

Rossby Wave-Zonal Flow Turbulence in a Tangled Magnetic field

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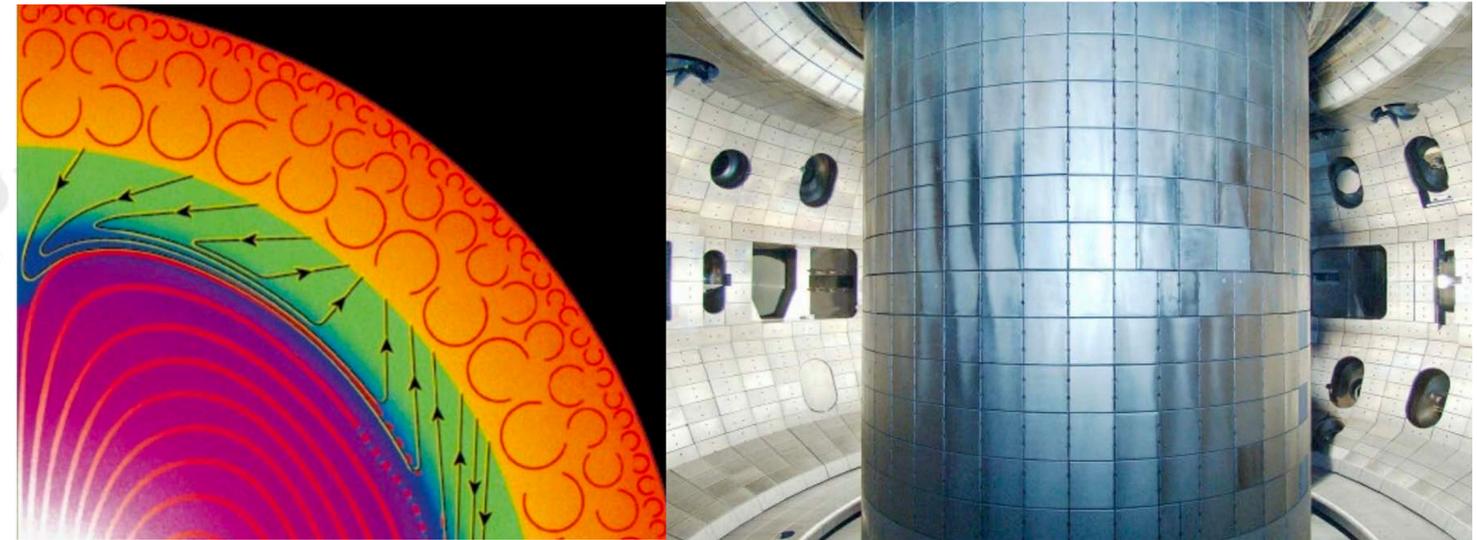
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Introduction

Why we study tangled magnetic fields?

1. Astro: Interstellar Medium and Solar tachocline.
2. Fusion: energy and momentum transport problem in L-H transition at edge of tokamaks with resonant magnetic perturbation (RMP).

Turbulent transport processes still poorly understood.



We focus on the solar tachocline:

1. Between the convective and radiative zone.
2. Strongly stratified/Pancake-like structures.

Incompressible rotating fluid in 2D layers— β -plane model

3. Zonal Flow and Rossby Waves— as in the Jovian Atmosphere.

4. A weak mean field— large magnetic Kubo number:

$$\mathbf{B} = B_0 + \widetilde{\mathbf{B}}$$

$$Ku_{mag} \equiv \frac{\delta_l}{\Delta_{eddy}} = \frac{l_{ac} |\widetilde{\mathbf{B}}|}{\Delta_{eddy} B_0}$$

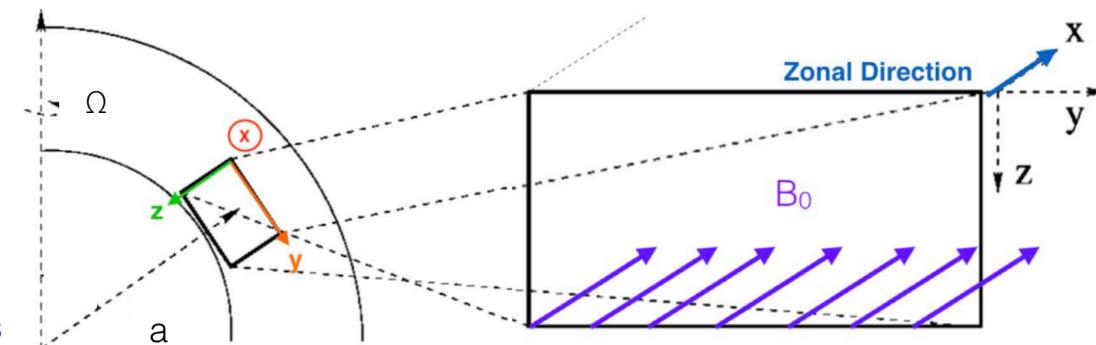
Quasi-Linear Theory (QLT) fails.

