

What is the mechanism for the emergence and sustainment of staircase structure?

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***This research was supported by the U. S. Department of Energy, Office of Science, Office of Fusion Energy Sciences, under Award Number DE-FG02-04ER54738**

***This project was also supported by Graduate School, HUST and NSFC Grant Nos. 11675059 and 11305071**

Gratitude

- **Mentors: P. H. Diamond and Lu Wang**
- **Teachers and fans during courses: Plasma Physics+Fluid Dynamics** 
- **Friends---in academic and daily life.**
 - ✓ **A. Ashourvan, Jiacong Li, Xiang Fan and R. J. Hajjar;**
 - ✓ **Qiming Hu, Jie Chen and Huiqian Wang.**

I. Research

■ Background

- **Patterns formation in Drift wave-Zonal flow turb.**
 - Staircase vs other patterns
- **Mechanism: feedback loops**
 - Rhines vs Shearing

- ### ■ Formation of staircase is sensitivity to
- ✓ Parameters
 - ✓ B.C.
 - ✓ Initial condition

■ Conclusions so far and future pursuing

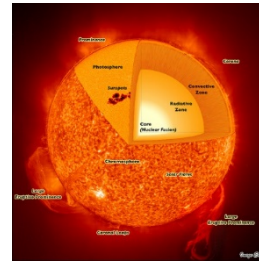
II. Experiences (live communication)

Suggestions on living in UCSD for the first time

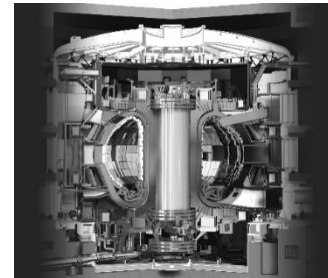
Background

- Self-organizing, non-equilibrium nonlinear systems, formation of **patterns** is a common feature.

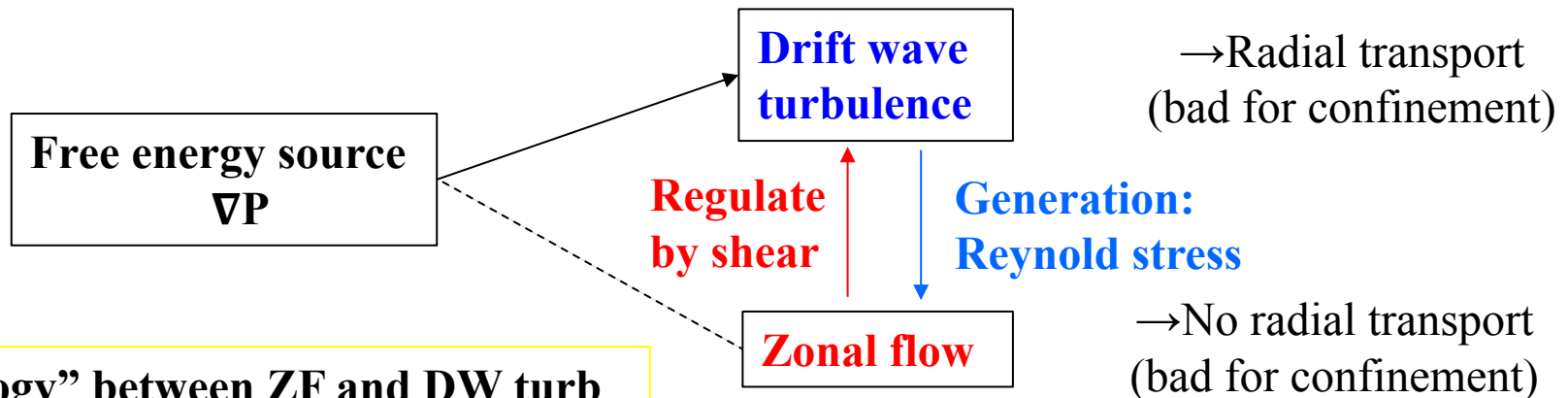
- ✓ Stratification layers
- ✓ Quasi-periodic flow
- ✓ Staircase
- ✓ $E \times B$ zonal flow (ZF) shear
- ✓



- ZF: regulate turbulent transport trigger L-H and ITBs



- Closing the feedback loop when **predators meet the preys**



“Ecology” between ZF and DW turb

What is the spatial structure?



■ Different spatial scale

- ✓ Drift wave, micro-scale ρ
 - ✓ Mean flow, macro-scale $L_n \sim a$
 - ✓ Zonal flow, meso-scale $\sqrt{\rho L_n}$
- ρ : Larmor radius
 - $L_n = -n/\nabla n$: density scale length

■ Predicting turbulence and transport in states evolving from saturated instability is the goal

→ Turbulent diffusivity **D scaling with ρ/a** is important in fusion!

$$D = D_{\text{Bohm}} \left(\frac{\rho}{a} \right)^\alpha, \quad 0 < \alpha < 1. \quad \begin{cases} \alpha = 0, \text{ Bohm scaling (Bad)} \\ \alpha = 1, \text{ Gyro-Bohm scaling (Good)} \end{cases}$$

■ The ecology of feedback should have some effects on scaling of spatial structure

$$|\delta\phi| \sim l_{\text{mix}} \sim \frac{1}{1 + v'_{E \times B}}$$

Staircase pattern in drift wave



✓ In drift wave, the $E \times B$ staircase (meso-scale).

Q: $E \times B$ staircase = Zonal flow ?

A: $E \times B$ staircase is primarily produced by the zonal flow generation, and enhanced by the inphase mean flow variation.

