

Studies of Turbulence-driven FLOWs:

a) V_{\perp} , V_{\parallel} Competition in a Tube

b) Revisiting Zonal Flow Saturation

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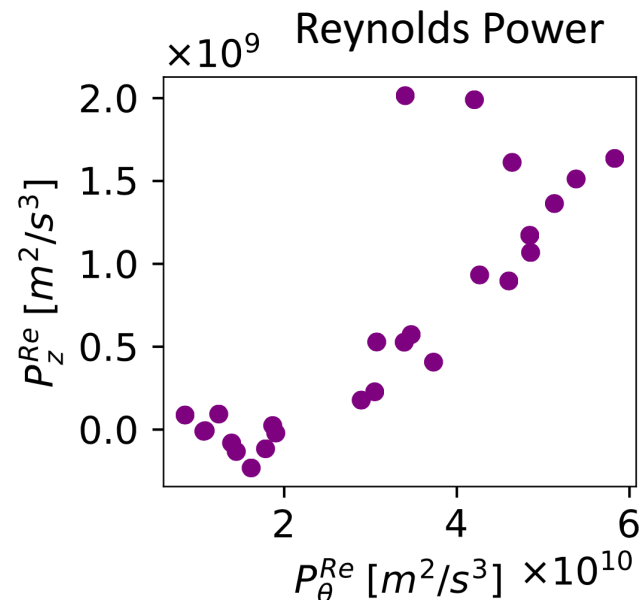
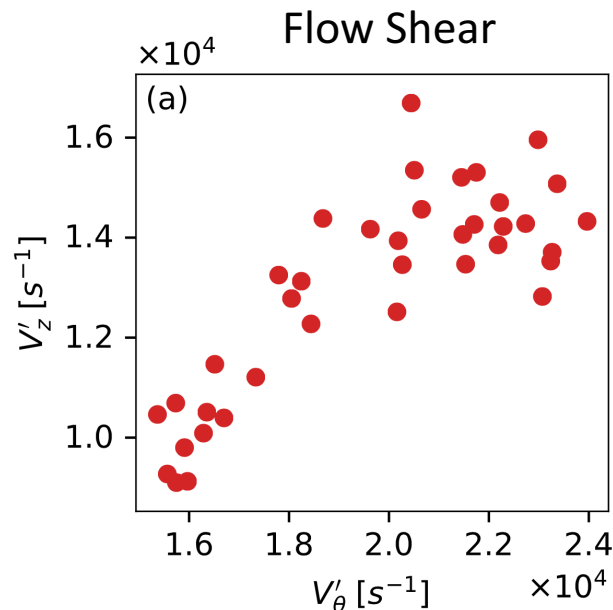
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Outline

- Background: turbulence driven V_{\perp} and V_{\parallel} observed in CSDX
- Questions: How do they interact? How do they saturate?
- Increment study of V_{\perp} , V_{\parallel} competition
 - Analogous to perturbation experiments
- Zonal flow saturation by wave-flow resonance
 - Wave-flow resonance effects on linear stability
 - ***Flow shear enhances instability via resonance***
 - Collisionless ZF saturation by resonance
 - ***Derives mesoscopic ZF scale, i.e. $L_{ZF} \sim \sqrt{\rho_s L_n}$***
 - Extended predator-prey model, compared to old model

Background

- **Intrinsic** axial and azimuthal flows observed in linear device (CSDX)
 - Increase $B \rightarrow$ scans mean flows-both V_{\perp} and V_{\parallel}
- **Dynamical** competition between perpendicular and parallel flows
- V_{\perp} and V_{\parallel} exchange energy with the turbulence, and each other.
 - \rightarrow Study energy apportionment between V_{\perp} and V_{\parallel}
 - \rightarrow Tradeoff between V_{\perp} and V_{\parallel}



Key Questions

- What's the coupling between *mean* perpendicular and parallel flows (V_{\perp} and V_{\parallel})?
 - How do they compete for energy from turbulence?
- How/Why do flows saturate, especially in *collisionless* regime?
- Why?
 - Linear device (CSDX) studies suggest apportionment of turbulence energy between V_{\perp} and V_{\parallel}
 - Relevant to L-H transition
 - Both V'_{\perp} and V'_{\parallel} increase, during transition.
 - The coupling of the two is relevant to transition threshold and dynamics.

