

Entrainment in Drift Wave Turbulence - A Basic Study

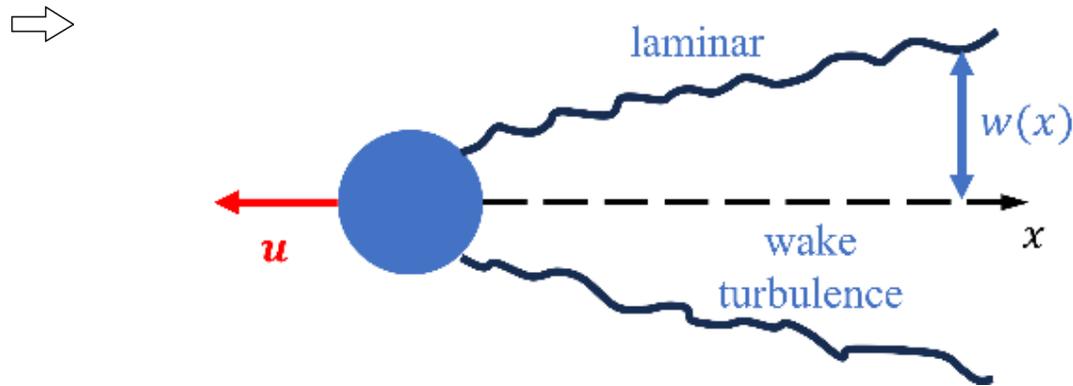
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Wake-Classic Example of Turbulence Spreading



Similarity Theory }
 Mixing Length Theory }
 $W \sim (F_d / \rho U^2)^{1/3} X^{1/3},$
 $F_d \sim C_D \rho U^2 A_s$
 C_D independent of viscosity at high Re

⇒ Physics: Entrainment of laminar region by expanding turbulent region.
 Key is turbulent mixing. ⇒ Wake expands

⇒ Townsend '49:

- Distinction between momentum transport — eddy viscosity—and fluctuation energy transport
- Failure of eddy viscosity to parametrize spreading

— Jet Velocity: $V = \frac{\langle V_{perp} * V^2 \rangle}{\langle V^2 \rangle}$ ⇒ spreading flux FOM See Ting Long for measurement

Spreading in MFE Theory

⇒ Numerous gyrokinetic simulations

N.B. Basic studies absent ...

i. e.

$$\partial_t \xi = \gamma \xi (1 - \xi) + \partial_x D(\xi) \partial_x \xi + D_0 \partial_x^2 \xi$$

⇒ Diagnosis primarily by: - color VG

$$\gamma \sim 0(\varepsilon)$$

 - tracking of “Front”

⇒ Theory ⇒ Nonlinear Intensity diffusion models

 ⇒ Reaction-Diffusion Equations - especially Fisher + NL diffusion

 ⇒ Continuum DP Models - Later.....

Recently:

⇒ Renewed interest in context of λ_q broadening problem, cf. P.D., Z. Li, Xu Chu

⇒ Simulations measure correlation of spreading $\langle \tilde{V}_r \tilde{p} \tilde{p} \rangle$ with λ_q broadening, cf. N. Li +

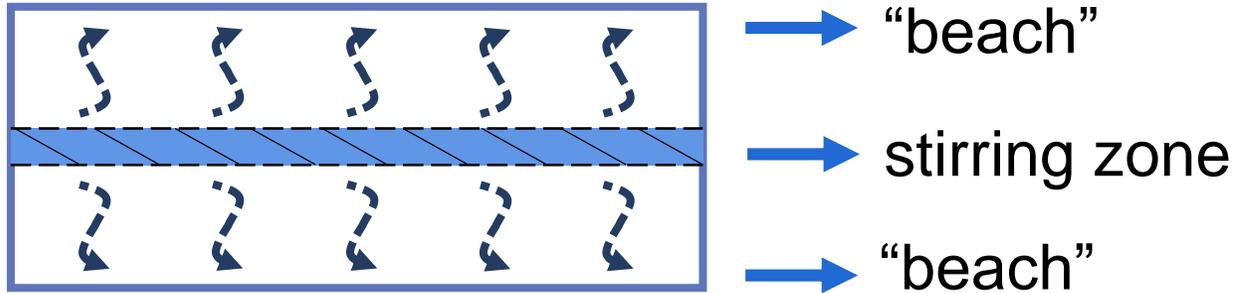
⇒ Intermittency effects T. Wu, P. D. + 2023, A. Sladkomedova 2024



Especially blobs, voids

Spreading Studies - Numerical Experiments

⇒ 2D Box, Localized Stirring Zone



⇒ Comparison of:

<u>System</u>	<u>Features</u>
2D Fluid	Selective Decay, Vortices How to Measure Spreading?
2D MHD with weak B_0 perp.	Alfvenization, Vortex Bursting, Zeldovich number
Forced Hasegawa-Mima with Zonal Flow	Waves + Eddies + ZF Multiple regimes and Mechanisms



N.B. Clear distinction between “spreading” and “avalanching”

Numerics: 2D Dedalus simulation

Box Characteristics:

- Dedalus Framework
analogous to BOUT++

- Grid Size: 512×512
- Doubly Periodic boundary condition, beach regulates expansion

Forcing Characteristics:

- Superposition of Sinusoidal Forcing, vorticity
- Spectrum: Constant $E(k)$, ensuring uniform energy distribution across wave numbers.
- Correlation Length: Approximately $1/10$ of the box scale, some room for dual cascade.
- Localized through a Heaviside step function.
- Phase of forcing randomized every typical eddy turnover time

Summary: 2D Fluid + 2D MHD Studies (R.X., P.D. submitted '24)

- Spreading mediated by dipole vortices
 - 2 components
 - Free flyer vortices
 - Turbulent gas/patch
- $W_{FF} \sim t$ (ballistic) – expansion of turbulent layer by dipoles
- FOM – Enstrophy distribution
- No clear ‘front’ - fractalization
- Weak transverse B-field can disrupt vortices, terminate ballistic spreading
- Zeldovich # is good FOM $Z = R_m V_{A,0}^2 / V_{eddy}^2$, $Z > 1$ for disruption
- Disrupted vortices → Alfvén waves

