

# **Microphysics and Macrophysics of the Tokamak (Edge) Density Limit**

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## Two Themes:

- Microphysics: Edge turbulence as  $n/n_G$  increases – shearing layers, Kubo number, Production Ratio
- Macrophysics: Radiative condensation → Strong particle flux → Scalings beyond Greenwald (especially ‘power’)

- Principal Collaborators:

Rongjie Hong, Ting Long, Rameswar Singh,  
Jingtong (Peter) Sheng, Ting Wu

- Ackn Collaboration of:

DIII-D, HL-2A, J-TEXT programs

# General Ideas

- Mixed experimental + theoretical approach
- Emphasis on dynamics, micro-macro connection, fluctuation structure
- Focus on L-mode DL
- Toward DL as a confinement transition
- N.B. Focus on edge fueling (relevance)

# Preview: From there to here of DL

- Greenwald (1988) scaling → Wrong → Power Scaling, Multiple Limits

$$\bar{n} \sim I_p / \pi a^2$$

- Sudo (1990) Scaling (stellarator) → mainstream, albeit incomplete → G + S unification

$$\bar{n} \sim P^{1/2} \text{ Transport Physics Central}$$

- DL as MHD + Disruption phenomenon → frequently, even usually, not  
Rebut, Gates, White
- DL as a 'Back-Transition' → from heresy to convention
- Radiation triggers MHD → Rad. triggers transport...
- Better Density 'Saturation' than 'Limit' !

# Basic Ideas

# Why Study Density Limits?

- Interested in max density - both line averaged and edge
- Constraint on operating space

$$\text{Lawson \#} \sim n T \tau_E$$

- Fusion power gain  $\rightarrow P \sim n^2 \sim \beta^2 B_0^4 R^3$  (MIT favorite)
- Emerging attractive feed back loop for burning plasma

$$P_{\text{fusion}} \sim n^2 \quad (\beta^2 + \text{ITG limits on } \nabla T_i \text{ steepening})$$

$$n_{\text{max}} \sim P^\alpha \quad (0 < \alpha < 1, \text{ but which } P \text{ in BP?})$$

# A Brief History of Density Limits

## → Conventional Wisdom

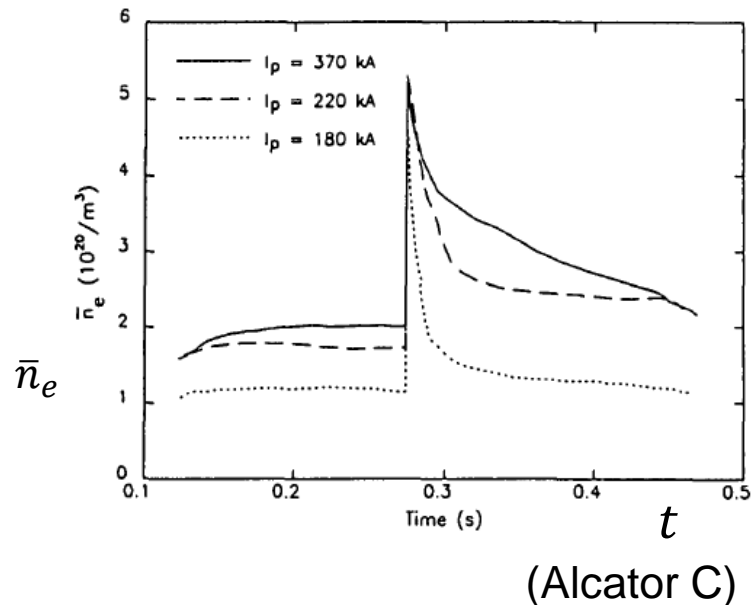
- Greenwald  $\bar{n}_G \sim I_p / \pi a^2$  (dimensions?)
- High density → edge cooling (radiation, transport?)
- Cooling edge → MARFE (Multi-faceted Asymmetric Radiation from the Edge) by Earl Marmor and Steve Wolfe

MARFE = Radiative Condensation Instability in Strong  $B_0$

after G. Field '64, via J.F. Drake '87 : Anisotropic conduction is key

- MARFE → Contract J-profile → Tearing, Island ... → Disruption  
after: Rebut, Hugon '84, ... , Gates ...
- But: more than macroscopics going on...

- Conventional Wisdom: Radiation + MHD (Rebut → Gates...)
- Argue: **Edge Particle Transport is fundamental**
  - ‘Disruptive’ scenarios secondary outcome, largely consequence of edge cooling, following fueling vs. increased particle transport → “Causality” issue
  - $\bar{n}_g$  reflects fundamental limit imposed by particle transport
- An Important Experiment (Greenwald, et. al. ‘88)



- Density decays without disruption after shallow pellet injection
- $\bar{n}$  asymptote scales with  $I_p$
- **Density limit enforced by transport-induced relaxation**
- Relaxation rate not studied
- Fluctuations?

# Shear Layer in L-mode? – Universal Feature of Edges

- Shear layer impacts/regulates edge turbulence even in Ohmic/L-mode, enhanced in H-mode