Current status of coupling model

- Conventional wisdom of $V_{\perp} \rightarrow V_{\parallel}$ coupling:
 - V_{\perp}' breaks the symmetry in k_{\parallel} , but requires finite magnetic shear
 - **Not applicable** in linear device (straight magnetic field)
- $V_{\parallel} \rightarrow V_{\perp}$ coupling via parallel compression:
 - 3D coupled drift-ion acoustic wave system [Wang et al, PPCF 2012]
 - Coupling between fluctuating PV and parallel compression $\langle \tilde{q} \nabla_{\parallel} \tilde{v}_{\parallel} \rangle$
breaks PV conservation
 - Sink/source for fluctuating potential enstrophy density
 - Zonal flow generation

V_{\perp} and V_{\parallel} competition

- Increment study of V_{\perp} and V_{\parallel} effects on Reynolds powers
 - ***Turbulent energy branching*** between V_{\parallel} and V_{\perp}
 - Reynolds power ratio $P_{\parallel}^R / P_{\perp}^R$ decreases as V_{\perp} increases
 - tradeoff between V_{\perp} and V_{\parallel}
 - $P_{\parallel}^R / P_{\perp}^R$ maximum occurs when $|\nabla V_{\parallel}|$ is below the PSFI (parallel shear flow instability) threshold
 - saturation of intrinsic V_{\parallel}

Exploration of V_{\perp} - V_{\parallel} Coupling

- Goal: *How do extrinsic flows affect powers?*
 - **Turbulent energy branching** between intrinsic V_{\perp} and V_{\parallel}
 - How does V_{\perp} affect intrinsic V_{\parallel} generation?
 - Analogous to perturbation experiments, i.e. fix one flow and increase the other through external momentum source
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- Collisional drift wave
 - near adiabatic electron:

$$\tilde{n} = (1 - i\delta)\phi, \delta \ll 1$$

- Slab geometry

$$\frac{D}{Dt}\tilde{n} + \tilde{v}_r \frac{\nabla n_0}{n_0} + \nabla_{\parallel}\tilde{v}_{\parallel} = D_{\parallel}\nabla_{\parallel}^2(\tilde{n} - \tilde{\phi}),$$

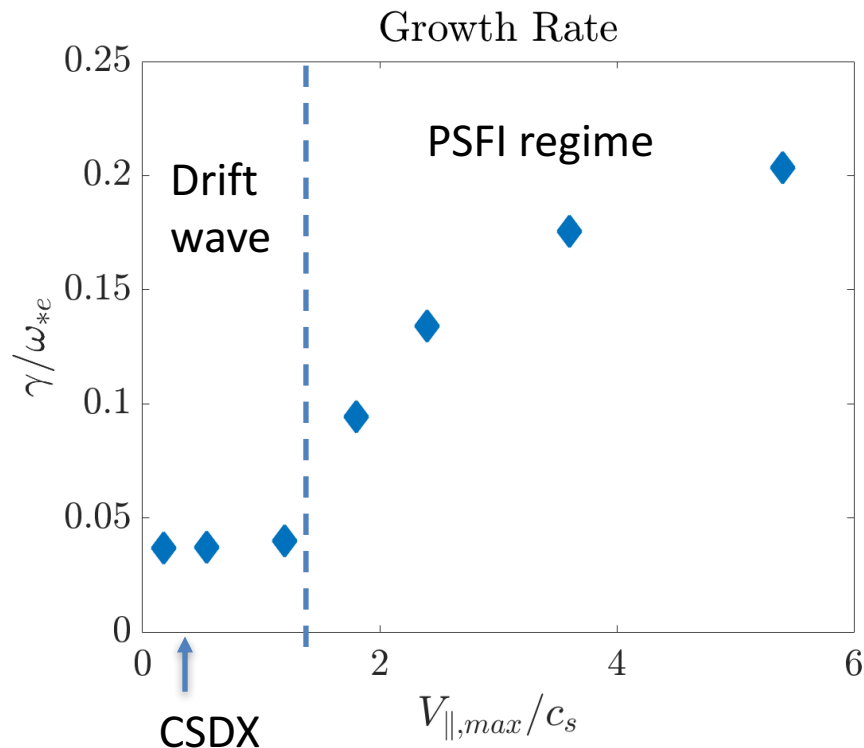
$$\frac{D}{Dt}\nabla_{\perp}^2\tilde{\phi} + \tilde{v}_r V_{\perp}'' = D_{\parallel}\nabla_{\parallel}^2(\tilde{n} - \tilde{\phi}),$$