Need a model beyond QLT.

$$\beta = \frac{df}{dy} \Big|_{\phi_0} = \frac{2\Omega \cos(\phi_0)}{a}$$

↑ rotation ↑ radius
↑ latitude

Derivative of angular frequency f (Coriolis parameter)



Outline

- * **A model for PV transport in Strong mean magnetic fields**
- * A Model for PV Transport in Random, small-scale Magnetic fields

Model— Strong mean field & The Quasi-Linear Method

Notations we have:

Stream Function $\psi = \psi(x, y, z)$
 Velocity field $\mathbf{u} = \left(\frac{\partial\psi}{\partial y}, -\frac{\partial\psi}{\partial x}, 0\right)$
 Fluid Vorticity $\zeta = (0, 0, \zeta)$
 Potential Field $\mathbf{A} = (0, 0, A)$
 Magnetic Field $\mathbf{B} = \left(\frac{\partial A}{\partial y}, -\frac{\partial A}{\partial x}, 0\right)$

Two main equations:

➤ QL closure

➤ Linear response of perturbations

$$\left(\frac{\partial}{\partial t} + \mathbf{u}_\perp \cdot \nabla_\perp\right)\zeta - \beta \frac{\partial\psi}{\partial x} = -\frac{(\mathbf{B}\nabla)(\nabla^2 A_z)}{\mu_0\rho} + \nu \nabla^2 \zeta$$

$$\left(\frac{\partial}{\partial t} + \mathbf{u}_\perp \cdot \nabla_\perp\right)A = B_0 \frac{\partial\psi}{\partial x} + \eta \nabla^2 A,$$

Quasi-Linear Approximation:

$$\zeta = \langle \zeta \rangle + \tilde{\zeta}$$

$$\psi = \langle \psi \rangle + \tilde{\psi}$$

$$A = \langle A \rangle + \tilde{A}$$

Perturbations produced by turbulences

, where $\langle \rangle = \frac{1}{L} \int dx \frac{1}{T} \int dt$

ensemble average over the zonal scales

PV flux:

➤ By using Taylor Identities we have the Reynolds and Maxwell stress.