✓ Quasi-regular (spatial) and long-lived (temporal) $E \times B$ flows with temperature corrugations coexist are observed numerically + experimentally

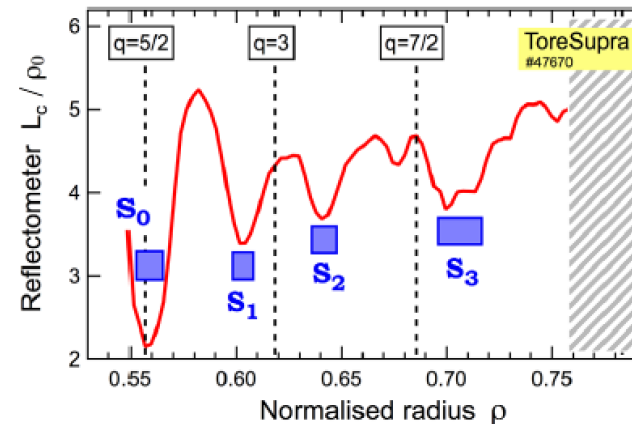
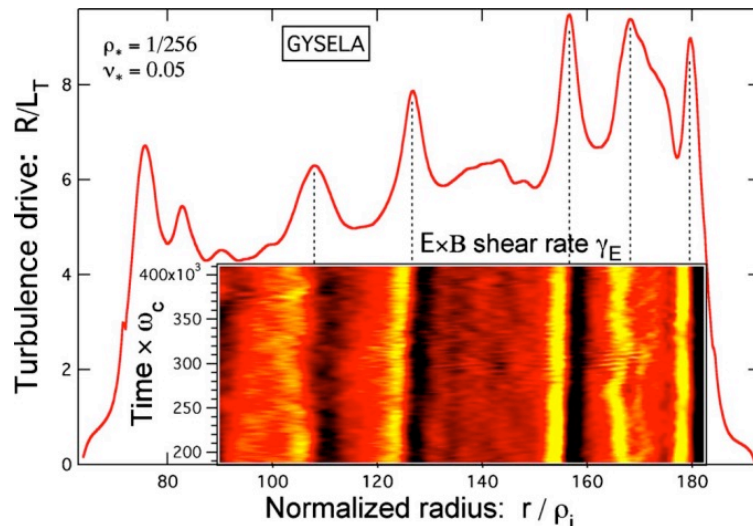


FIG. 3 (color online). The reflectometer coherence length plotted against radius shows clear experimental evidence of a staircase at locations S_1 , S_2 and S_3 , possibly also at S_0 .

Note: Coherent quasi-periodic shearing pattern !!!

[G. Dif-pradalier, *et al*, Physical Rev. E **82**, 025401 (2010)]
 [G. Dif-pradalier, *et al*, , Phys. Rev. Lett. **114**, 085004 (2015).]

Beauty of staircase pattern

- ✓ $E \times B$ staircase patterning \rightarrow segregating regions where **avalanching** is dominant from regions where **zonal flow concentrates** (forming the staircase microbarriers)

“A natural and dynamic means for the simultaneous existence of these two antagonistic trends”

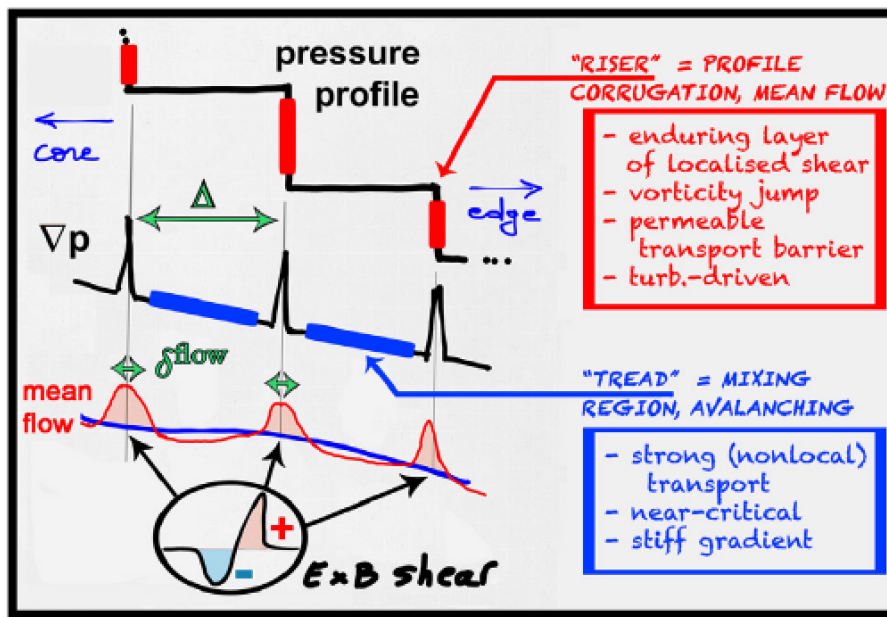


Figure 1. The $E \times B$ staircase, schematic view.

[G. Dif-Pradalier et al Nucl. Fusion 57 (2017) 066026]

- These zonal mean flows, of typical radial extent

$$\delta^{\text{flow}} \sim 10 \rho_s$$
- The most probable step size of avalanches

$$\Delta^{\text{stat}} \sim 40 \rho_s$$
- ∇T corrugations temporally-averaged over 0.53 ms between 1030 a/c_s and 1343 a/c_s

Mechanism for the emergence and sustainment of staircase structure?



Mechanism for the formation and sustainment of staircase structure from H-W equation



- **Reduced models** that self-consistently relate variations in mean plasma fields to fluctuation intensity !!! [A. Ashourvan and P. H. Diamond, PoP 2017, PRE 2016]

- ✓ lower computational cost, supply **essential** understanding.....