$$\frac{D}{Dt}\tilde{v}_{\parallel} + \tilde{v}_r V_{\parallel}' = \nabla_{\parallel}\tilde{n},$$

∇n_0 is the Primary Instability Drive

- Other potential drives:
 - $V_{\perp}'' \rightarrow$ Kelvin-Helmholtz instability
 - $\nabla V_{\parallel} \rightarrow$ Parallel shear flow instability

- KH is not important
 - V_{\perp}'' drive weaker than ∇n_0 drive, i.e. $|k_y \rho_S^2 V_{\perp}''| \ll \omega_{*e}$

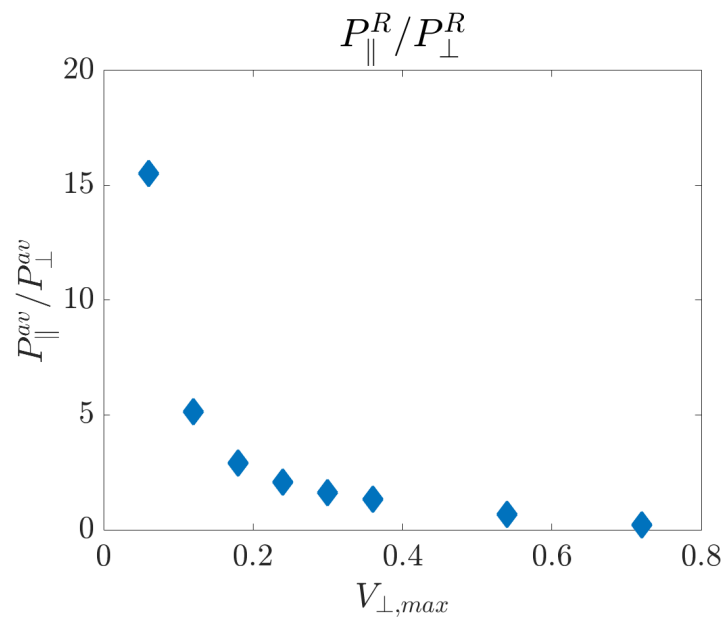


- ∇V_{\parallel} in CSDX is well below the PSFI linear threshold
 \rightarrow **PSFI stable** in CSDX