Forced Hasegawa – Mima + Zonal Flows

H-M + Zonal Flow System

— System:

$$\frac{d}{dt} (\tilde{\phi} - \rho_s^2 \nabla_{\perp}^2 \tilde{\phi}) + v_* \frac{\partial \tilde{\phi}}{\partial y} + v_{*u} \frac{\partial \tilde{\phi}}{\partial y} = \frac{\partial}{\partial r} \rho_s^2 \langle \tilde{v}_r \nabla_{\perp}^2 \tilde{\phi} \rangle + \nu \nabla^2 \nabla^2 (\tilde{\phi}) + \tilde{F} \text{ -Waves, Eddys}$$

PV forced
↓

$$\frac{d}{dt} = \frac{\partial}{\partial t} + \bar{v}_z \frac{\partial}{\partial y} - \nabla \tilde{\phi} \times \hat{\mathbf{z}} \cdot \nabla$$

$$\frac{\partial}{\partial t} \nabla_x^2 \bar{\phi}_z + \frac{\partial}{\partial r} \langle \tilde{v}_r \nabla_{\perp}^2 \tilde{\phi} \rangle + \mu \nabla_x^2 \bar{\phi}_z = 0 \text{ -Zonal Flow (Axisymmetric)}$$

N.B. $\bar{\phi}_z = \bar{\phi}_z(x)$, only.

N.B. : Electrons Boltzmann for waves, not for Zonal Flow

— viscosity controls small scales

— drag controls zonal flow - μ (large scale)

— conserved: Energy $\rightarrow \langle \tilde{\phi}^2 + \rho_s^2 (\nabla \tilde{\phi})^2 \rangle + \langle \rho_s^2 (\nabla \phi_z)^2 \rangle$
 Potential Enstrophy $\rightarrow \langle (\tilde{\phi} - \rho_s^2 \nabla_{\perp}^2 \tilde{\phi})^2 \rangle + \langle (\rho_s^2 \nabla^2 \phi_z)^2 \rangle$

\downarrow Waves \downarrow ZF

N.B. Energy, Pot Enstr. exchange between Waves and ZF possible.

Channels, cont'd:

⇒ **Spreading in presence of fixed, externally prescribed shear layer**

⇒ **Here:** → Forcing → $\left\{ \begin{array}{l} \text{Waves} \\ \text{Eddies} \end{array} \right\}$ → Zonal flow (self-generated)

∴ forcing (\tilde{v}_{rms}, Re) + drag ⇒ control parameters

⇒ “weak” and “strong” Turbulence Regimes

$$v_{gr} \text{ VS } v_r \rightarrow \frac{\langle \tilde{v}_r \xi \rangle}{\sum_{\mathbf{k}} v_{gr}(\mathbf{k}) \xi_{\mathbf{k}}} \rightarrow \frac{\tilde{v}_r \tau_c f}{\Delta_c} \rightarrow Ku$$

coherency factor

⇔ 2nd vs 3rd order energy flux

$\Delta_c \sim v_{gr} \tau_c$

⇒ $Ku < 1$ → wave dominated spreading

$Ku > 1$ → mixing dominated spreading ⇒ ~ 2D fluid

→ Dipoles gone – density gradient

Channels, cont'd

But → Enter the ZONAL FLOW

- Multiple channels for NL interaction
- But with ZF ↔ eddy, wave coupling to ZF dominant
- ZF is the mode of minimal inertia, damping, transport

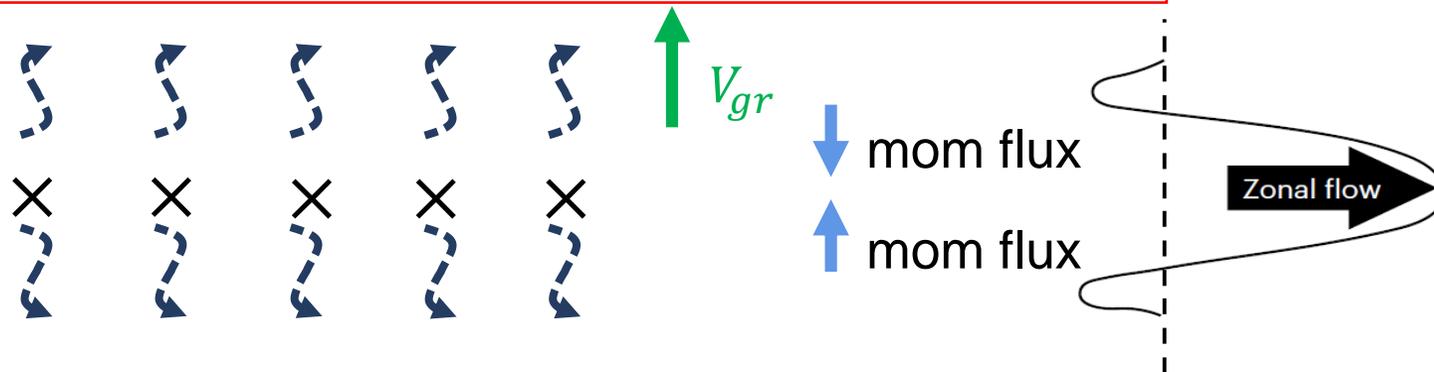
⇒ energy coupled to ZF ($\tilde{v}_r = 0$) cannot “spread”, unless recoupled to waves

Waves:

$$\frac{\partial}{\partial t} (1 + k_{\perp}^2 \rho_s^2) \tilde{\phi} = \dots$$

ZF:

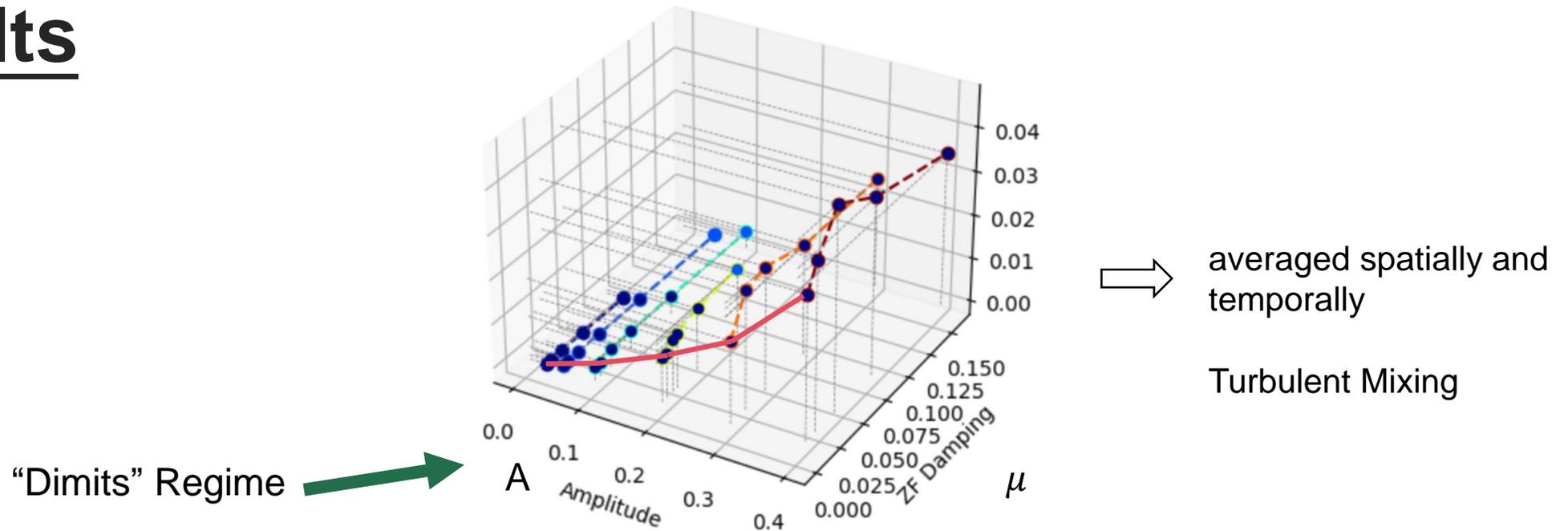
$$\frac{\partial}{\partial t} (k_r^2 \rho_s^2) \bar{\phi}_z = \dots$$



- Degradation of ZF (back transfer) is crucial to spreading
- ∴ μ must regulate spreading. What of $\mu \rightarrow 0$ regimes?
- Revisit collisionless NL dissipation problem

Results

FOM – Fluctuation Potential Enstrophy Flux



- Potential enstrophy flux generally increases as drag increases. “Dimits regime” for turbulence spreading. Spreading diminishes with power coupled to Z.F. (Fixed, spatially)
- Z.F. is self-generated barrier to spreading
- For A increasing, PE flux rises sharply for weak ZF damping. Fate of ZF?
“KH-type” mechanism loss of Dimits regime at higher A? Characterization??

N.B. “Dimits Regime”= Condensation of energy into ZF for weaker forcing.

Results, Cont'd

Wave Energy Flux

- Dimits regime at low forcing and ZF damping
- Increases with ZF damping and forcing amplitude
- Dominant K_x increases due ZF decorrelation
- Spectrum condensation towards low k with inverse cascade



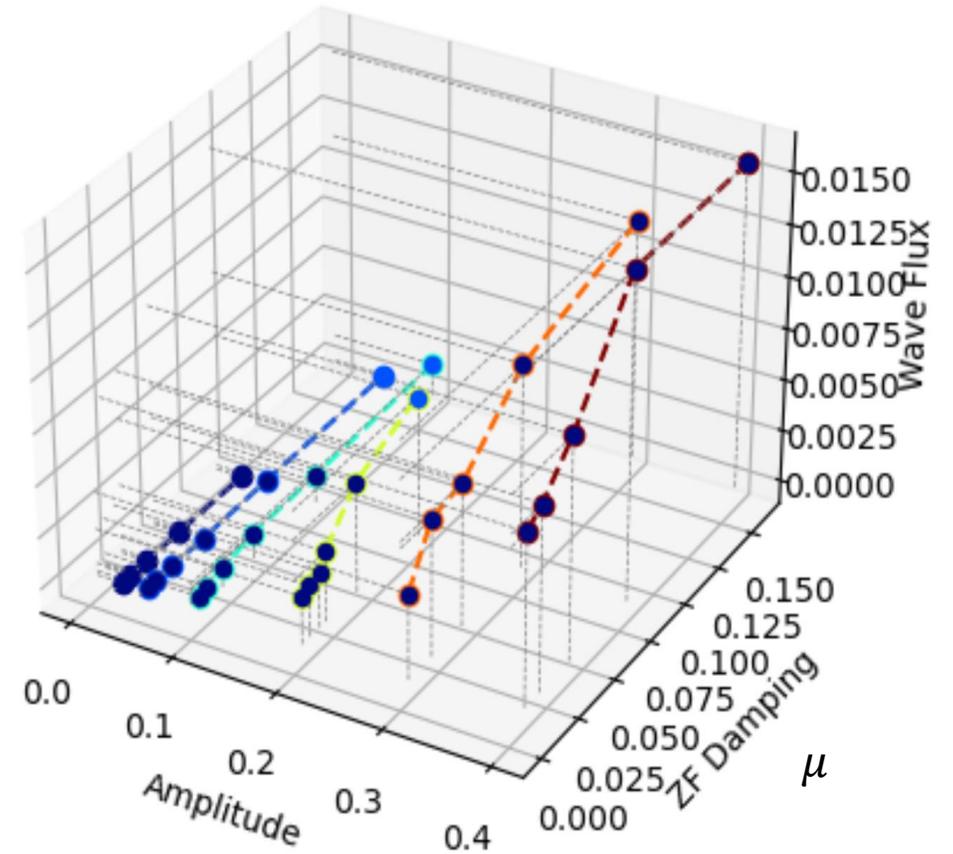
implication for v_{gr} and $\sum_{\mathbf{k}} v_{gr}(\mathbf{k})E_{\mathbf{k}}$

- Take note of increasing W.E.flux as $\mu \rightarrow 0$,
A increases.