- Two popular drift wave models : Hasegawa–Mima (H–M) and Hasegawa–Wakatani (H–W) model [Hasegawa A and Mima K 1978 PoF 21 87]
[Hasegawa A and Wakatani M 1983 PRL. 50 682]

(1) H-M equation: adiabatic and collisionless limit

- ✓ 2D fluid, H-M → potential vorticity (PV) conservation along fluid trajectories.

$$q = n - u \quad \begin{array}{l} n : \text{normalized density} \\ u : \text{normalized vorticity} \end{array}$$

(2) H-W equation: collisional

- ✓ Usually neglect the dynamical coupling along the magnetic field line → 2D system describing the **collisional** drift wave instability driven turbulence, which conserves the energy and PE

$$(\text{PE: potential enstrophy } \varepsilon = q^2 / 2)$$

- ✓ Conservation of PE leads to the spontaneous generation of ZF by turbulence (Reynolds stress)

Derivation details about H-W equation and physical understanding



- In strong magnetic field $B=B_z\bar{e}_z$, assume cold ion, the **ion** force balance Eq.

$$\frac{d\bar{u}}{dt} = -\frac{e}{M}\nabla\phi + \frac{e}{M}\bar{u}\times\bar{B} - \cancel{\nabla p_i} - \nabla\Pi + \cancel{F}, \quad \nabla\Pi = \mu\nabla^2\bar{u}$$

- Treat the drift velocity order by order \rightarrow the lowest order is the $E\times B$ drift and the first order is the polarization drift and drift caused by viscosity

$$\mathbf{u}_0 = \mathbf{u}_E = -\nabla\phi \times \frac{\mathbf{B}}{B_0^2} \quad \mathbf{u}_1 = \mathbf{u}_p + \mathbf{u}_{visc} = -\frac{1}{\omega_{ci}B_0}\frac{d\nabla\phi}{dt} - \frac{\mu}{\omega_{ci}B_0}\nabla^2(\nabla\phi)$$

- From the divergence free condition $\nabla\cdot\bar{J}=0$ with $\bar{J}=\bar{J}_\perp+\bar{J}_\parallel$

$$\bar{J}_\perp = n|e|\bar{u}_1$$

$$\nabla_\parallel\phi + \eta\bar{J}_\parallel - \frac{1}{en_0}\nabla_\parallel p_e = 0$$

- Finally, the equation for evolution of $\nabla_\perp^2\phi$

$$\rho_s^2\frac{d}{dt}\nabla_\perp^2\phi = -D_\parallel\nabla_\parallel^2\left(\phi - \frac{n}{n_0}\right) + \nu\nabla_\perp^2\nabla_\perp^2\phi$$

Derivation details about H-W equation and physical understanding



- In the force balance equation for electron $\nabla \Pi = 0$ and treat $m_e n_e \frac{d\vec{u}}{dt} = 0$, but keeps electron-ion friction

$$\mathbf{F}_{e\parallel} = -m_e n_e \nu_{ei} \mathbf{u}_{e\parallel} = \frac{m_e \nu_{ei}}{e} \mathbf{J}_{\parallel}$$

- From the Ohm's law

$$\nabla_{\parallel} \phi + \eta \bar{\mathbf{J}}_{\parallel} - \frac{1}{en_0} \nabla_{\parallel} p_e = 0$$

and the electron continuity equation

$$\frac{dn}{dt} + n \nabla_{\parallel} \cdot \vec{u}_{e,\parallel} = 0$$



$$\frac{d}{dt} n - D_0 \nabla^2 n = -D_{\parallel} \nabla_{\parallel}^2 \left(\phi - \frac{n}{n_0} \right)$$



$$\rho_s^2 \frac{d}{dt} \nabla^2 \phi = -D_{\parallel} \nabla_{\parallel}^2 \left(\phi - \frac{n}{n_0} \right) + \nu \nabla^2 \nabla^2 \phi$$

**H-W
equations
!!!!**

Derivation details about H-W equation and physical understanding



- When $D_{\parallel} k_{\parallel}^2 / \omega \gg 1 \rightarrow \frac{\delta n}{n_0} = \frac{e\delta\phi}{T_e}$ i.e., electron is adiabatic. In the adiabatic and collisionless limit \rightarrow reduce to the **Hasegawa-Mima (H-M)** equation