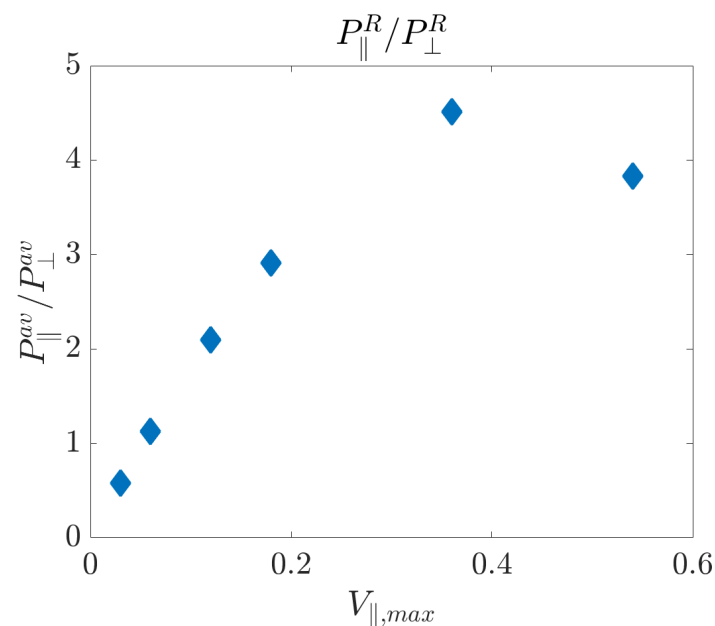


Coupling of V_{\perp} and V_{\parallel} \leftrightarrow Ratio of Reynolds Powers

- Ratio $P_{\parallel}^R / P_{\perp}^R$ decreases with V_{\perp}
 - Energy branching of V_{\parallel} reduced
 - V_{\perp} reduces generation of V_{\parallel}
 - Suggest **competition** between V_{\perp} and V_{\parallel}



- Increase $V_{\parallel} \rightarrow P_{\parallel}^R / P_{\perp}^R$ turnover **before** ∇V_{\parallel} hits PSFI threshold
 - Max energy branching of V_{\parallel} below PSFI threshold
 - Suggest V_{\parallel} saturates **below** PSFI threshold



Partial Summary 1

- CSDX experiments suggest **energy apportionment** between mean V_{\perp} and V_{\parallel}
- Increment study on Reynolds power ratio $P_{\parallel}^R / P_{\perp}^R$
 - Analogous to perturbation study
 - $P_{\parallel}^R / P_{\perp}^R$ decreases with $V_{\perp} \rightarrow$ tradeoff between V_{\perp} and V_{\parallel}
 - $P_{\parallel}^R / P_{\perp}^R$ maximum occurs **before** ∇V_{\parallel} hits PSFI threshold

Collisionless zonal flow saturation

- Wave-flow resonance prominent in linear device (CSDX)
 - Enters turbulence regulation, both linearly and nonlinearly
 - Flow shear is not the exclusive control parameter
- Resonance suppresses linear instability by wave absorption
 - Are shear suppression “rules” correct?
 - V'_{\perp} weakens resonance \rightarrow *flow shear enhances instability via resonance*

Collisionless zonal flow saturation (cont'd)

- Collisionless Zonal flow saturation by resonant PV mixing
 - Model of **resonant PV mixing**
 - Resonant diffusion of vorticity saturates zonal flow in collisionless regime
 - Incorporated in an extended predator-prey model
 - Drift wave mixes PV at zonal flow shear below that for KH/tertiary excitation

Wave-flow resonance effect on linear stability

- Resonance: $\omega_k - k_y V_\perp - k_\parallel V_\parallel$
 $|k_\parallel|/k_y \ll 1 \rightarrow$ Resonance set by $\omega_k - k_y V_\perp$
- Hasegawa-Wakatani drift wave model, with extrinsic V_\perp

$$\left(\frac{d}{dt} + \tilde{\mathbf{v}}_E \cdot \nabla \right) \tilde{n} + \tilde{v}_x \frac{\nabla n_0}{n_0} = D_\parallel \nabla_\parallel^2 (\tilde{n} - \tilde{\phi}) + D_c \nabla^2 \tilde{n},$$

$$\left(\frac{d}{dt} + \tilde{\mathbf{v}}_E \cdot \nabla \right) \tilde{\rho} + \tilde{v}_x \langle \rho \rangle' = D_\parallel \nabla_\parallel^2 (\tilde{n} - \tilde{\phi}) + \chi_c \nabla^2 \tilde{\rho},$$