Physics: ZF shears refract waves

$$\text{Wave Energy Flux} < -\frac{\partial \phi}{\partial t} \nabla \phi > \longleftrightarrow \sum_{\mathbf{k}} v_{gr}(\mathbf{k})E_{\mathbf{k}}$$

for drift waves

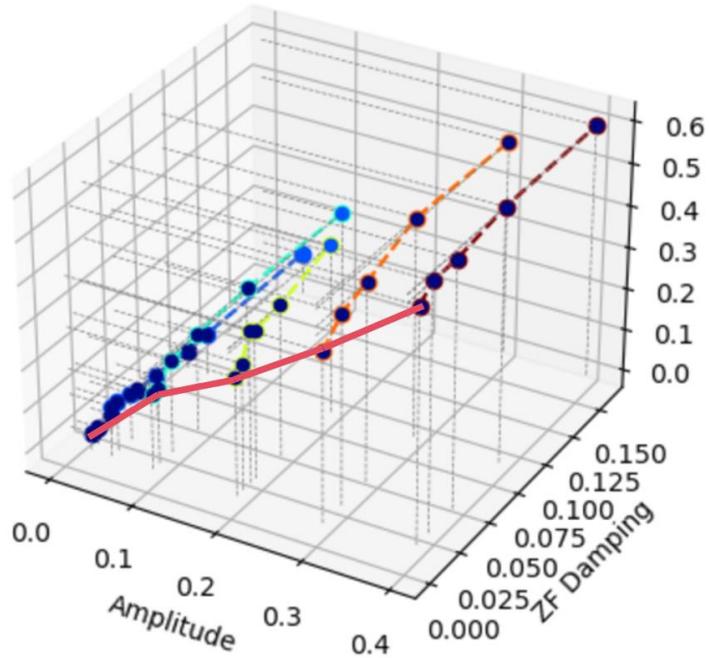


Results, Cont'd

$$\frac{\tilde{v}_r \tau_c f}{\Delta_{cc}} \text{ where } \Delta_c \sim \langle K_x^2 \rangle^{-1/2}$$



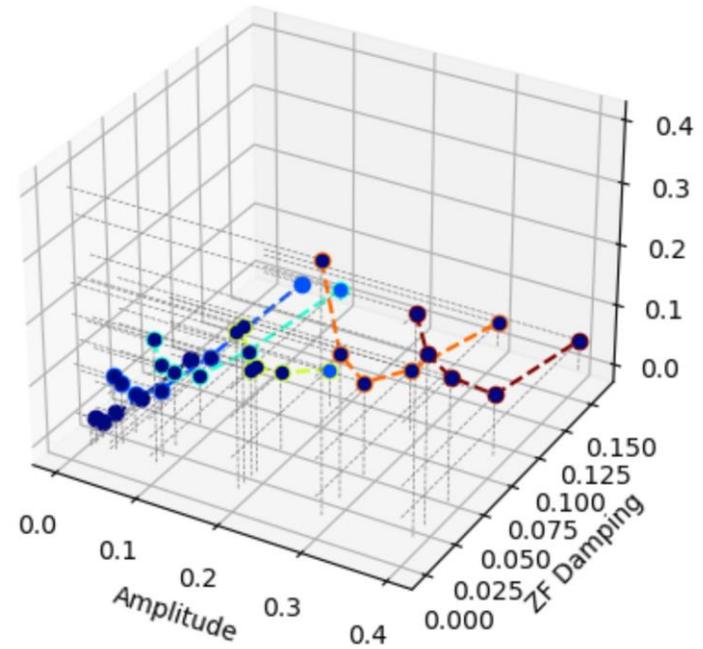
Kubo Number



Fluctuation intensity increases as drag increases, A increases

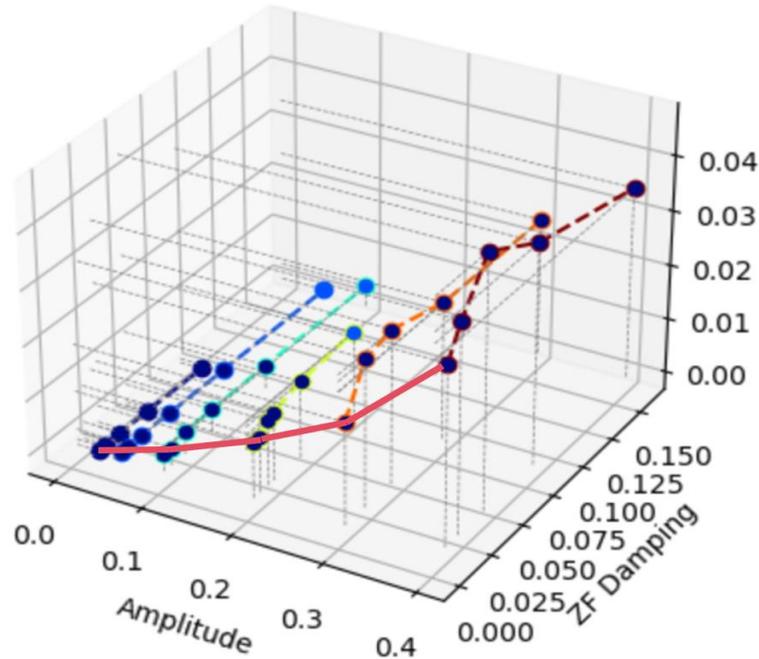
Kubo # tracks mixing
Control parameters set Ku

zonal_velocity



Zonal velocity decreases with increasing drag (clear)

→ Spreading and Fate of Zonal Flows



- Spreading rises for increased forcing, even for $\mu \rightarrow 0$
- Limits regime destroyed. How?
- ⇒ NL back-coupling from ZF necessary for spreading in systems with ZF