- Ritz, et. al. 1990

$v_{ph}$  - closed

$v_{pl}$  - open

density

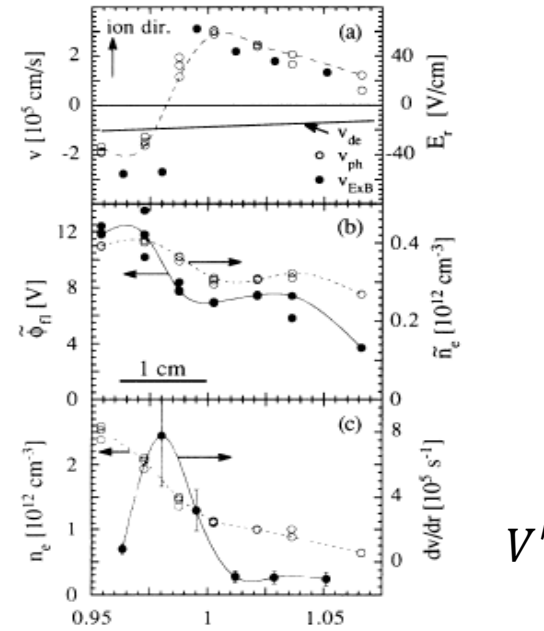


FIG. 1. Radial profiles for a discharge with  $B_e = 2$  T, plasma current of 200 kA, and chord-averaged density of  $n_{chord} = 2 \times 10^{13} \text{ cm}^{-3}$ . (a) Phase velocity of the fluctuations  $v_{ph}$  (closed circles),  $v_{E \times B}$  plasma rotation (open circles), and drift velocity  $v_{de}$ . (b) Density and floating potential fluctuations. (c) Density and velocity shear. The statistical error for individual shots is of order the symbol size and shot-to-shot reproducibility is given by the individual symbols. The systematic error in the plasma position is 0.5 cm or  $r/a \approx 0.02$ .

Shear layer

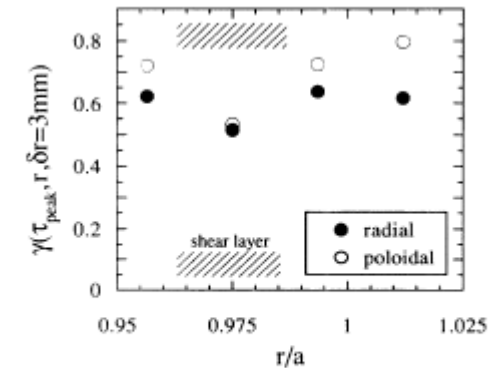


FIG. 3. Peak values of the normalized two-point correlation function for poloidally and radially separated probes with fixed separations of  $\delta r = 3$  mm.

Peak correlation

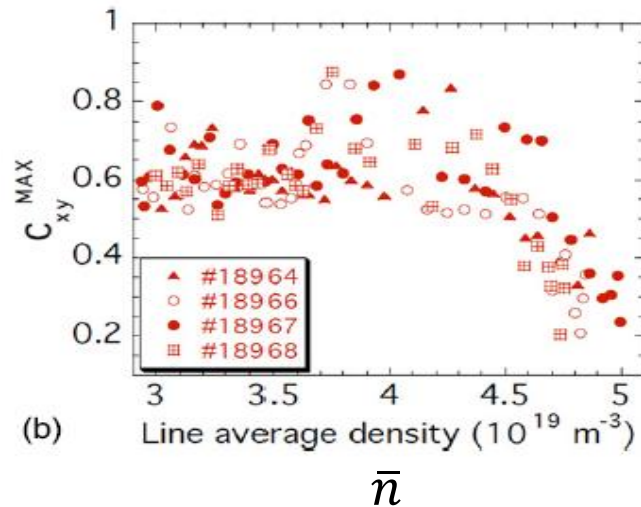
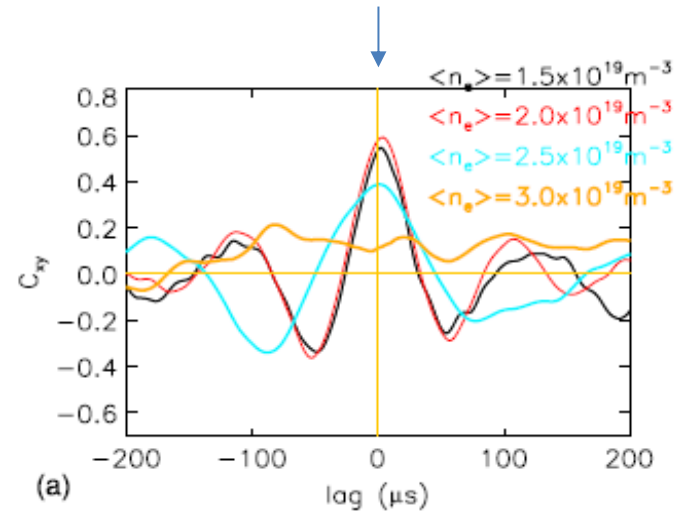
Title: “Evidence for Confinement Improvement by Velocity Shear Suppression of Edge Turbulence”

n.b. not H-mode!

➔ Role of Shear Layer in L → DL ?

# Toward Microphysics: Recent Experiments - 1

(Y. Xu et al., NF, 2011)



## LRC vs $\bar{n}$

- Decrease in maximum correlation value of LRC (i.e. **ZF strength**) as line averaged density  $\bar{n}$  increases at the edge ( $r/a=0.95$ ) in both TEXTOR and TJ-II.
- The reduction in LRC due to increasing density is also accompanied by a reduction in edge mean radial electric field (**Relation to ZFs**).

Is density limit related to edge shear decay?!