➤ Express the PV flux with two diffusivities

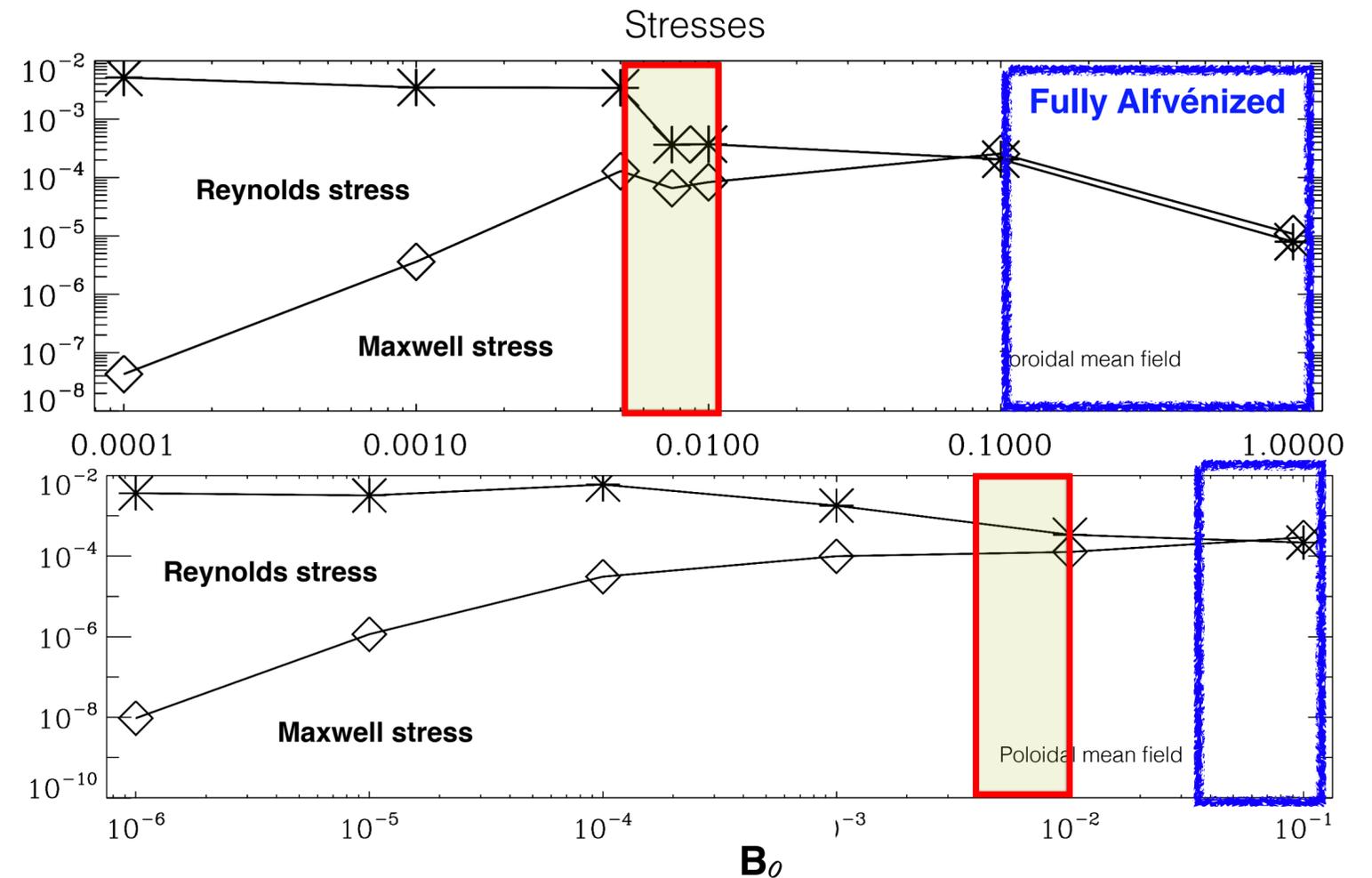
$$D_{fluid} = \sum_k |\tilde{u}_{y,k}|^2 \frac{\nu k^2 + \omega_A^2 \frac{\eta k^2}{\omega^2 + \eta^2 k^4}}{\left(\omega - \omega_A^2 \frac{\omega}{\omega^2 + \eta^2 k^4}\right)^2 + \left(\nu k^2 + \omega_A^2 \frac{\eta k^2}{\omega^2 + \eta^2 k^4}\right)^2}$$

$$D_{mag} = \sum_k |\tilde{u}_{y,k}|^2 \frac{\omega_A^2 \left(\nu k^2 (\omega^2 + \eta^2 k^4) + \omega_A^2 \eta k^2\right)}{\omega^2 \left(\omega^2 + \eta^2 k^4 - \omega_A^2\right)^2 + \left(\nu k^2 (\omega^2 + \eta^2 k^4) + \omega_A^2 \eta k^2\right)^2}$$

The Reynolds stress drops occurs at weaker B_0 than that for which the system is fully Alfvénized!



The cross phase effect suppresses the Reynolds stress when mean field is weak!