$$\frac{d}{dt}(\phi - \rho_s^2 \nabla^2 \phi) + v_* \partial_y \phi = 0$$

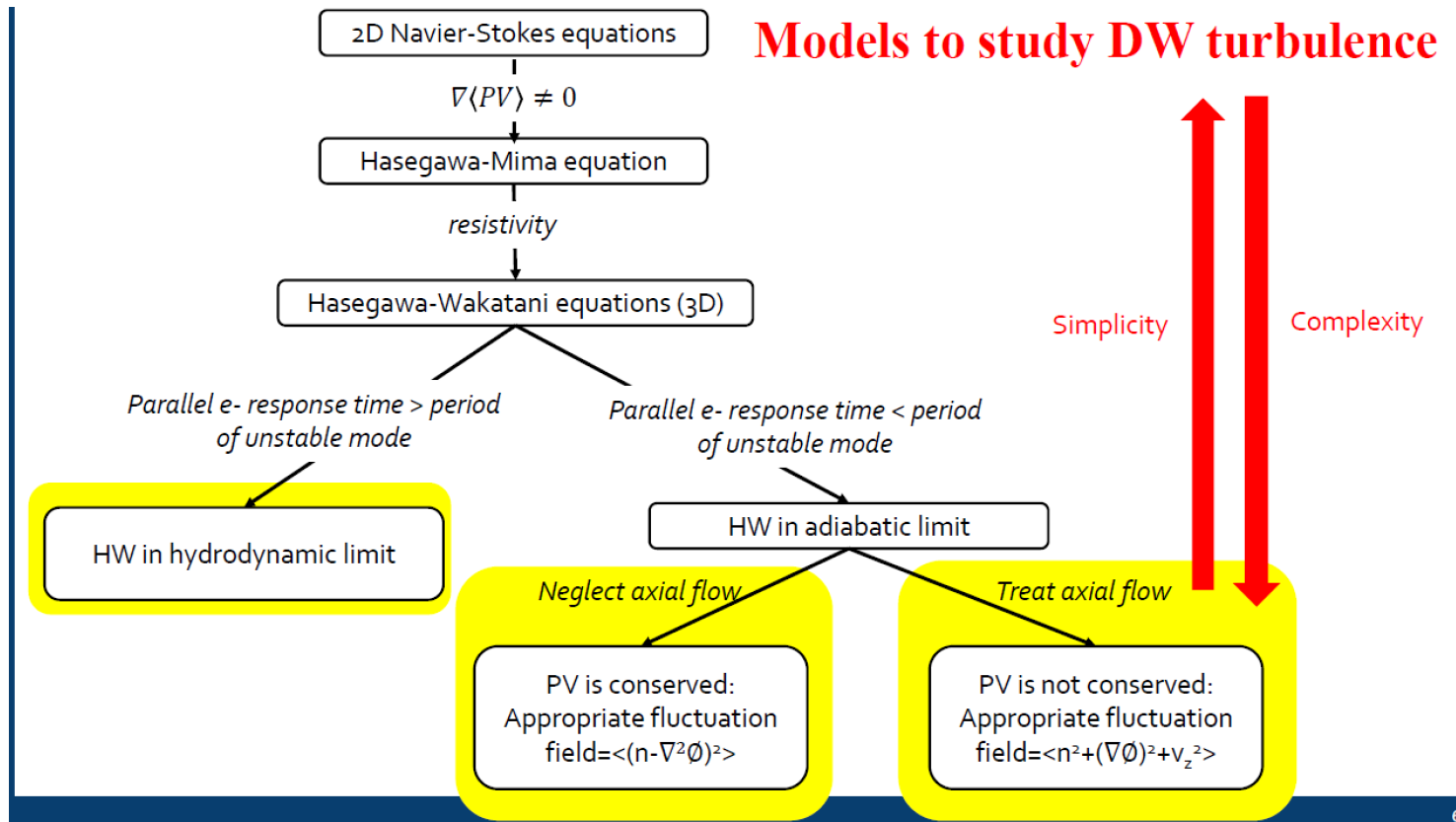
\rightarrow PV is conserved

- ✓ H-M model and the H-W) model are often used to investigate drift wave-ZF system
- ✓ H-W equation can be used to explain the generation of zonal-flow under the use of Taylor identity $\partial \langle v_r v_y \rangle / \partial r = \langle v_r \nabla_{\perp}^2 \phi \rangle$

$$\partial \langle v_y \rangle / \partial t = - \langle v_r \nabla_{\perp}^2 \phi \rangle - \nu \langle v_y \rangle$$

- ✓ The developing of this reduced model were widely investigated by professor Diamond

Examples: (1) works done recently by doctor Rima Hajjar



Examples: (2) works done recently by Jiacong Li

- ✓ How electron DW generate the intrinsic $v_{||}$ when $s = 0$
- ✓ The interaction between zonal flow and $v_{||}$

Bulid the reduced model to understand the mechanism of staircase based on H-W equation



■ The vorticity equation and the continuity equation

$$\left\{ \begin{array}{l} \frac{d}{dt} (\nabla^2 \varphi) = \eta \nabla_{\parallel}^2 (\log N - \varphi) + \mu_c \nabla_{\perp}^4 \varphi, \\ \frac{d}{dt} \log N = \eta \nabla_{\parallel}^2 (\log N - \varphi) + D_c \nabla_{\perp}^2 \log N. \end{array} \right. \quad u \equiv \nabla_{\perp}^2 \varphi,$$

$$\frac{d}{dt} = \frac{\partial}{\partial t} + \mathbf{v}_E \cdot \nabla, \quad \mathbf{v}_E = \hat{z} \times \nabla \varphi. \quad u = \langle u \rangle + \delta u(x, y, z, t) \quad \mathbf{v} = \langle \mathbf{v} \rangle \hat{y} + \delta \mathbf{v}(x, y, z, t),$$

■ Quasi-linear equations for the mean density n and vorticity u

$$\partial_t \langle u \rangle + \partial_x \langle \delta v_x \delta u \rangle - \mu_c \nabla_{\perp}^2 \langle u \rangle = 0, \quad (11) \quad \langle \delta v_x \delta u \rangle = (\chi - D_n) \partial_x \langle n \rangle - \chi \partial_x^2 \langle v \rangle$$

$$\partial_t \langle n \rangle + \partial_x \langle \delta v_x \delta n \rangle - D_c \nabla_{\perp}^2 \langle n \rangle = 0. \quad (12) \quad \langle \delta v_x \delta n \rangle = -D_n \partial_x \langle n \rangle + V_{\text{pinch}},$$

■ The fluctuation equations are obtained as

$$(\partial_t + \langle v \rangle \partial_y) \delta n + \{ \delta \varphi, \delta n \} + \delta v_x \partial_x \langle n \rangle - \mu_c \nabla_{\perp}^2 \delta n = \eta \nabla_{\parallel}^2 (\delta n - \delta \varphi),$$

$$(\partial_t + \langle v \rangle \partial_y) \delta u + \{ \delta \varphi, \delta u \} + \delta v_x \partial_x^2 \langle v \rangle - D_c \nabla_{\perp}^2 \delta u = \eta \nabla_{\parallel}^2 (\delta n - \delta \varphi).$$

■ Fluctuation PE $\varepsilon = \langle \delta q^2 \rangle / 2$. are obtained as

$$\partial_t \langle \delta q^2 \rangle + \partial_x \langle \delta v_x \delta q^2 \rangle = - \langle \delta v_x \delta q \rangle \partial_x \langle q \rangle - \mu_{\varepsilon} \langle |\nabla_{\perp} \delta q|^2 \rangle + P, \quad \langle \delta v_x \delta q \rangle = -\chi \partial_x \langle q \rangle,$$

Build the reduced model to understand the mechanism of staircase based on H-W equation



- Finally: mean density n , vorticity u , + turbulent potential enstrophy ε

$$\partial_t n = \partial_x D_n \partial_x n + D_c \partial_x^2 n, \quad (1)$$

$$\partial_t u = \partial_x (D_n - \chi) \partial_x n + \chi \partial_x^2 u + \mu_c \partial_x^2 u, \quad (2)$$