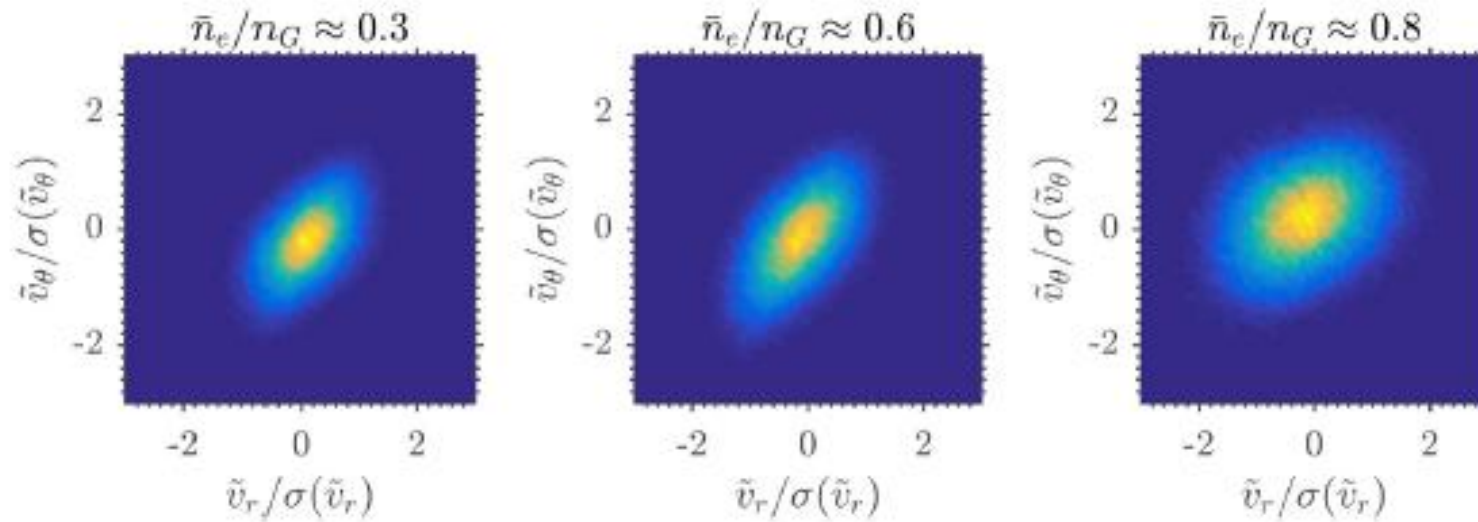
↓  
**Yes !**

See also: Pedrosa '07, Hidalgo '08 ...

Reynolds work (Flow production) drops as  $n \rightarrow n_G$  (Hong+ '18)

# Fluctuation + $n/n_G$ scan, R. Hong et. al. (NF 2018)

Distribution  
Fluctuating  
Velocities



- Joint pdf of  $\tilde{V}_r, \tilde{V}_\theta$  for 3 densities,  $\bar{n} \rightarrow n_G$
- $r - r_{sep} = -1cm$
- Note:



- Tilt lost, symmetry restored as  $\bar{n} \rightarrow \bar{n}_g$
- Consistent with drop in  $P_{Re}$  observed

→ Weakened shear flow  
production by Reynolds stress  
as  $n \rightarrow n_G$

# Reynolds Power (Flow Production)

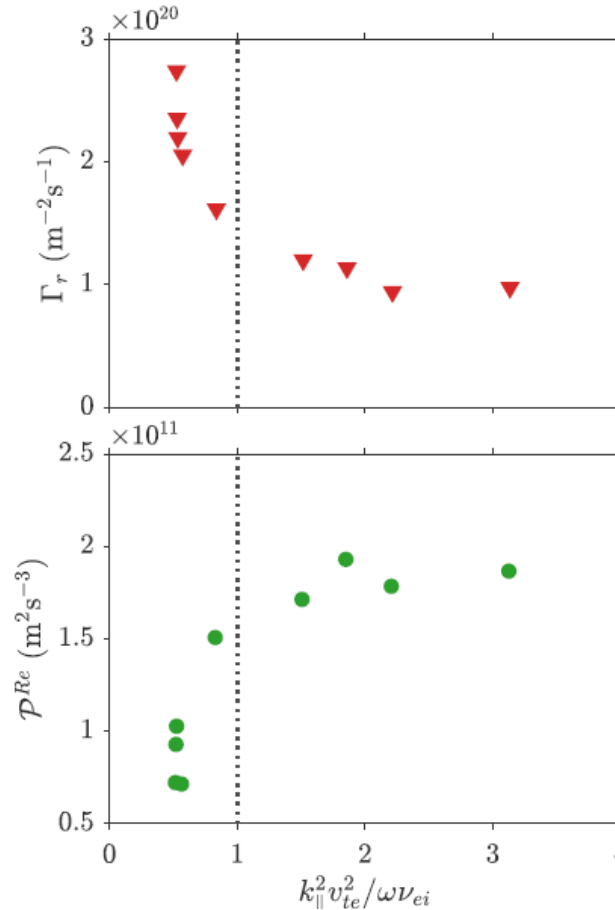
- Studies of  $P_{Re} = -\langle \tilde{v}_r \tilde{v}_\theta \rangle \partial \langle V_E \rangle / \partial r$  vs  $n/n_G$

$$\alpha = k_{\parallel}^2 V_{the}^2 / \omega \nu$$

adiabaticity

$\alpha \ll 1 \rightarrow$  cellular, hydro turbulence  $\sim$  2D

$\alpha \gg 1 \rightarrow$  drift wave turbulence



Particle flux  
surges for  $\alpha > 1$

$P_{Re}$  drops for  $\alpha < 1$

$\rightarrow$  Is DL evolution linked to degradation of edge shear layer ?