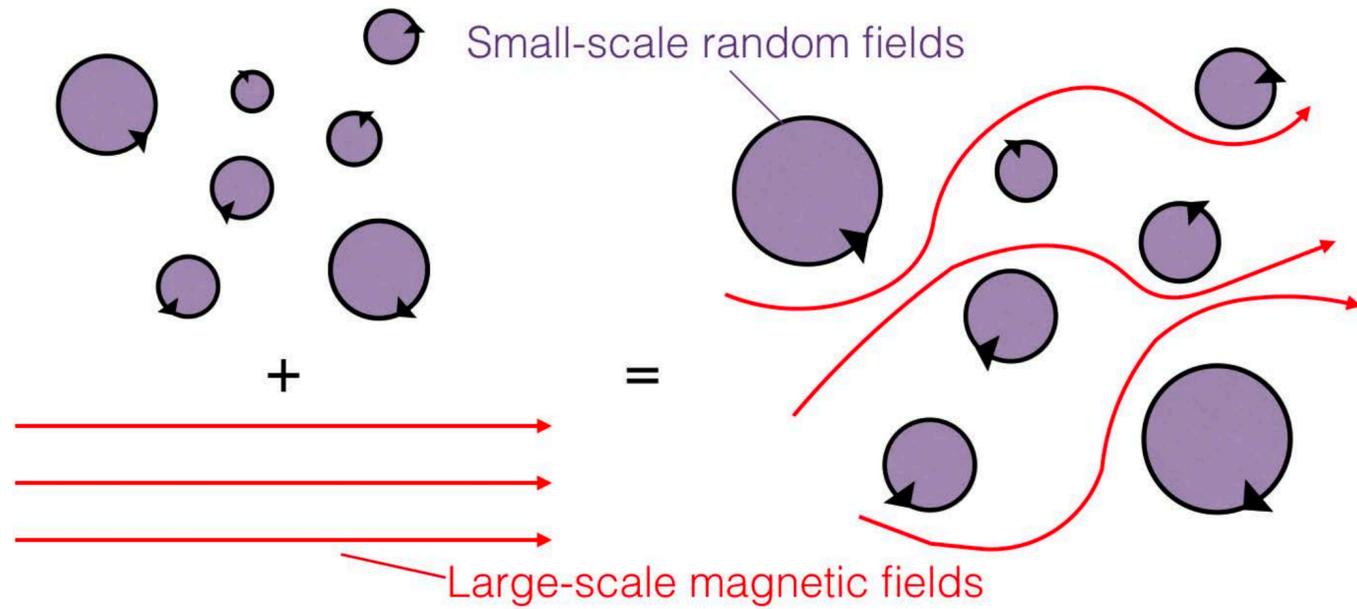


(Tobias et al. in preparation)

Outline

- * A model for PV transport in Strong mean magnetic fields
- * **A Model for PV Transport in Random, small-scale Magnetic fields**

Model— Random fields: Order of Scales



The large-scale magnetic field is distorted by the small-scale fields. The system thus is the 'soup' of cells threaded by sinews of open field line (Zeldovich, 1957).

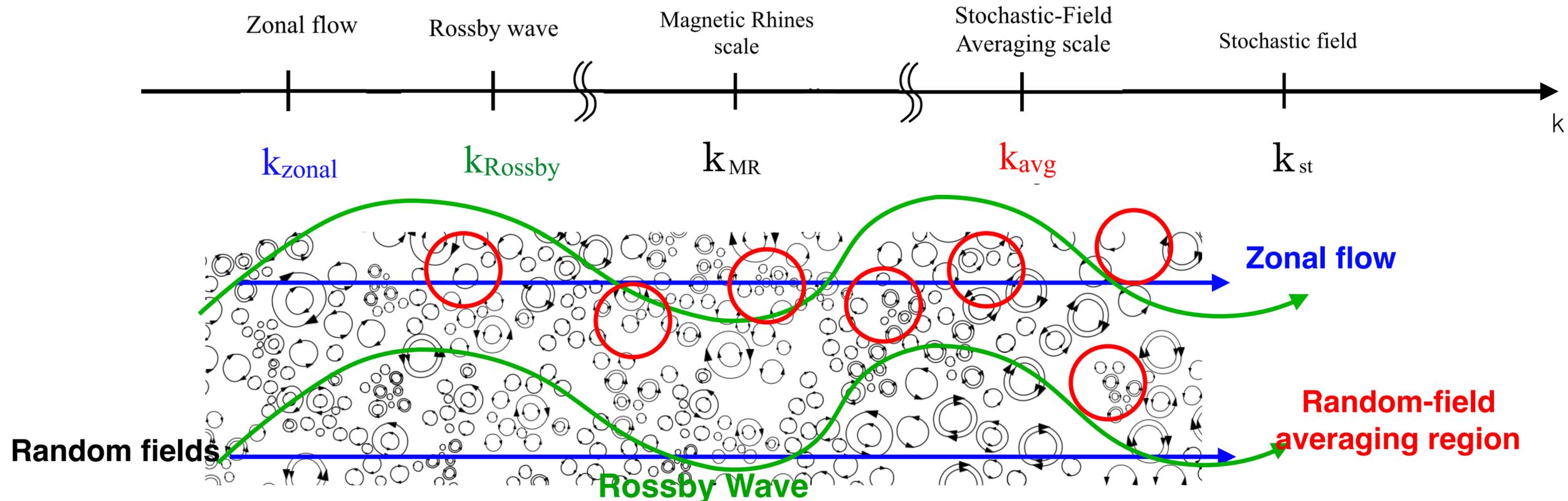
● **Properties of random fields are (Rechester & Rosenbluth 1978):**