→ Animal Hunt for linear instabilities (KH, Tertiary ...) seems pointless in turbulence

→ Instead, $P_{Re} = -\langle \widetilde{V}_x \widetilde{V}_y \rangle \cdot \frac{\partial \overline{V}_y}{\partial x} \rightarrow$ Power transfer [fluctuations \rightarrow flow]

$P_{Re} < 0$: Wave \rightarrow ZF transfer

$P_{Re} > 0$: ZF \rightarrow Wave transfer \Rightarrow ZF decay

Aside:

- Of course, evokes ‘happy memories’ of studies of limitation of Dimits shift in G.K.
 - But mere identification of ‘Tertiary Instability’, “R-K.” etc not useful
 - Seek insight to and quantification of return of energy from Z.F. to turbulence, as control parameters scanned → Reynolds Power density
 - Goal is nonlinear ZF decay model for improved Predator-Prey system
- N.B. Reynolds power density used widely in data analysis

Quantifying Wave-ZF Power transfer

$$1/2 * \frac{\partial \bar{V}_y^2}{\partial t} = \omega_Z \langle \widetilde{v_x \widetilde{v_y}} \rangle - drag * \bar{V}_y^2$$



Reynolds power

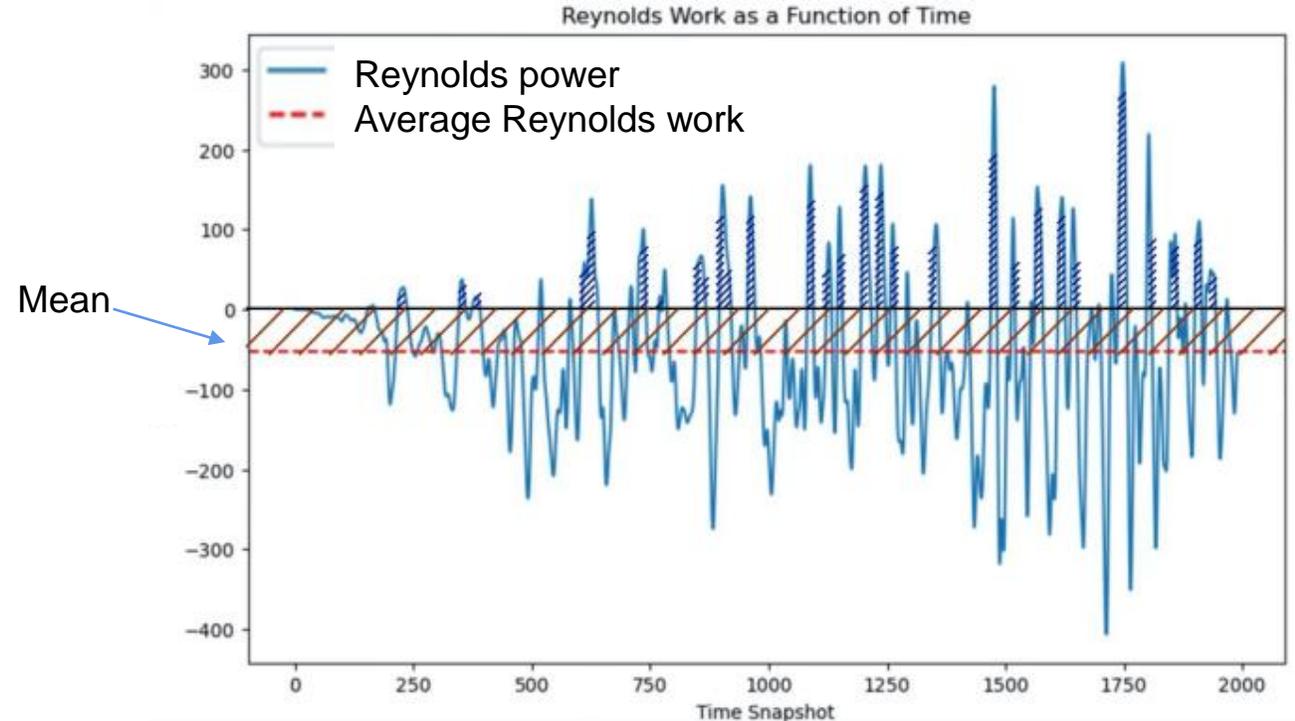
We quantify ZF → Waves Power Transfer as the ratio of the area above the axis to mean work done on the zonal flow.

N.B.:

$$P_{Re} = -\langle \widetilde{V_x \widetilde{V_y}} \rangle \cdot \frac{\partial \bar{V}_y}{\partial x}$$

‘Turbulent viscosity’ model fails capture 2 signs

Positive Reynolds power spikes → zonal flow vortex shedding events ?!



Reynolds power vs time

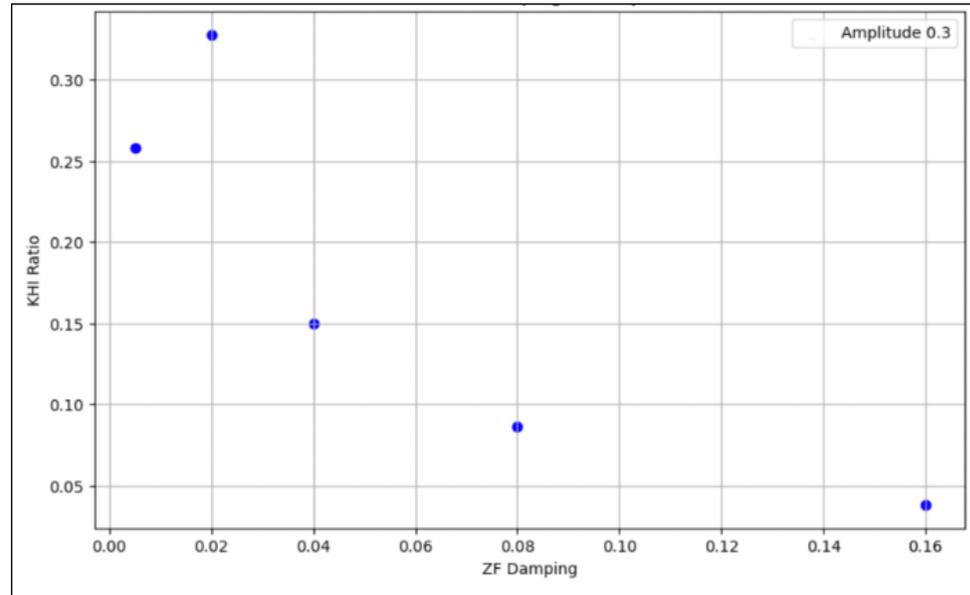
$P_{Re} < 0 \Rightarrow$ Wave → ZF transfer

