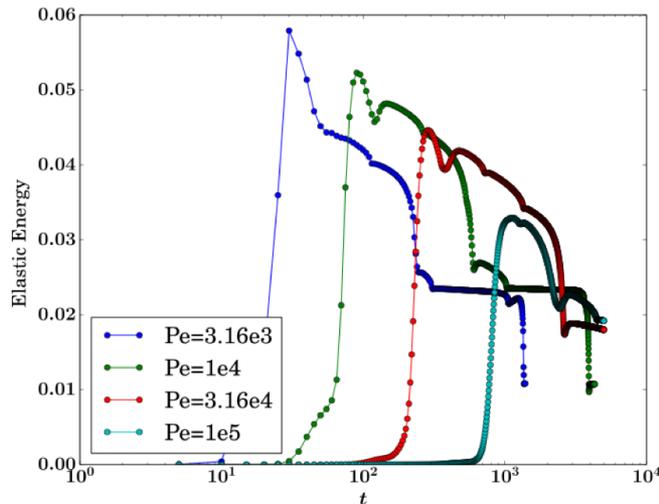
$$\frac{d}{dt} \delta y \sim s \delta r$$
 - So $\langle \delta y^4 \rangle \sim s^4 D \xi^2 t^5$.
 - Note that $Pe \sim L_y v / D$
 - So $\tau_{mix} \sim Pe^{1/5} Ch^{-2/5} t_0$



Time Scales

- Time to reach the maximum elastic energy:

$$\tau_m \sim PeCh^2 t_0$$



Conclusions

- Even kinematic single eddy mixing can exhibit unexpected nontrivial dynamics.
- 3 stages: (A) the “*jelly roll*” stage, (B) the *topological evolution* stage, and (C) the *target pattern* stage.
- Band merger process occurs on a time scale exponentially long relative to the eddy turnover time.
- Band merger process resembles step merger in drift-ZF staircases.
- Multi time-scale process: the $Pe^{1/5}$ and the Pe^1 time scale.