1. Smaller scale, Static
2. Randomness in space
3. Auto-correlation length of fields is small. $\longrightarrow Ku_{mag} \equiv \frac{\tilde{u}\tau_{ac}}{\Delta_{eddy}} = \frac{l_{\parallel} |\tilde{\mathbf{B}}|}{\Delta_{\perp} B_l} < 1$
4. Amplitudes of random fields distributed statistically. (assumption: PDF Gaussian)

● **Two-average Method:**

1. $\bar{F} = \int dR^2 \int dB_{st} \cdot P_{(B_{st,x}, B_{st,y})} F$
2. $\langle\langle \rangle\rangle = \frac{1}{L} \int dx \frac{1}{T} \int dt$ ensemble average over the zonal scales

● **Order of Scales:** Function of fields $\mathbf{F} = \mathbf{F}_0 + \tilde{\mathbf{F}} + \mathbf{F}_{st}$



Model— Random fields: Assumption, Derivation & Results

● Assumptions:

1. The collective field at Rossby-scale is NOT large enough to alter the structure of the random fields: $\overline{B_{st}} \rightarrow 0$

2. We approximate the correlation matrix as diagonal: $\overline{B_{r,x}B_{r,y}} = 0$

● Derivation:

Two main equations:
$$\begin{cases} \frac{\partial}{\partial t} \bar{\zeta} - \beta \frac{\partial \bar{\psi}}{\partial x} = \frac{\overline{(\mathbf{B} \cdot \nabla) J}}{\mu_0 \rho} + \nu \nabla^2 \bar{\zeta} \\ \frac{\partial}{\partial t} A = \mathbf{B} \cdot \nabla \psi + \eta \nabla^2 A. \end{cases}$$

Key term:
Average effect of $\mathbf{J} \times \mathbf{B}$

Linear response of the vorticity:
$$\tilde{\zeta}_k = \left(\frac{i}{\omega + i\nu k^2 + \frac{i\overline{B_{st,j}^2} k_j^2}{\mu_0 \rho \eta k^2} + \frac{i}{\mu_0 \rho} \frac{B_0^2 k_x^2}{\eta k^2 - i\omega}} \right) \tilde{u}_{y,k} \left(-\frac{\partial}{\partial y} \bar{\zeta} - \beta \right)$$

● Dispersion relation of the Rossby-Alfvén wave in random magnetic fields:

$$\left(\omega - \omega_R + \frac{i\overline{B_{st,y}^2} k_y^2}{\mu_0 \rho \eta k^2} + i\nu k^2 \right) \left(\omega + i\eta k^2 \right) = \frac{B_{0,x}^2 k_x^2}{\mu_0 \rho}$$

↑
↑

Dissipative response to Random magnetic fields
AW of the large-scale magnetic field

(mean square)
(square mean)

Random fields: Result I— Cross-Phase Modification

● **Mean PV Flux ($\bar{\Gamma}$) and PV diffusivity (D_{PV}):**

$$\bar{\Gamma} = -\frac{\partial}{\partial y} \langle \widetilde{u}_x \widetilde{u}_y \rangle \quad (\text{Reynolds stress force})$$

$$\bar{\Gamma} = - \sum_k |\widetilde{u}_{y,k}|^2 \frac{\nu k^2 + \left(\frac{B_0^2 k_x^2}{\mu_0 \rho}\right) \frac{\eta k^2}{\omega^2 + \eta^2 k^4} + \frac{\overline{B_{st,y}^2} k_y^2}{\mu_0 \rho \eta k^2}}{\left(\omega - \left(\frac{B_0^2 k_x^2}{\mu_0 \rho}\right) \frac{\omega}{\omega^2 + \eta^2 k^4}\right)^2 + \left(\nu k^2 + \left(\frac{B_0^2 k_x^2}{\mu_0 \rho}\right) \frac{\eta k^2}{\omega^2 + \eta^2 k^4} + \frac{\overline{B_{st,y}^2} k_y^2}{\mu_0 \rho \eta k^2}\right)^2} \left(\frac{\partial}{\partial y} \bar{\zeta} + \beta\right)$$

$$\bar{\Gamma} = -D_{PV} \left(\frac{\partial \bar{\zeta}}{\partial y} + \beta\right)$$

PV Diffusivity D_{pv}

Large-scale field small-scale random fields

The large- and small-scale magnetic fields have a synergistic effect on the cross-phase in the Reynolds stress.