$\rightarrow$  Adiabaticity as key parameter

# An In-depth Look at More Recent Experiments

Ting Long, P.D. et. al. 2021 NF

Rui Ke, P.D., T. Long et. al. 2022 NF

See also: PD + Phil Trans 2024, T. Long, P.D.+, NF 2024

N.B. These experiments were “theoretically motivated”

# J-TEXT – Ohmic

- $B_T \sim 1.6 - 2.2 T$       $\frac{n}{n_G} \sim 0.7$       $n_G \sim 6.4 \rightarrow 9.3 \times 10^{19} m^{-3}$

- $I_p \sim 130 - 190 kA$       $\bar{n} \sim 2.0 - 5.3 \times 10^{19} m^{-3}$

- Principal Diagnostics: Langmuir Probes

- Shear layer collapses as  $n/n_G$  increases (1)

- Turbulence particle flux increases (3)

- Reynolds stress decays (2)

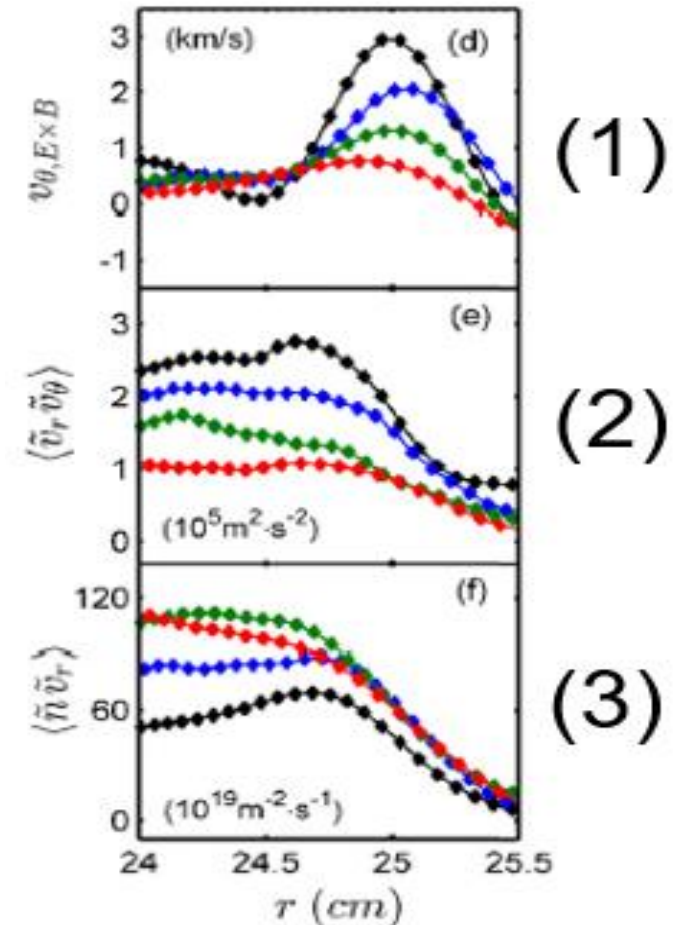
- Velocity fluctuation PdF  $\rightarrow$  symmetry

Black -  $0.3n_G$

Blue -  $0.34n_G$

Green -  $0.6n_G$

Red -  $0.63n_G$



# Mean-Turbulence Couplings

- In standard CDW model:

Production  $\equiv$  Input from  $\nabla n$

$$\delta n = \tilde{n}/n_0$$

$$P_I = -c_s^2 \langle \tilde{V}_r \delta n \rangle \left( \frac{1}{n_0} \frac{\partial \langle n \rangle}{\partial r} \right)$$

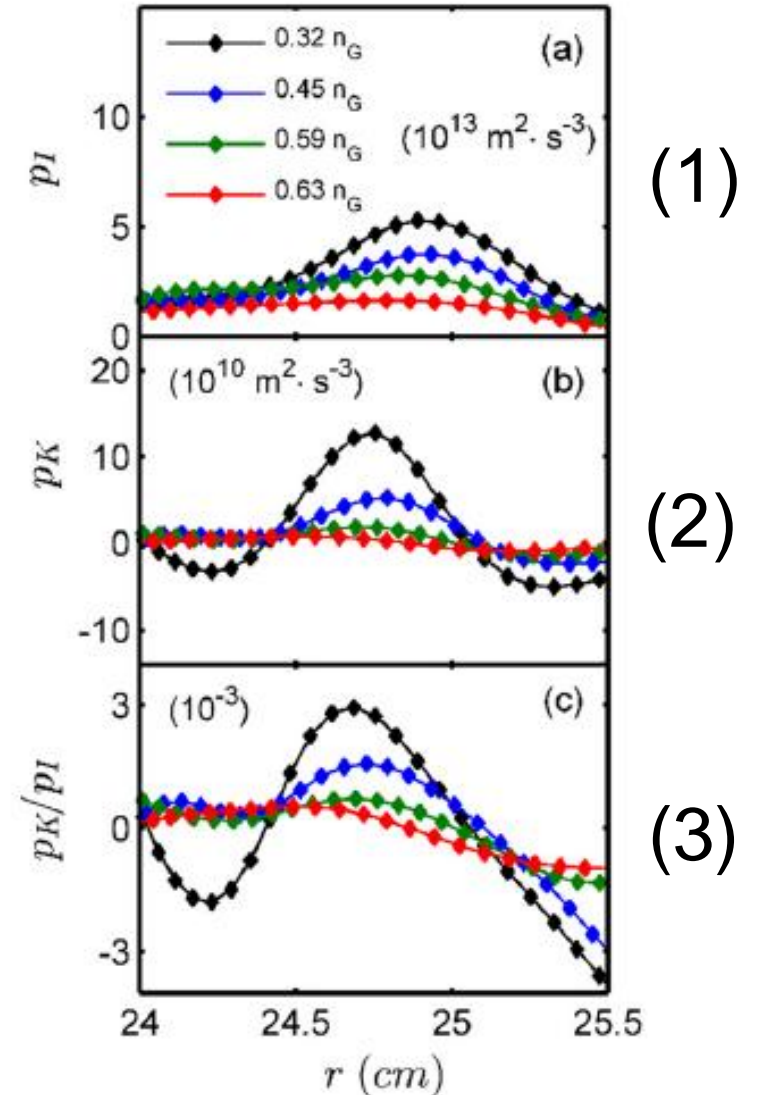
Reynolds Power  $\equiv$  Coupling to Zonal Flow

$$P_K = -\langle \tilde{V}_r \tilde{V}_\theta \rangle \langle V_E \rangle'$$

- Reynolds power drops as  $n/n_G$  rises (see Hong+, '18) (2)
- $P_K/P_I$  drops as  $n/n_G$  rises (3)

➔ Fate of the Energy ?