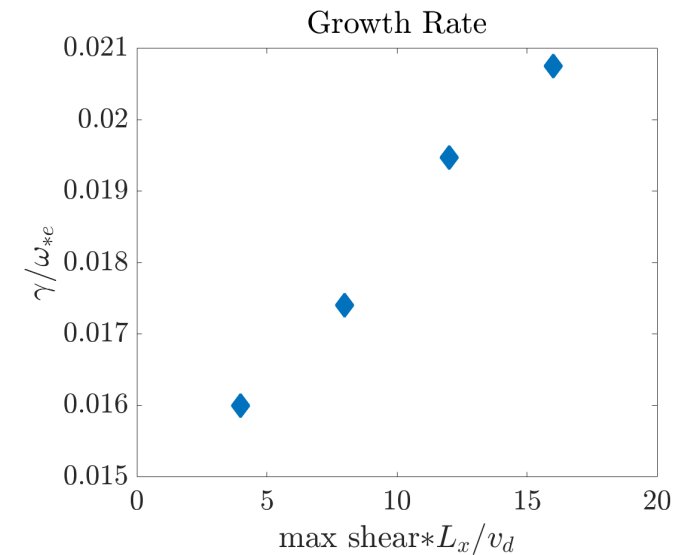
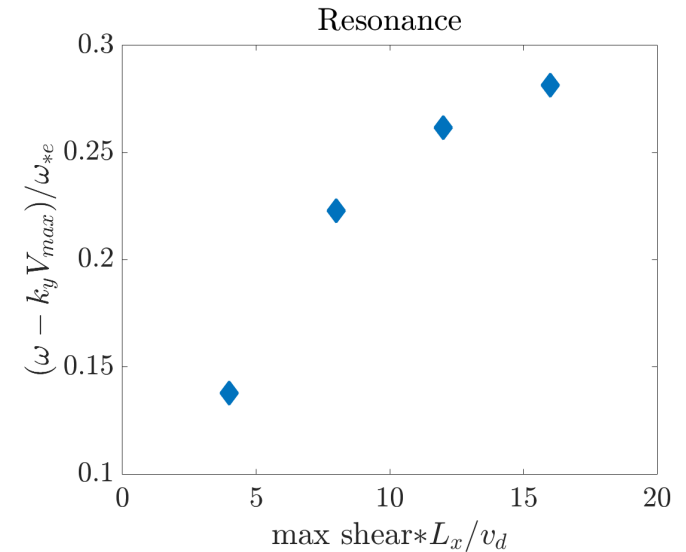
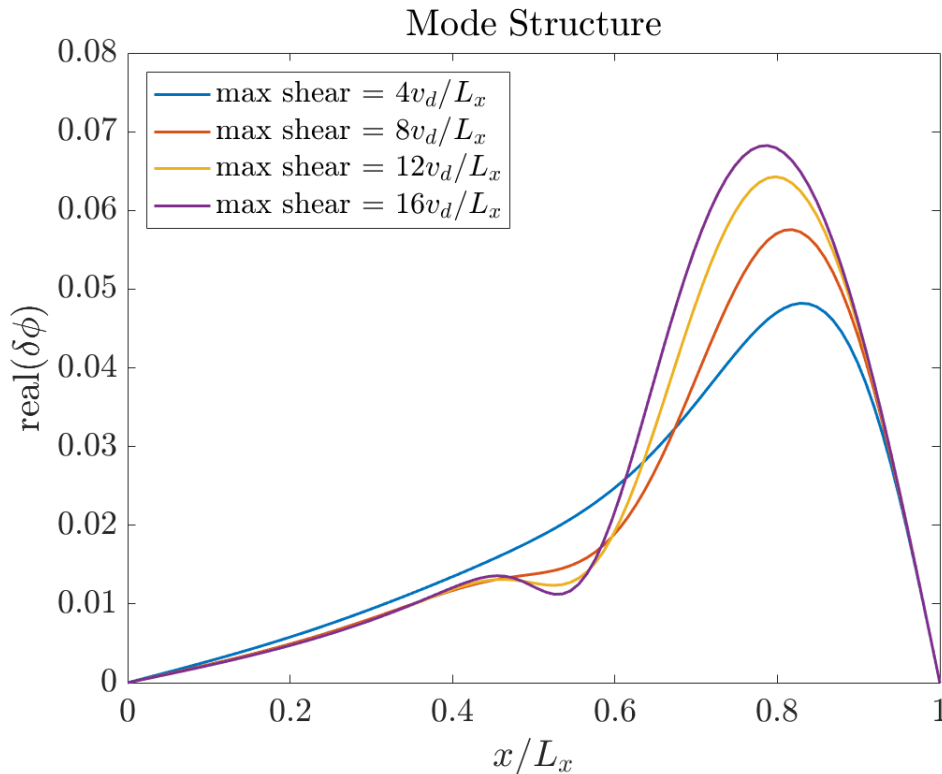
- **KH drive negligible**, i.e. $|k_y \rho_S^2 \langle v_y \rangle''| \ll \omega_{*e} \rightarrow$ Drift wave instability dominant
 - Near adiabatic electron: $\tilde{n} = (1 - i\delta)\phi$, $\delta \ll 1$
 - $\delta = (\omega_{*e} - \omega_k + k_y V_\perp) / k_\parallel^2 D_\parallel^2 = v_{ei} (\omega_{*e} - \omega_k + k_y V_\perp) / k_\parallel^2 v_{The}^2$