In a certain limit where $B_0^2 \ll \overline{B_{st,y}^2}$:

$$D_{PV} = \sum_k |\widetilde{u}_{y,k}|^2 \frac{\nu k^2 + \frac{\overline{B_{st,y}^2} k_y^2}{\mu_0 \rho \eta k^2}}{\omega^2 + \left(\nu k^2 + \frac{\overline{B_{st,y}^2} k_y^2}{\mu_0 \rho \eta k^2}\right)^2}$$

Cross-phase effect in the Reynolds stress **decreases for stronger random fields.**

Random fields: Result II— Resisto-Elastic Medium

● Evolution of Zonal Flow

➤ Random magnetic fields suppress the Reynolds stress and increase the drag.

$$\frac{\partial}{\partial t} \langle u_x \rangle = \langle \bar{\Gamma} \rangle - \frac{1}{\eta \mu_0 \rho} \langle \overline{B_{st,y}^2} \rangle \langle u_x \rangle + \nu \nabla^2 \langle u_x \rangle$$

Cross-phase effect on Reynolds Stress Force

Magnetic drag force ($J_{st} \times B_{st}$)

Random magnetic fields have an effect on both the PV flux and the magnetic drag.

Two effects!

● Resisto-Elastic medium:

If we turnoff the Rossby frequency, we have a 2D non-rotating plane:

$$\omega^2 + i(\alpha + \eta k^2)\omega - (\alpha \eta k^2 + \chi) = 0$$

↑ drag + resistivity dissipation
 ↑ effective spring constants

$$\alpha \equiv \frac{\overline{B_{st,j}^2} k_j^2}{\mu_0 \rho \eta k^2} \propto \frac{\text{spring constant}}{\text{dissipation}}$$

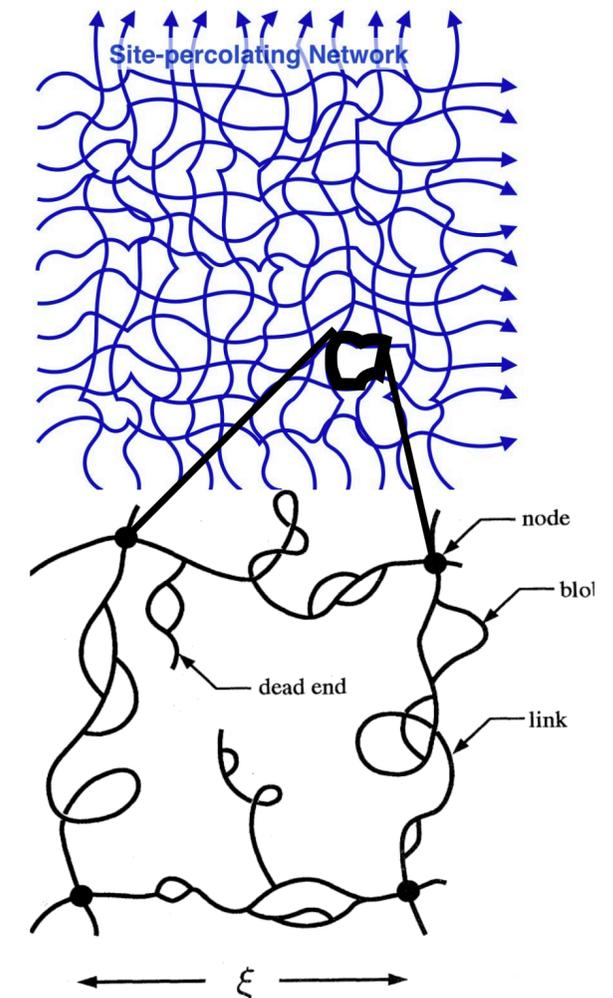
$$\chi \equiv \frac{B_0^2 k_x^2}{\mu_0 \rho}$$

➔ **Alfvén waves propagate in a resisto-elastic medium.**
Energy dissipates due to the coupling of the drag and the resistivity.

● More can be done:

Fractal Network (Site-percolating) —

Calculate the effective spring constant, effective Young's Modulus of elasticity, and effective “conductivity” of vorticity (such as encountered in amorphous solids).



Schematic of the nodes-links-blobs model (Nakayama & Yakubo 1994).