$P_{Re} > 0 \Rightarrow$ ZF → Wave transfer

Results, Cont'd

P_{Re} ratio vs ZF damping

Dimits Regime



— The ratio generally decreases as a function of ZF damping

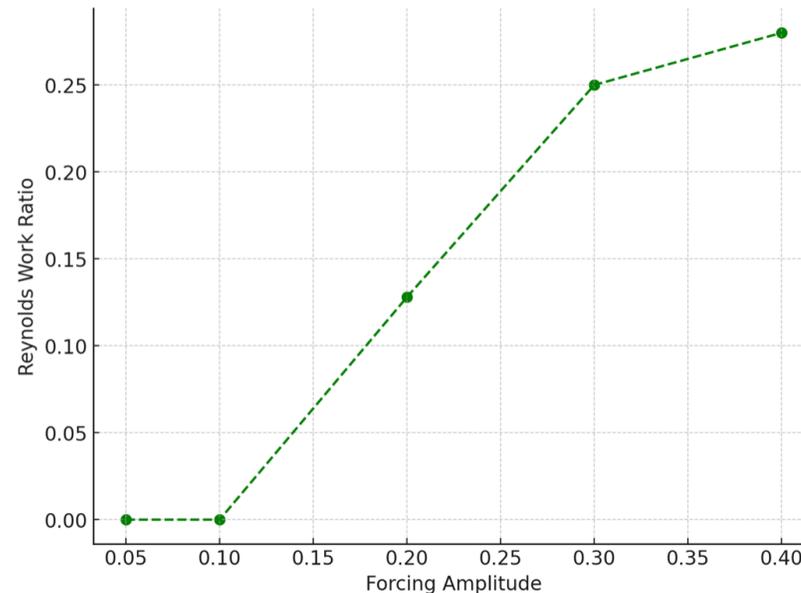


Damped Zonal Flow More Stable.

∴ Fewer Re spikes, as fewer vortex shedding events

Results, Cont'd, P_{Re} Ratio vs Forcing Strength

P_{Re} ratio vs forcing amplitude

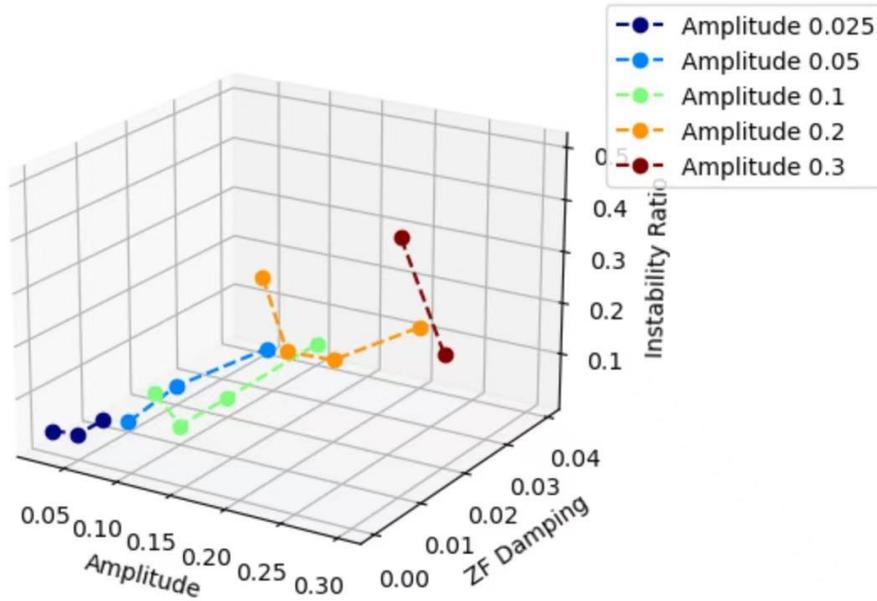


Preliminary
→ Explore other FOMs

- The ratio increases as a function of forcing strength
- Indicates that re-coupling of ZF energy to turbulence increases for stronger forcing
- This approach avoids instability morass → amenable to parametrization
- ➔ Significant nonlinear recoupling of energy to waves

P_{Re} Ratio vs A, μ

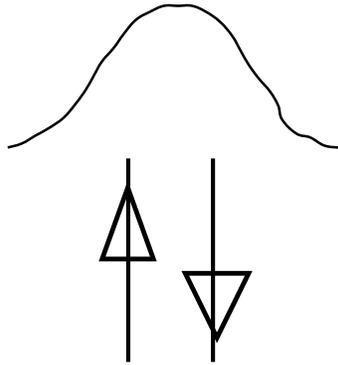
Instability Ratio vs Amplitude and ZF Damping



- P_{Re} back transfer increases with forcing, and as μ decreases
- Further analysis required

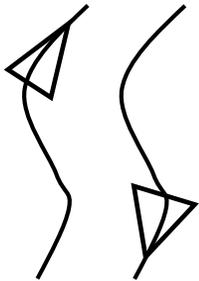
Related Problem: Jet Migration (Laura Cope)

i.e. - Here: ϵ



⇒ turbulence patch propagates,
drags ZF/Jet along, by generation
Near, but not at, Dimits regime?