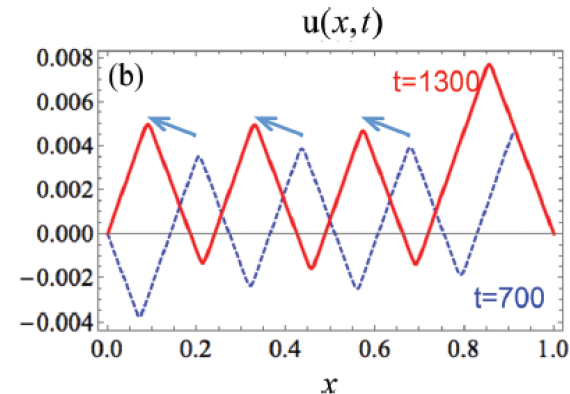
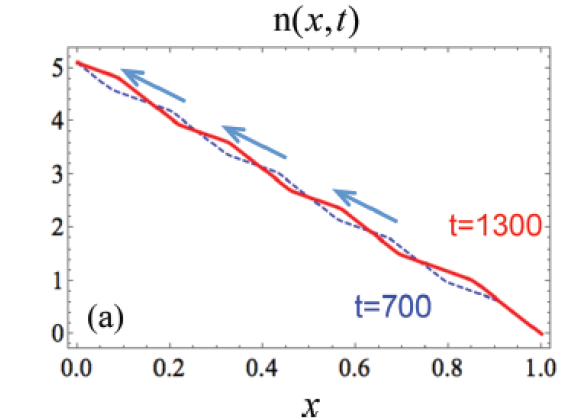
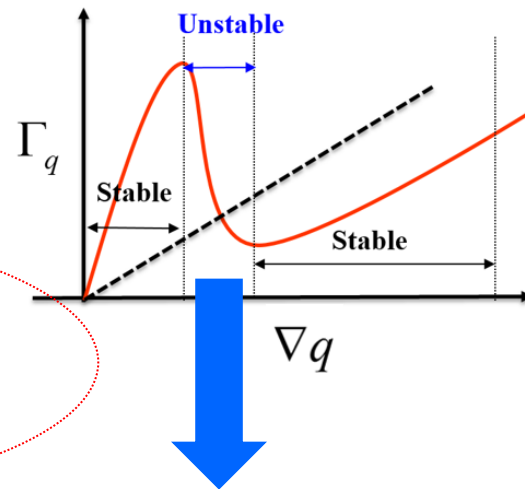
$$\partial_t \varepsilon = \partial_x D_\varepsilon \partial_x \varepsilon + \chi [\partial_x (n - u)]^2 - \varepsilon_c^{-1} \varepsilon^{3/2} + \mathcal{P}. \quad (3)$$

- $D_n \approx l^2 \frac{\varepsilon}{\alpha}$ $\chi(x) = c_\chi l^2 \frac{\varepsilon}{\alpha}$ $D_\varepsilon(x) \cong \beta l^2 \varepsilon^{1/2}$

$$l \rightarrow l_{\text{mix}} = \frac{l_0}{\left(1 + l_0^2 \left[\partial_x (n - u)\right]^2 / \varepsilon\right)^{\kappa/2}}$$

Steepen of mean PV gradient

Further drop of l ← Drop in local ε



- The positive feedback loop in the **negative diffusion region** drives the instabilities which lead to nonlinear feature formation in the mean profile. i.e., formation of **density staircase and vorticity corrugation**

- **Existence of the density staircase and vorticity corrugation is depends on the parameter and B.C.**

What kills the nonlinear features?

→**Scan the parameters.**

→**Change the B.C.**

- **Is there exist other new feedback loop?**

→**Shearing feedback loop**

- **Does the staircase structure can still be sustained or not if the external shear (macro-scale) exist in the plasma?**

→**Modify the model**

→**Run the modified model**

→**Flux driven**

Parameter dependence of staircase !!

Variation	Number of jumps
Parameters	
Suppression parameter $\kappa \neq 2$	Reduce and kill
Increase flow viscosity	Reduce
Increase particle viscosity	Reduce
Increase PE turbulence spreading multiplier β	Reduce
Increase Initial Density gradient	Firstly increase and then reduce
Increase production rate	Increase
Increase initial mean vorticity	Decrease

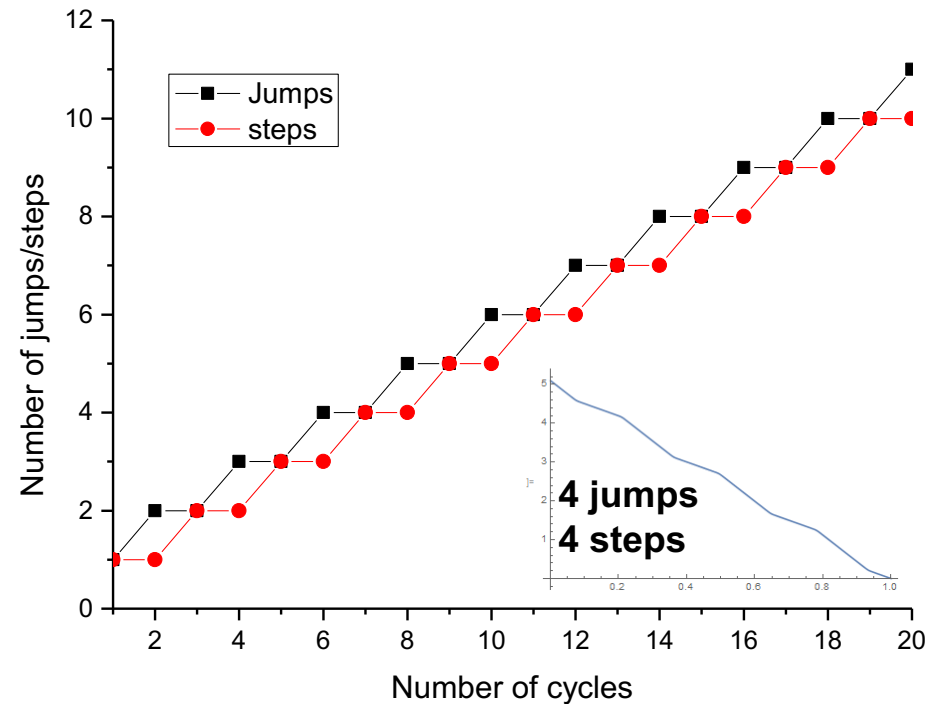
- **Vorticity corrugation disappear \implies Kill the staircase!**
- **Then, if we give a seed corrugation structure by changing the B.C. condition, what will happen ?????**