➔ Where does it go?



# Fate of the Energy ?

- Turbulence Energy Budget

$$\frac{\partial \varepsilon}{\partial t} + \frac{\partial}{\partial r} \langle v_r \varepsilon \rangle = P_I - \text{Dissipation}$$

↑ Triplet
↑ Production

↑ Spreading

$$\varepsilon = \varepsilon_k + \varepsilon_I \quad \varepsilon_I = \frac{c_s^2}{2} \langle (\tilde{n}/n_0)^2 \rangle \quad (\text{Internal Energy})$$

- Then  $P_S \rightarrow$  Power coupled to fluctuation energy flux  $\rightarrow$  Turbulence spreading

$$P_S = -\partial_r \langle \tilde{v}_r \varepsilon_I \rangle = -\partial_r \langle \tilde{v}_r \tilde{n}^2 c_s^2 \rangle / 2n^2$$

→ Turbulence Spreading Power

- Turbulence Spreading encompasses “Blob” and “Void” propagation

# Fate of the Energy, Cont'd

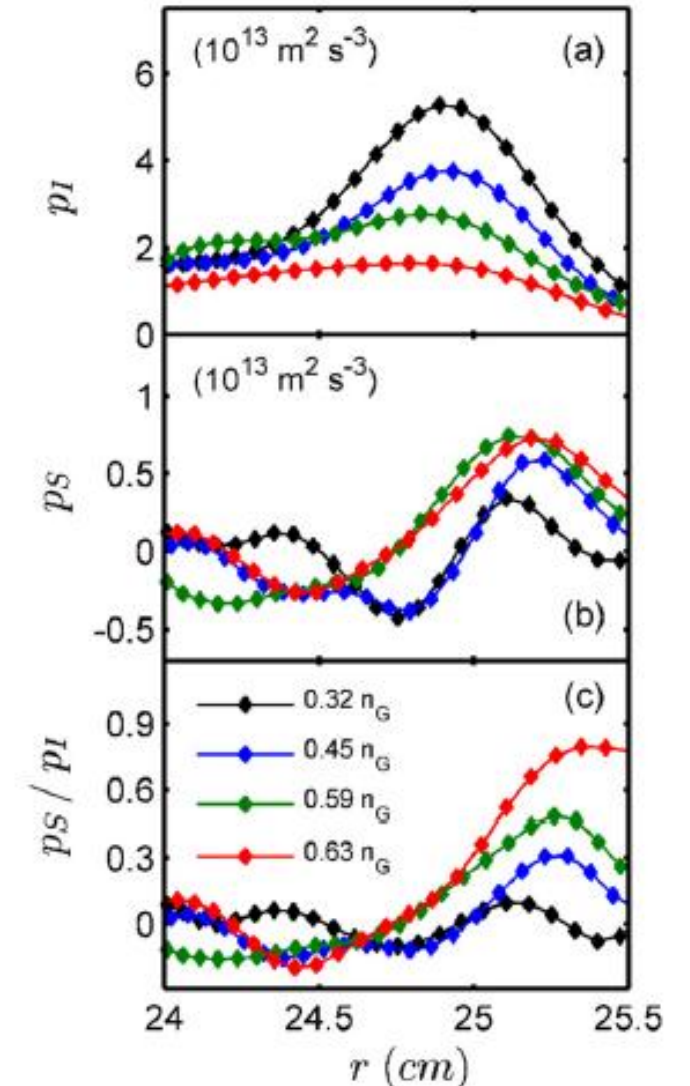
- Turbulence Spreading !
  - Reynolds power drops
  - $P_s$  increases; transitions  $P_s < 0$  to  $P_s > 0$
- Where does the shear layer energy go?

$$(P_k/P_I)_{peak} \times (P_s/P_I)_{peak} \sim 0.3, 0.5, 0.4, 0.4 \times 10^{-3} \text{ as } n/n_G \uparrow$$