Takeaways I

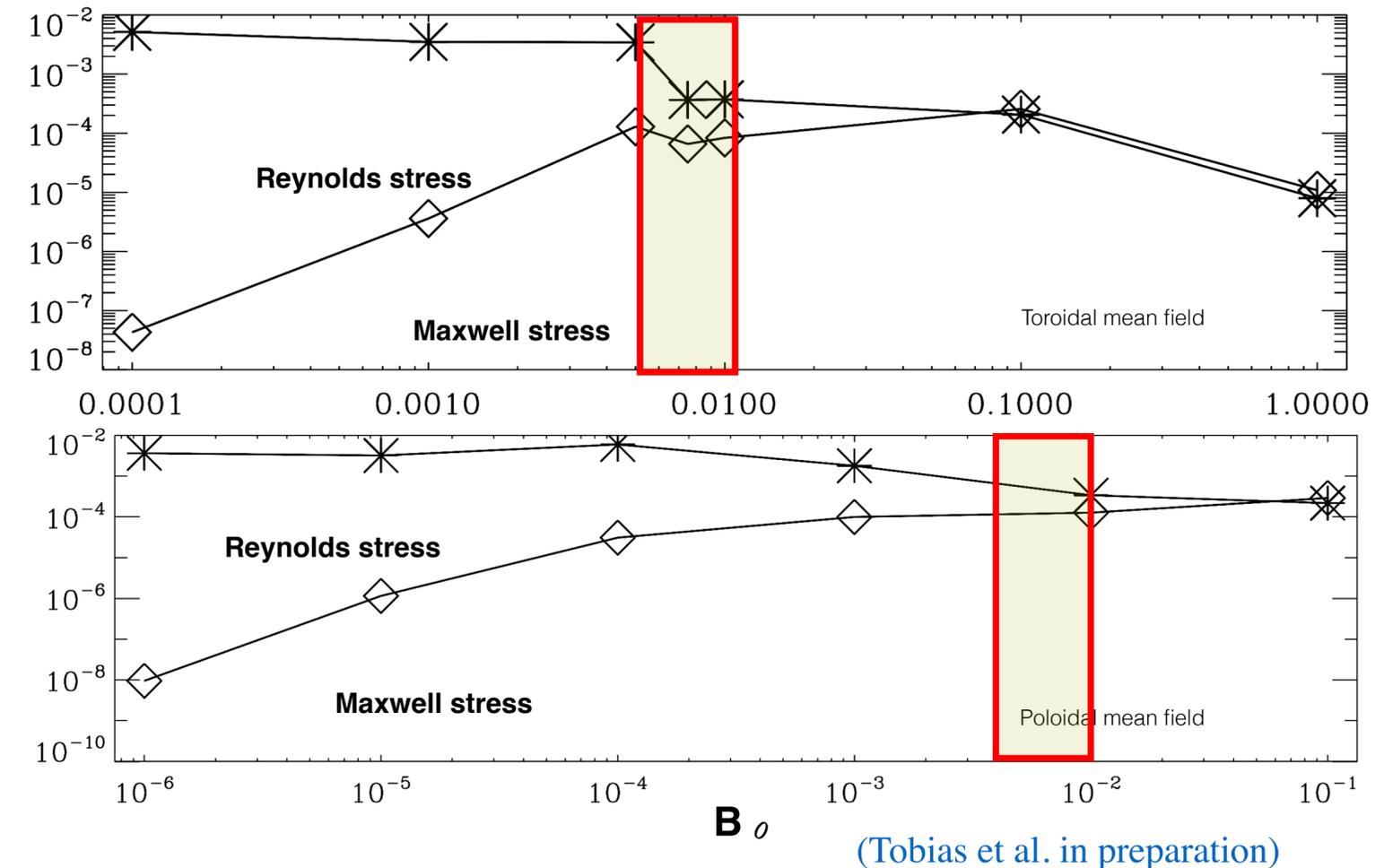
- What studies have shown and what we obtained:

Reynolds stress will be suppressed at levels of field intensities **well below that of Alfvénization**, where Maxwell stress balances the Reynolds stress.

The flow generated by PV mixing/Reynolds force are reduced by:

- ◆ Coupling of resisto-elastic waves, which is $\overline{B_{st}^2}$ dependence.
- ◆ Increase of the magnetic drag.

Cross-phase effect and the magnetic drag reduce shear flow generation.