$$\begin{aligned}
 u(x, t = 0) &= 0; \\
 u(0, t) &= u(1, t) = 0.
 \end{aligned}
 \implies
 u = a \sin(n\pi x) + b$$

a ---background mean zonal shear;
 b ---constant mean shear

(1) Rhines: initial $u = a \sin(n\pi x) + b$
 (a) $a=1$, $b=0.001$, scan n

	Number of jumps N_{jump}	Number of steps N_{step}
$\sin(\pi x) + 0.001$	1	1
$\sin(2\pi x) + 0.001$	2	1
$\sin(3\pi x) + 0.001$	2	2
$\sin(4\pi x) + 0.001$	3	2
$\sin(5\pi x) + 0.001$	3	3
$\sin(6\pi x) + 0.001$	4	3
$\sin(7\pi x) + 0.001$	4	4
$\sin(8\pi x) + 0.001$	5	4
$\sin(9\pi x) + 0.001$	5	5
$\sin(10\pi x) + 0.001$	6	5
$\sin(11\pi x) + 0.001$	6	6
$\sin(12\pi x) + 0.001$	7	6

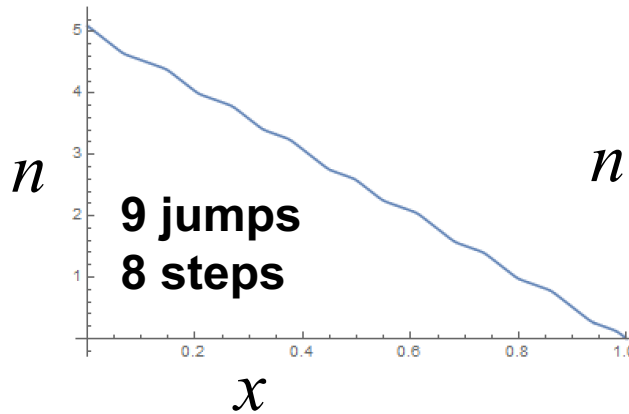


- Specially, when $n=7$, Same to the **wavelength** in A&D paper
- What's more, also three stages!
Microscale instabilities → **NL mesoscale stru.** → **Merger** → **Migration**

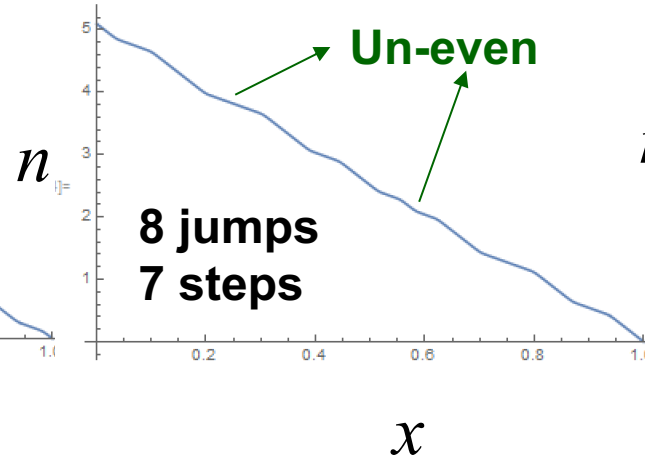
(1) Rhines: initial $u = a \sin(n\pi x) + b$
 (b) $n = 4$, $b = 0.001$, scan a



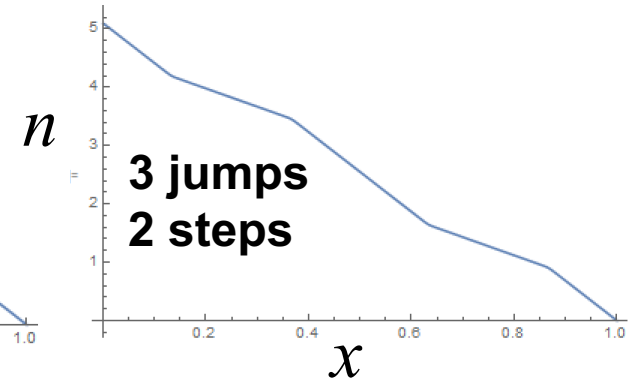
$a = 0.00001$



$a = 0.001$



$a = 0.1$

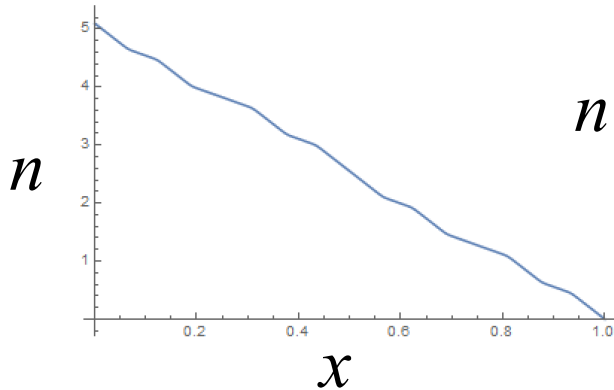


- $a \nearrow \implies$ both N_{jump} and $N_{step} \searrow$
- **Un-even** staircase and be evident when $a \sim b$, which suggests the interplay between mean zonal shear and mean constant shear.
- Stronger zonal background shear a will suppress the drift wave and finally kills the staircase