- Resonance reduces the **eigenmode scale** \rightarrow Suppresses instability

(Width of eigenmode)

Perpendicular flow shear *destabilizes* turbulence

- Mean perpendicular flow shear increases mode scale L_m/ρ_s
 - Weakens resonance
 - **Enhances instability**
- KH drive **negligible** compared to ∇n_0



Zonal Flow Saturation: Motivation

- Why?
 - Crucial to understand Dimits state physics
 - Collisionless zonal flow saturation, i.e. collisional damping $\rightarrow 0$
- Tertiary instability does not work
 - Severely damped by magnetic shear
 - Observed mean flow shear is always below the threshold for tertiary instability excitation

Nonlinear Model: *Resonant* PV Mixing

- Density:
$$\frac{\partial}{\partial t} \langle n \rangle = \frac{\partial}{\partial x} D_{n,\text{turb}} \frac{\partial}{\partial x} \langle n \rangle + D_c \nabla^2 \langle n \rangle,$$
- Vorticity:
$$\frac{\partial}{\partial t} \langle \rho \rangle = \frac{\partial}{\partial x} \left[(D_{n,\text{turb}} - D_q^{\text{res}}) \frac{\partial}{\partial x} \langle n \rangle + D_q^{\text{res}} \frac{\partial}{\partial x} \langle \rho \rangle \right] - \mu_c \langle \rho \rangle - \mu_{NL} \langle \rho \rangle + \chi_c \nabla^2 \langle \rho \rangle,$$
- PE:
$$\frac{\partial}{\partial t} \Omega = D_\Omega \frac{\partial}{\partial x} \Omega + D_q^{\text{res}} \left[\frac{\partial}{\partial x} (\langle n \rangle - \langle \rho \rangle) \right]^2 - \varepsilon_c \Omega^{3/2} + \gamma_L \Omega.$$

PE = Potential Enstrophy, i.e. $\Omega \equiv \langle \tilde{\rho}^2 \rangle$

- $\mu_{NL} = \mu_{NL}(\langle v_y \rangle)$: nonlinear damping rate ← driven by tertiary mode

Irrelevant to most cases we have encountered

- D_c, μ_c, χ_c : collisional particle diffusivity, flow damping, vorticity diffusivity → vanishing in collisionless regime