● Related: Experiment on edge of DIII-D:

1. Stochastic magnetic fields due to RMP inhibits Reynolds stress and shear suppression.

This work is relevant— close analogy between Rossby-Alfvén wave and EM drift wave turbulence.

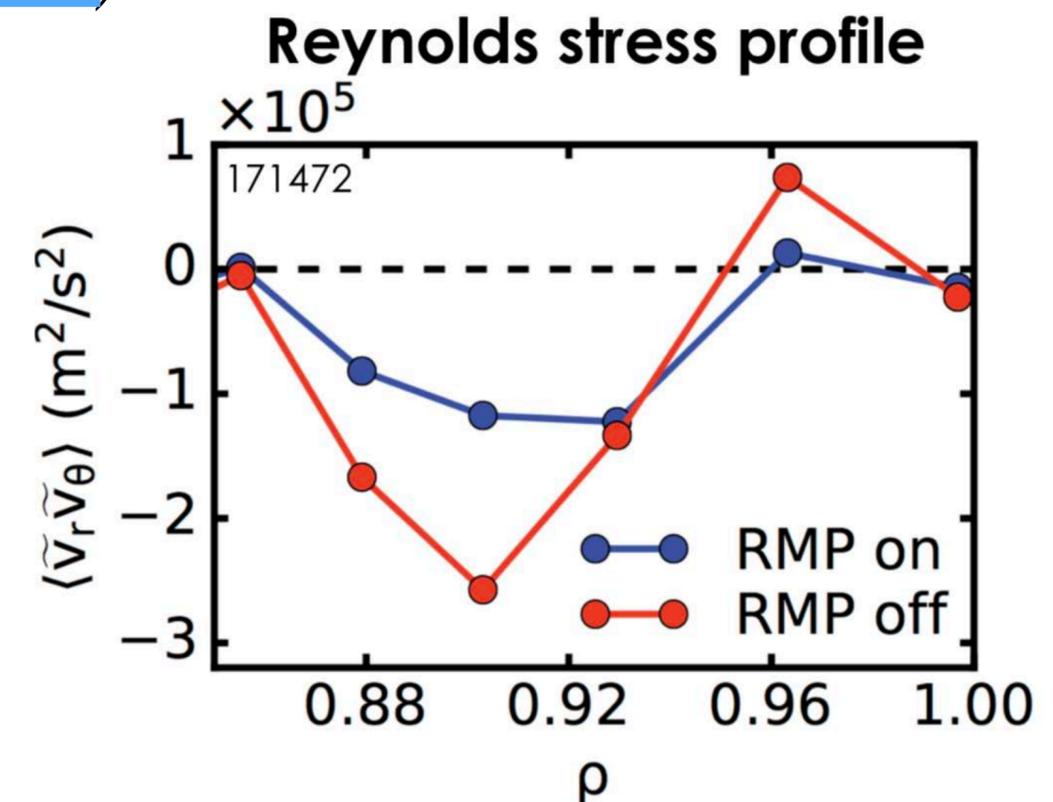
$$\bar{\Gamma} = - \sum_k |\tilde{u}_{y,k}|^2 \frac{\nu k^2 + \left(\frac{B_0^2 k_x^2}{\mu_0 \rho}\right) \frac{\eta k^2}{\omega^2 + \eta^2 k^4} + \frac{\overline{B_{st,y}^2} k_y^2}{\mu_0 \rho \eta k^2}}{\left(\omega - \left(\frac{B_0^2 k_x^2}{\mu_0 \rho}\right) \frac{\omega}{\omega^2 + \eta^2 k^4}\right)^2 + \left(\nu k^2 + \left(\frac{B_0^2 k_x^2}{\mu_0 \rho}\right) \frac{\eta k^2}{\omega^2 + \eta^2 k^4} + \frac{\overline{B_{st,y}^2} k_y^2}{\mu_0 \rho \eta k^2}\right)^2} \left(\frac{\partial}{\partial y} \bar{\zeta} + \beta\right)$$

2. Our results imply: Reduction of Reynolds stress, due to the effect of the stochastic field $\overline{B_{st}^2}$ in cross-phase, inhibits the growth of zonal flow, by coupling energy into elastic waves.

Hence, the **nonlinear energy transfer** from turbulence to flow and thus **shear suppression** at the edge are reduced.

(See experiments from D. M. Kriete et al.)

3. Conjecture: Stochasticity of fields might raise the power threshold, by weakening Reynolds stress trigger.



(D. M. Kriete et al., TTF 2019)

Thank you!