- There:



⇒ Jet migrates
but Migration enabled by dynamics of fluctuation
field, especially zonalon

⇒ Zonalon → low mode # fluctuation co-located with Z.F.
→ necessary broken symmetry → propagation

$$V_{drift} \sim \mu^{0.7} \epsilon^{0.3} \quad \text{Microscopics?}$$



So Jet Velocity !?

→ As waves/eddies drag along zonal flow, Jet velocity(ala' Townsend) is related to Jet Migration.

so

→ Enstrophy Jet Velocity?!

$$V_{jet} = \langle v_r \tilde{u}^2 \rangle / \langle \tilde{u}^2 \rangle - \text{not experimentally accessible}$$

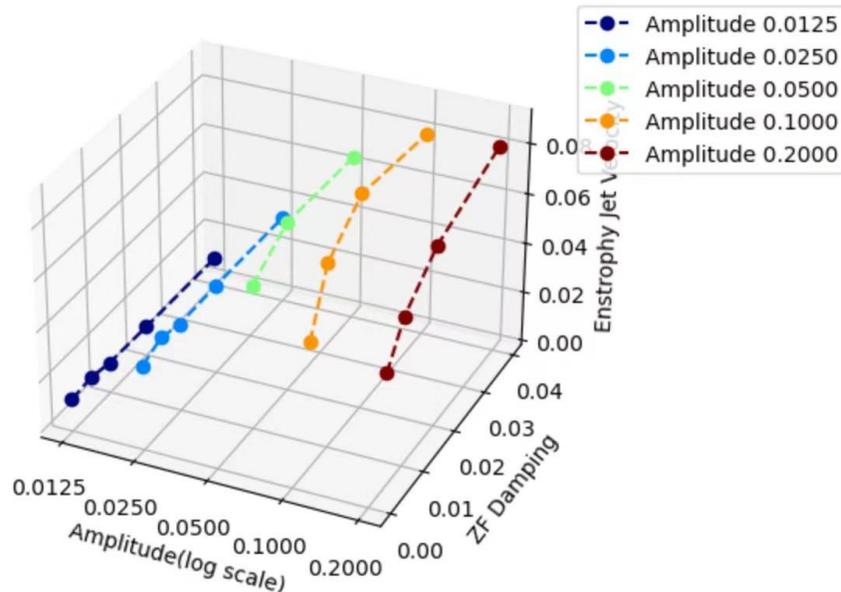
Long +

$$V_{jet} = \langle v_r \tilde{n}^2 \rangle / \langle \tilde{n}^2 \rangle$$

- Now familiar trends

- Seems semi-quantitatively consistent with Cope results.

Enstrophy Jet Velocity vs Amplitude and ZF Damping



Summary - Drift Wave Turbulence

- Spreading fluxes mapped in forcing, ZF damping parameter space
- Dominant mechanism \longleftrightarrow Ku (waves vs mixing) , Both waves and mixings in play.
What of Ku ~ 1 ? Interplay ?!
- Dimits-like regime discovered. Stationary ZF pattern.
- ZF disruption and quenching intimately linked to spreading
- $P_{Re} > 0$ bursts track breakdown of Dimits regime and onset turbulent mixing
Spreading increases.
- $P_{Re} > 0$ bursts likely due vortex shedding by zonal flows

→ General Summary

→ Coherent structures dipoles frequently mediate spreading

←→ underpin “ballistic scaling”

→ Spreading dynamics non-diffusive; Conventional wisdom misleading, or worse.

→ In DWT, wave propagation and turbulent mixing both drive spreading

→ ZF quenching critical to spreading in DWT. Power coupling most useful to describe ZF quench.

→ Closely related to jet migration.

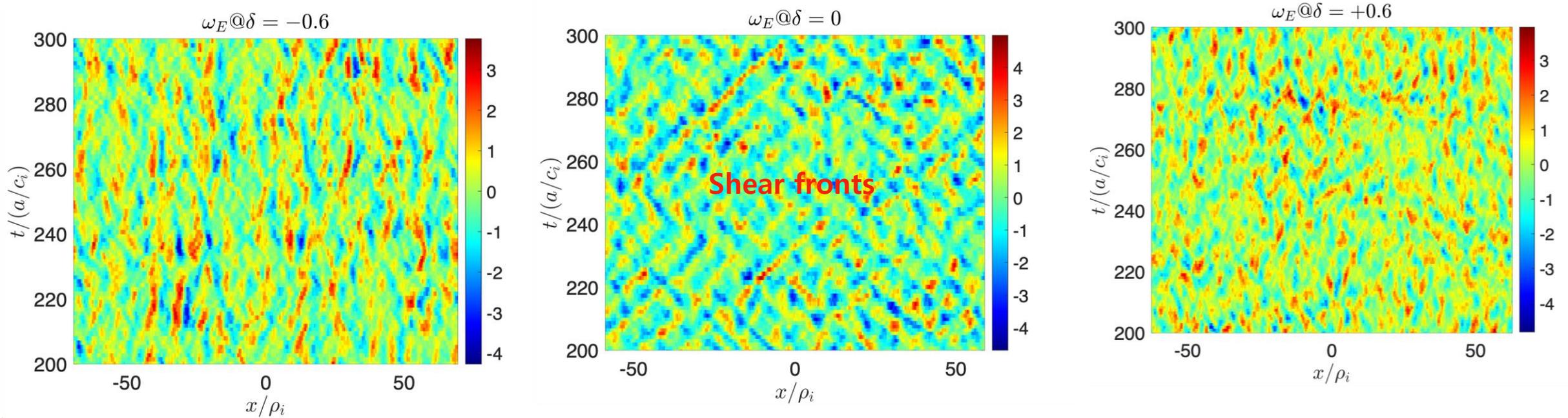
So: “ The more things change, the more they stay the same” – J.-B. Karr (1984)

- Collisionless, nonlinear damping/saturation of Z.F. remains poorly understood.
- Little progress beyond linear zoology, circa 2000. “Undead” theoretical question.
- Improved confinement in N.T. is related (R. Singh, P.D. submitted ‘24).