$\approx$  constant

Energy diverted from shear layer to spreading at  $L \rightarrow DL$

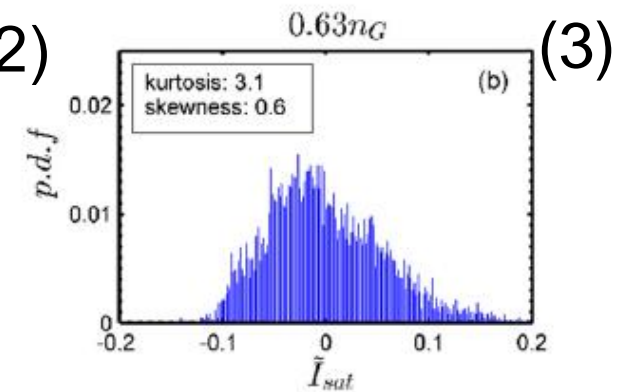
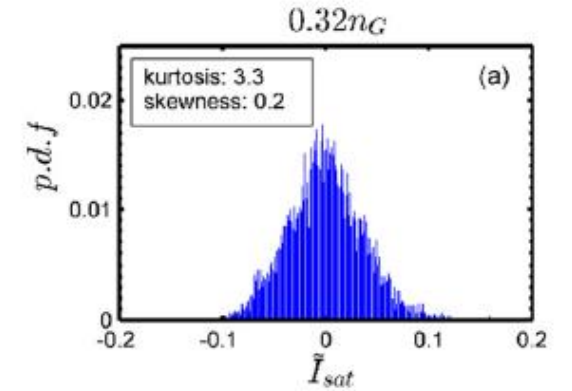
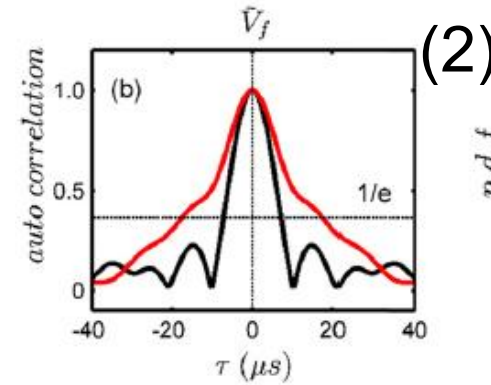
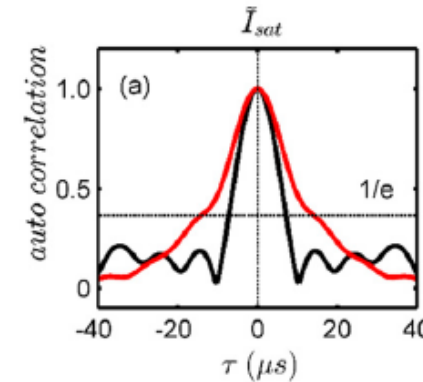
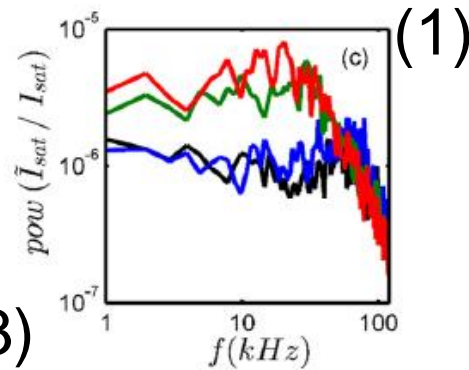
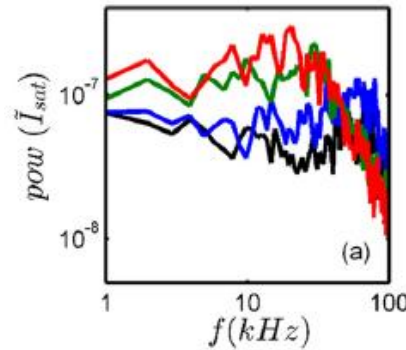
- N.B. Recent result (Long + 2024, N.F.):  $\delta$ (spreading flux) is more robust indicator of DL than  $\delta$ (particle flux)



# Characteristics of Spreading

- Low frequency content of  $\tilde{I}_{sat}/I_{sat}$  increases (1)
- $\tilde{I}_{sat}$  autocorrelation time increases (2)

Pdf  $\tilde{I}_{sat}$  develops positive skewness as  $n/n_G$  increases (3)



See also T. Long, P.D. NF 2024 for  $\tilde{n}$  skewness  $\leftrightarrow$  spreading correlation and in  $\rightarrow$  out symmetry breaking

# Characteristics of Spreading, Cont'd

- Enhanced turbulent particle transport events accompany L→DL back transition
- Events are quasi-coherent density fluctuations. Diffusive model of spreading  
dubious
- Localized over-turning events, small avalanches, “blobs”, ...

N.B. “The limits of my language means the limits of my world.”

- Ludwig Wittgenstein

- Blob ejection → recycling → cold neutral influx → cooling + MHD trigger

# Is there a key parameter? – Adiabaticity!

- Adiabaticity  $\alpha = k_{\parallel}^2 V_{the}^2 / \omega \nu$

$\alpha$  drops  $< 1$  as  $n/n_G$  increases

- $V'_E$  rises with  $\alpha \uparrow$

$\tau_{ac}$  decreases with  $\alpha \uparrow$

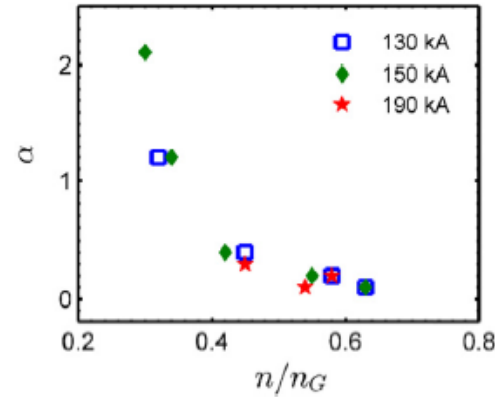
$\sigma(\tilde{I})/I$  decreases with  $\alpha \uparrow$

$P_S/P_I$  decreases with  $\alpha \uparrow$

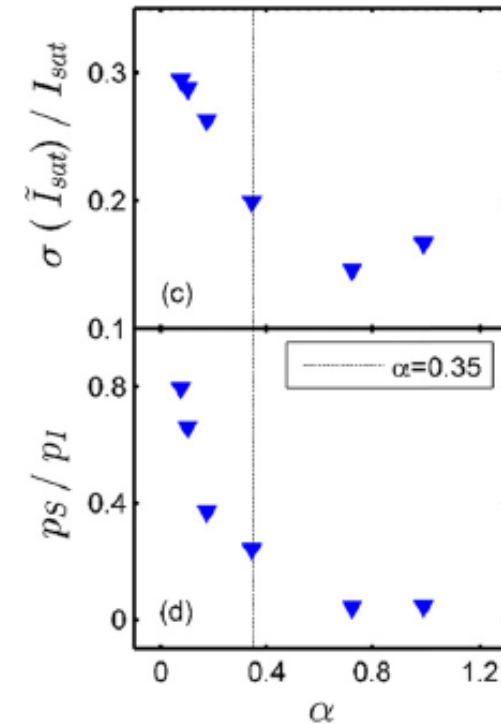
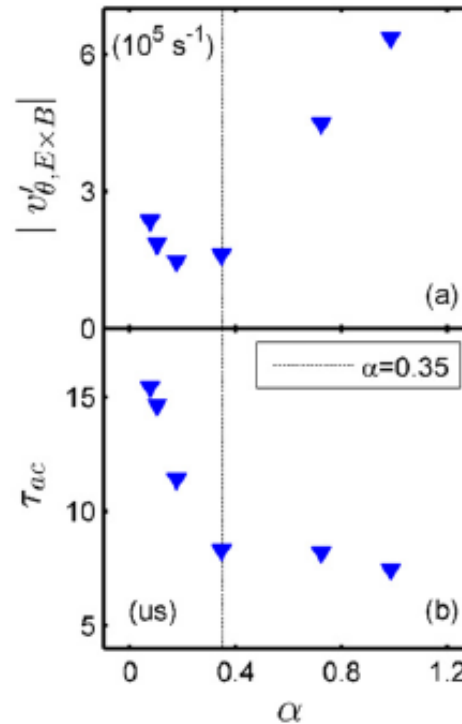
re: Adiabaticity