Resonant PV diffusion

- PV flux \rightarrow turbulent PV diffusion: $\langle \tilde{v}_x \tilde{q} \rangle = -D_{q,turb} \frac{\partial}{\partial x} \langle q \rangle$



$$D_{q,turb} = \text{Resonant} + \text{Non-resonant}$$

- Resonant PV diffusivity:

$$D_q^{res} = \sum_k |\tilde{v}_x|^2 \pi \delta(\omega_k - k_y V_\perp) \sim \sum_k \tau_{c,k} k_y^2 \rho_s^2 c_x^2 |\phi_k|^2$$

$$\tau_{c,k} \sim [|v_{g,y} - v_{ph,y}| \Delta k_y + v_{g,x} \Delta k_x]^{-1}$$

- Non-resonant PV diffusivity:

$$D_q^{non-res} = \sum_{\omega_k \neq k_y \langle v_y \rangle} k_y^2 \rho_s^2 c_s^2 |\phi_k|^2 \frac{|\gamma_k|}{|\omega_k - k_y \langle v_y \rangle|^2} \sim \sum_{\omega_k \neq k_y \langle v_y \rangle} \frac{k_y^2 \rho_s^2 c_s^2}{k_\parallel^2 D_\parallel} \frac{k_y^2 \rho_s^2 + L_m^{-2} \rho_s^2}{1 + k_y^2 \rho_s^2 + L_m^{-2} \rho_s^2} |\phi_k|^2$$

Resonant diffusivity exceeds non-resonant part:

$$D_q^{non} / D_q^{non-res} \sim \tau_{c,k} k_\parallel^2 v_{The}^2 / \nu_{ei} \gg 1$$

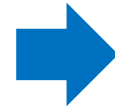
Collisionless saturation by resonant diffusion of vorticity

- Zonal flow evolution ← Mean enstrophy equation:

$$\frac{\partial}{\partial t} \int dr \frac{\langle \rho \rangle^2}{2} = \int dr \langle \tilde{v}_r \tilde{\rho} \rangle \frac{d\langle \rho \rangle}{dr} - \nu_i \int dr \langle \rho \rangle^2 + \dots$$

Vorticity ($\rho \equiv \nabla_{\perp}^2 \phi$) flux:

$$\langle \tilde{v}_r \tilde{\rho} \rangle = -D_q^{res} \frac{d\langle \rho \rangle}{dr} + \Gamma_{\rho}^{Res}$$



Conserves enstrophy between mean flow and fluctuations



$$\frac{\partial}{\partial t} \int dr \frac{\langle \rho \rangle^2}{2} = - \int dr D_q^{res} \left(\frac{d\langle \rho \rangle}{dr} \right)^2 + \int dr \Gamma_{\rho}^{Res} \frac{d\langle \rho \rangle}{dr} - \nu_i \int dr \langle \rho \rangle^2 + \dots$$

- $\nu_i \rightarrow 0 \rightarrow$ Dimits shift regime \rightarrow **Resonant diffusion saturates ZF**
- Collisionless damping by turbulent viscosity: $d\langle \rho \rangle / dr \sim \Gamma_{\rho}^{Res} / D_q^{res}$
 - Resonant vorticity diffusivity $D_q^{res} \rightarrow$ ZF saturation

Mesosopic stationary zonal flow

- Balance vorticity flux: $\langle \tilde{v}_x \tilde{\rho} \rangle = -D_q^{res} \frac{d\langle \rho \rangle}{dx} + \Gamma_\rho^{Res} = 0$

$$\rightarrow \langle v_y \rangle'' = d\langle \rho \rangle / dx \sim \Gamma_\rho^{Res} / D_q^{res}$$

- Vorticity flux driven by ∇n : $\Gamma_\rho^{Res} = (D_{n,turb} - D_q^{res}) \frac{\partial}{\partial x} \langle n \rangle$

- Resonant PV diffusivity:

$$D_q^{res} = \sum_k \tau_{c,k} k_y^2 \rho_s^2 c_x^2 |\phi_k|^2 \text{ with } \tau_{c,k} \sim [(v_{g,y} - v_{ph,y}) \Delta k_y + v_{g,x} \Delta k_x]^{-1}$$

- Stationary flow:

$$\langle v_y \rangle'' = \langle \rho \rangle' = \left(1 - \frac{D_{n,turb}}{D_q^{res}} \right) \frac{\partial \langle n \rangle}{\partial x} \sim -\frac{c_s}{\rho_s L_n} \left(1 - \frac{1}{\tau_{c_k} k_{||}^2 D_{||}} \frac{k_y^2 \rho_s^2 + L_m^{-2} \rho_s^2}{1 + k_y^2 \rho_s^2 + L_m^{-2} \rho_s^2} \right)$$

$$\rightarrow \text{Zonal flow scale: } L_{ZF} \sim \sqrt{\rho_s L_n} \rightarrow \rho_s \ll L_{ZF} \ll L_n \quad \downarrow$$

This *derives* the standard ordering, which is just invoked, in ad hoc way.