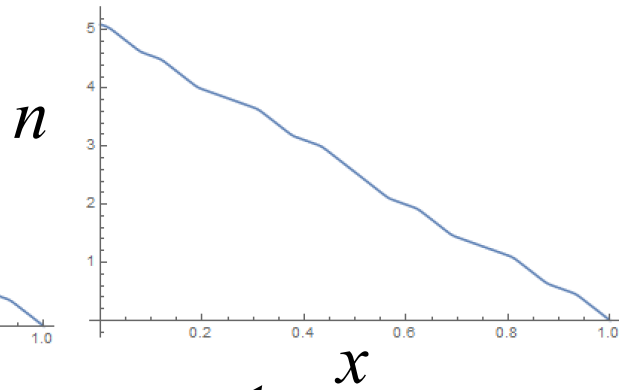
(1) Rhines: initial $u = a \sin(n\pi x) + b$
 (c) $n = 4$, $a = 0.01$, scan b



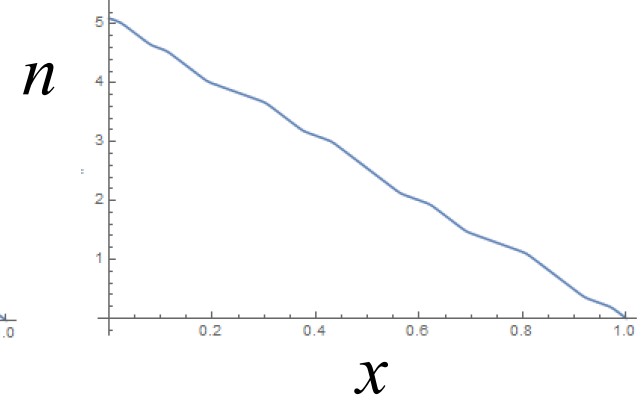
$b = 0.0001$



$b = 0.01$

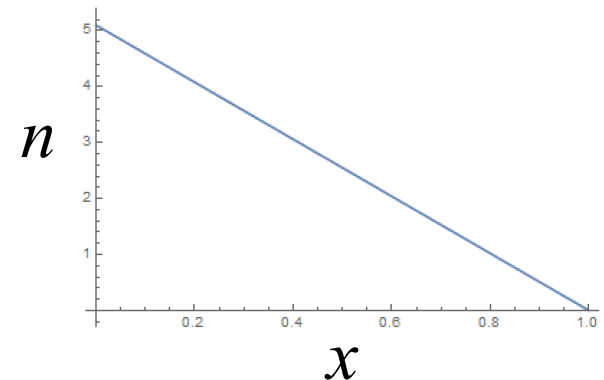


$b = 1$



- Different with a , When $b \leq 1$, b \implies do not have the significant effects on staircase structure. It indicates that the density staircase structure is determined by the zonal mean shear.
- Both stronger mean zonal background shear a and mean constant shear b will finally kill the density staircase by suppressing the drift wave

$b = 3$



(2) Shearing feedback loop

- **In the case of mean poloidal shears, the following form of mixing length is suggested** [H. Biglari, P. H. Diamond, and P. W. Terry, 2(1), 1–4 PoP 1990]

$$l_{mix}^2 = \frac{l_0^2}{\left[1 + (\bar{v}'_y)^2 \tau_c^2\right]^\kappa}$$

- l_0 : the mixing scale without v_y
- τ_c : fluctuation correlation time
- \bar{v}'_y : poloidal shear rate

- **The correlation time is given** [R. J. Hajjar, P. H. Diamond, G. R. Tynan, PoP **25**, 022301 (2018)]

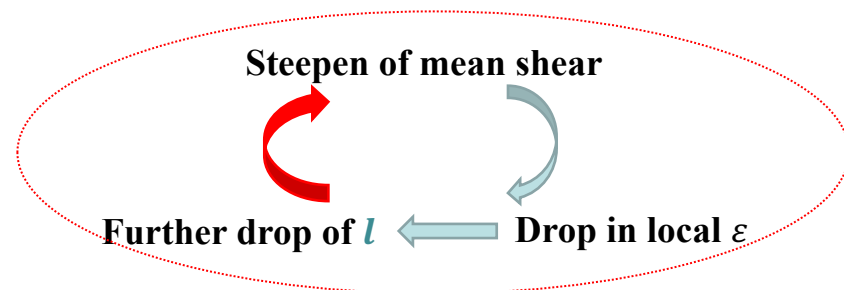
$$\tau_c = \left[\frac{(\bar{v}'_y)^2 \varepsilon}{l_0^2} \right]^{-1/4}$$

- ε : turbulent PE

- **Then, we have the new mixing lengthen and new feedback loop as**

$$l_{mix} = \frac{l_0}{\left[1 + \frac{1}{l_0} \frac{u}{\sqrt{\varepsilon}}\right]^{\kappa/2}}$$

- **Note:** $u = \bar{v}'_y$



(2) Shearing feedback loop

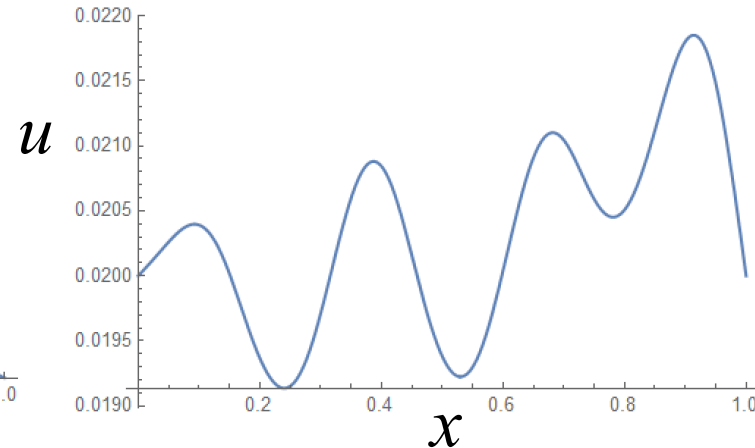
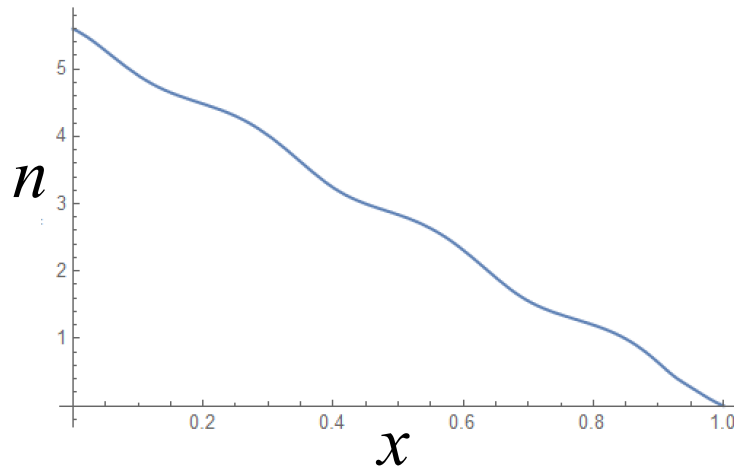