Collisionless GK for ITG using GENE →

NT enhanced confinement ↔ ZF resiliency

Zonal ExB shearing rates: spatiotemporal features

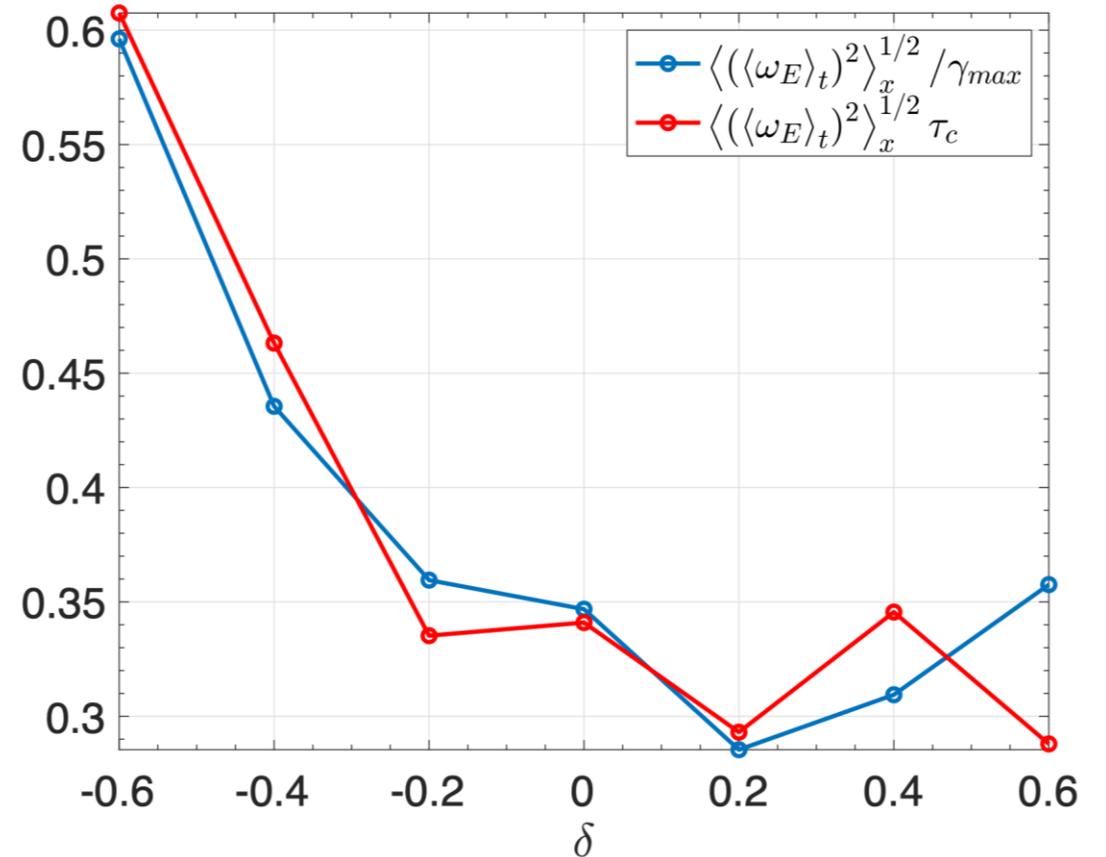
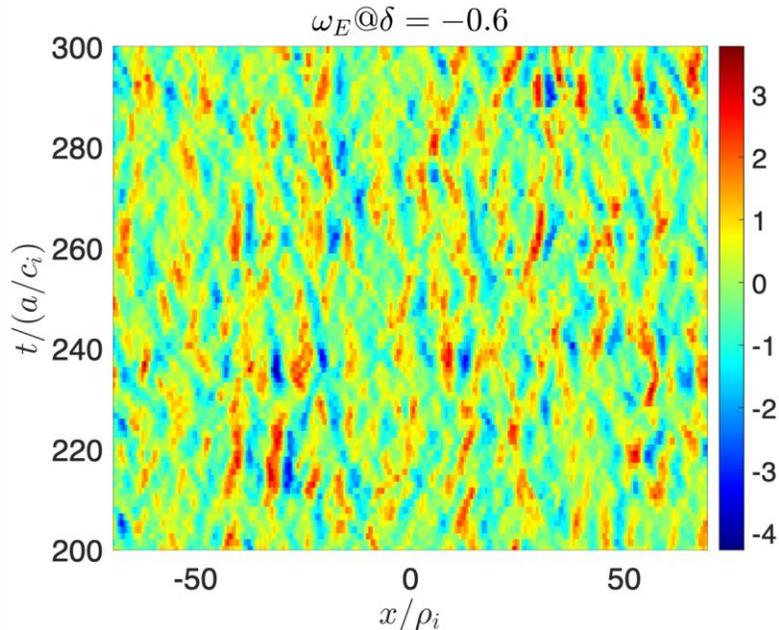


- Spatiotemporal patterns are highly sensitive to δ .
 - **Spatiotemporal shearing pattern more coherent for NT than for PT.**
 - Propagating shearing fronts \rightarrow dispersive feature for $\delta = 0$! Front speed $\sim 2.25\rho_* v_{th}$.
- **More coherent spatiotemporal shearing pattern for NT \rightarrow Stronger mean shearing effect for NT.**

Does NT reduce power transport out of Z.F.?

Figure of Merit

- All analyses point at the dimensionless parameter $\omega_E \tau_c$ or ω_E / γ_{max} as *figure of merit*.
- $\omega_E \tau_c$ higher for NT than for PT. Nicely correlates with the δ -trend of heat diffusivity.



→ Future Plans

- High resolution studies
- Understand ZF quenching physics and calculate power recoupling
- What is physics of $P_{Re} > 0$ bursts? - shedding? – how quantify?

- Spreading vs Avalanching. Relative Efficiency? Spreading and Transport?

More general:

- Is spreading mechanism universal? Seems unlikely
- Towards a model... $Ku \sim 1$ is an interesting challenge
- Relation/connection of DW+ZF spreading and Jet Migration (L. Cope)
- Is Directed Percolation of any use in this?

Details-??