Extended Predator-Prey Model

- Mean flow energy:

$$\frac{L_{ZF}^2}{2} \frac{dV''^2}{dt} = \alpha_1 |V''| E - \alpha_2 V''^2 E - \gamma_{NL} V''^2 - \mu_c V''^2.$$

new

L_{ZF} : zonal flow profile scale, $\rho_s \ll L_{ZF} \ll L_n$

- Turbulence energy (PE):

$$\frac{dE}{dt} = -\alpha_1 |V''| E + \alpha_2 V''^2 E - \varepsilon_c E^{3/2} + \gamma_L E.$$



Turbulence and flow states

- Compare by regime:

Regime	Collisionless	Weak Collisional	Strong Collisional
Collisional Damping Strength	$\mu_c \ll \alpha_2 E$	$\alpha_2 E \ll \mu_c \ll 4\gamma_L \alpha_1^2 / \varepsilon_c^2$	$\mu_c \gg 4\gamma_L \alpha_1^2 / \varepsilon_c^2$
Flow State	$\frac{\alpha_1}{\alpha_2}$	$\frac{\alpha_1 \gamma_L^2}{\mu_c \varepsilon_c^2}$	$\frac{\gamma_L}{\alpha_1}$
Turbulence Energy	$\frac{\gamma_L^2}{\varepsilon_c^2}$	$\frac{\gamma_L^2}{\varepsilon_c^2}$	$\frac{\gamma_L \mu_c}{\alpha_1^2}$

- Collisionless = collisional damping/viscosity $\rightarrow 0$
- Collisionless saturation compared to usual collisional damping:
 - Turbulence energy determined by linear stability and small scale dissipation
 \rightarrow Different from usual models, where turbulence energy \sim flow damping
 - Flow state basically independent of stability drive
 \rightarrow There can be flows in nearly marginal turbulence

Analogy to Landau Damping Absorption in Langmuir Turbulence

	Langmuir Turbulence Collapse	Collisionless ZF Saturation
Primary player	Plasmon-Langmuir wave	Drift wave turbulence
Secondary player	Ion- acoustic wave (caviton)	Zonal flow
Free energy source	Langmuir turbulence driver	$\nabla n, \nabla T$ drives
Final State	(Nearly) empty cavity	Saturated zonal flow and <i>residual</i> turbulence
Resonance	Landau damping	Resonant wave absorption
Other damping effects	Ion-acoustic radiation	Kelvin-Helmholtz relaxation

- Landau damping: flattens PDF (negative slope) in phase space
- Resonant PV mixing: homogenizes mean PV in real space

Partial Summary 2

- Resonance effects on linear stability
 - Wave-flow resonance suppresses instability
 - V'_\perp weakens resonance $\rightarrow V'_\perp$ **enhances** instability via resonance
- Resonant diffusion of vorticity saturates zonal flow in collisionless regime
 - Resonant PV mixing \leftarrow resonant diffusion of PV
 - Model shows that stationary zonal flow scale is mesoscopic, i.e. $\rho_s \ll L_{ZF} \ll L_n$, since $L_{ZF} \sim \sqrt{\rho_s L_n}$
 - Extended predator-prey model
 - \rightarrow **turbulence energy** $\sim \gamma_L^2 / \varepsilon_c^2$ not $\sim \gamma_L$
 - Flow independent of turbulence level/drive
 - \rightarrow **flow in marginal turbulence**