- **While, for the same parameters and same B.C. condition as in *Ashourvan and Diamond, PoP 2017, PRE 2016*, the staircase can not be reproduced !!!**
- **Rhines vs Shearing**
$$\partial_x(n - u) \text{ vs } u \text{ !!!!}$$
- **Only keep the density gradient**
- **Then, change the B.C. to $u = a \sin(n\pi x) + b$ to give a zonal mean shear**

Shearing feedback loop:

$\kappa = 2$, initial $u = 0.02\sin(n\pi x) + 0.01$

- ✓ Lower D_c , μ_c and ε_c vs Rhines, increase the B.C. for turbulent PE,
 $\varepsilon(x, t = 0) = \varepsilon(0, t) = \varepsilon(1, t) = \varepsilon_i = 0.2048$; $n=7$



- ✓ Scan initial density gradient $n(x, t = 0) = -g_i x$, and period n

	Number of jumps formed at $t=1$
Increase $g_i \in [1, 6]$ with $n=4$	<ul style="list-style-type: none"> ✓ Firstly increase and then decrease ✓ Widen the the initial density region for generating staircase by increasing the smallest g_i
Increase $n \in [1, 10]$ with $g_i=5.6$	<ul style="list-style-type: none"> ✓ Firstly increase and then decrease

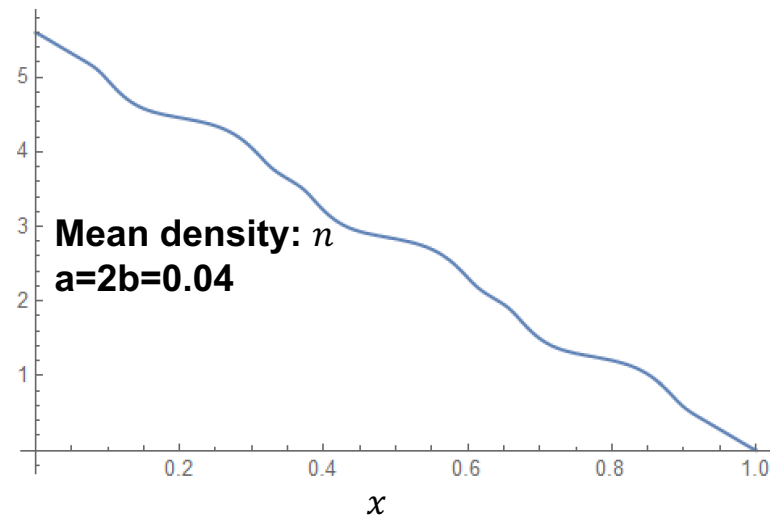
Shearing feedback loop:

$$\kappa = 2, \text{ initial } u = a \sin(7\pi x) + b$$



✓ Scan $u = a \sin(7\pi x) + b$

	Number of jumps N_{jump} formed at $t=1$
Increase $a \in [0.001, 0.01]$ with $b=0.01$	<ul style="list-style-type: none">✓ The increase of a is beneficial for staircase✓ Further increase of a just makes the jump and step characteristic more evident rather than increases N_{jump} too much
Increase $b \in [0, 0.05]$ with $a=0.04$	<ul style="list-style-type: none">✓ Finite $b \approx$ is necessary for the formation of staircase structure✓ The increase of b will increase N_{jump} and can also cause the staircase structure become not very uniform✓ When $b > a$, staircase disappear



- The the emergence and sustainment of staircase structure are sensitive to :
 - ✓ Feedback loop
 - Recovered and sustained the staircase for Rhines scale feedback
 - Sustained staircase **is not observed for shearing feedback**
 - **Density gradient is essential** to staircase formation:
 $\nabla n \uparrow \longrightarrow D(\nabla n) \downarrow$
 - ✓ Drive
 - Steepening ∇n and increasing the production rate **protect** staircase;
 - While, increasing viscosity, PE turbulence spreading β and mean vorticity **weaken** the staircase.
 - ✓ Initial condition
 - Increase initial mean vorticity (mean shear) **weakens** the staircase by suppressing drift wave;
 - **Oscillating** initial condition **persist footprint** on staircase pattern
 - Inhomogeneous staircase due to the interplay of zonal shear and mean shear

Future works---Looking forward for another possible opportunity



- Understand the feedback loops
→ Staircase emergence via Rhines scale with ∇n only??
- Why shearing feedback is not working?
- How to sustain the staircase?
→ Flux driven studies
- Understand these sensitivities more carefully and deeply!
- Why a wider zonal modulation pattern is not observed?
- Staircase vs ITB
- Finish the modification of the reduced model to study the external shear
.....

EVERYBODY
IS A
GENIUS.
BUT IF YOU
JUDGE A
FISH BY ITS
ABILITY
TO CLIMB
A TREE,
IT WILL LIVE
ITS WHOLE
LIFE
BELIEVING
THAT IT IS
STUPID.

— ALBERT EINSTEIN



**Thank you very much for
your attention!**

**Thanks for all the sweet
caring!**