$\alpha > 1 \rightarrow$  drift waves

$\alpha < 1 \rightarrow$  convective cells



N.B.  $k_{\parallel} = 1/Rq$  assumed



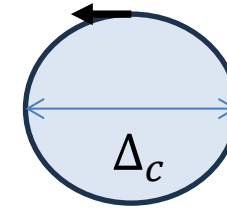
# Microphysics of Edge Turbulence vs $n/n_G$

(with J. (Peter) Sheng, T. Long, T. Wu)

→ How the Turbulence Evolves vs  $n/n_G, I_p$  etc.

# Microphysics of Edge Turbulence: A Tale of Two Numbers

- Kubo Number



$$\tau_{flight} \sim \Delta_c / \tilde{V}$$

- Also Strouhal Number

- $Ku \sim \tilde{V} \tau_{ac} / \Delta_c \sim \tau_{ac} / \tau_{flight}$       R. Kubo 1963

- Auto correlation time vs time to advect over scale  $\sim$  “time of flight”

- Derived from statistical theory of stochastic differential equations

# Kubo #, continued

$$\text{i.e. } \frac{\partial f}{\partial t} + \tilde{v} \cdot \nabla f + \dots = 0$$

Liouville, stochastic

$$\rightarrow f = \exp\left[-\int \tilde{v} \cdot \nabla dt\right] f_0$$

$$\langle f \rangle = \langle \exp\left[-\int \tilde{v} \cdot \nabla dt\right] \rangle f_0$$

Statistics,  $\tau_{ac}$  crucial

$$\rightarrow \langle f(1)f(2) \rangle = \left\langle \exp\left[-\int_{t_1}^{t_2} v \cdot \nabla dt\right] \right\rangle f_0(1)f_0(2)$$

- $Ku < 1$  (short  $\tau_{ac}$ )  $\rightarrow$  many kicks in  $\Delta_c$  scale
  - $\rightarrow$  suggestive of diffusion, QL
- $Ku > 1$  (long  $\tau_{ac}$ )  $\rightarrow$  long auto-correlation time vs time of flight
  - $\rightarrow$  coherency over many scatters
  - $\rightarrow$  trapping, percolation ... ? ( $\sim$  frozen scatters)

# Kubo # and MFE

- Reduced turbulence/transport models are unavoidable
- Virtually all reduced models are variations on the theme of quasilinear theory ...

$$\tau_{ac} < \tau_{TR}, \tau_{ev}$$

- Quasilinear theory  $\leftrightarrow$   $Ku < 1$

and/but:

- $Ku \sim 1$  is often said to be characteristic of plasma turbulence (simulations)

$$\tau_{ac} \sim \tau_{flight}$$

$\leftrightarrow$  “Mixing Length Estimates” -  $\tau_c \sim \tau_{flight}$

- What of  $Ku > 1$  ?

# “Contamination” vs Local Relaxation

## ↔ Turbulence Spreading, Avalanching, Blobs

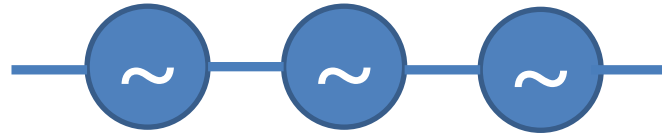
- Reduced models invariably posit local flux-gradient relation

yet:

- Intense consideration of turbulence spreading, avalanching etc. (cast of thousands)

↔ “Contamination Processes” – c.f. Y. Pomeau 1986

Re: Oscillator array



“... each oscillator if in a turbulent state may either relax spontaneously towards its quiescent state or contaminate its neighbors (if they are already turbulent, this interaction changes nothing).”

i.e. neighbor coupling ↔ drive

# “Contamination” vs Local Relaxation, cont’d

∴

- Must incorporate contamination processes – spreading, etc. → intensity fronts, etc
- Spreading = ‘contamination of stable region by unstable region’ is for too limiting a definition

∴

- Need measure of drive via contamination vs production by local relaxation

→ Production Ratio

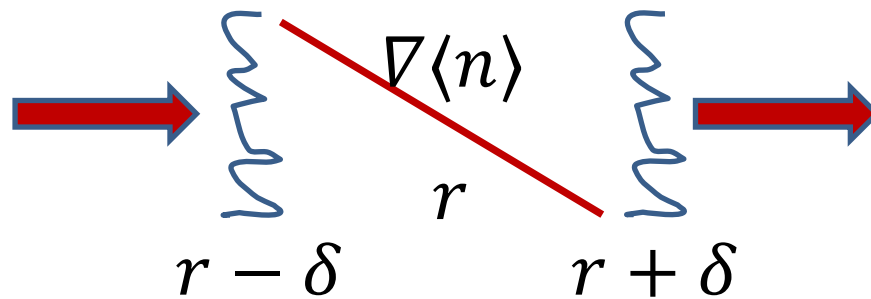
(Pomeau Number ?!)

# Production Ratio

(c.f. Manz; Wu, P.D.)

$$R_p(r, \delta) = \left[ \Delta \langle \tilde{v}_r \tilde{n}^2 \rangle / \left( - \int_{r-\delta}^{r+\delta} \langle \tilde{v}_r \tilde{n} \rangle \frac{\partial \langle n \rangle}{\partial r} \right) \right]$$
$$\Delta \langle \tilde{v}_r \tilde{n}^2 \rangle = \langle \tilde{v}_r \tilde{n}^2 \rangle \Big|_{r+\delta}^{r-\delta} \quad \delta > \Delta_c$$

- Ratio of net spreading / contamination drive to net local production in interval



- $R_p > 1 \rightarrow$  turbulence intensity front driven region  $\rightarrow$  contamination dominated
- $Ku \leftrightarrow R_p$  ?!

# Production Ratio

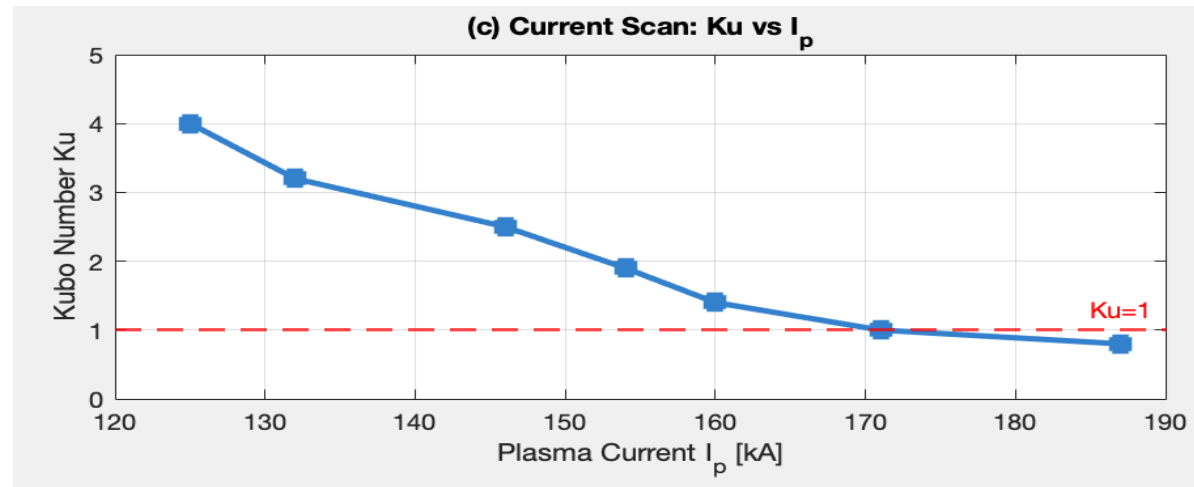
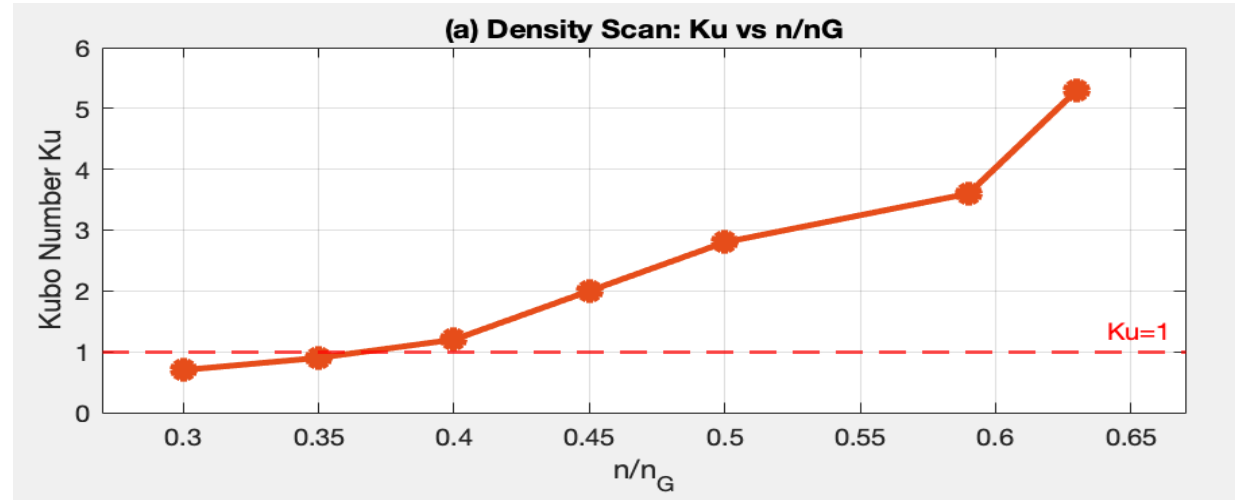
- Many versions of  $P_R$  possible, depending on quantity of interest
- For kinetic energy:

$$P_R = \Delta \langle \tilde{v}_r \tilde{v}^2 \rangle / \int_1^2 dr \left[ \frac{c_s^2}{R} \left\langle \frac{\tilde{v}_r \tilde{n}}{n_0} \right\rangle - \langle \tilde{v}_r \tilde{v}_\perp \rangle \frac{\partial}{\partial r} \langle v_\perp \rangle \right]$$

# The Data

- T. Long+, NF`22, `23, `24 – J-TEXT  $n/n_G$  study  
T. Wu+, NF`23 – HL-2A
- Edge turbulence in OH
- Emphasis on  $n/n_G$  scan, but not only
- Langmuir Probes – full characterization (BES velocimetry  $\rightarrow \tilde{v}_r$ , spreading)
- Also – Electrode bias data on J-TEXT
  - Control of shearing  $\leftrightarrow$  Fate of Ku,  $P_r$

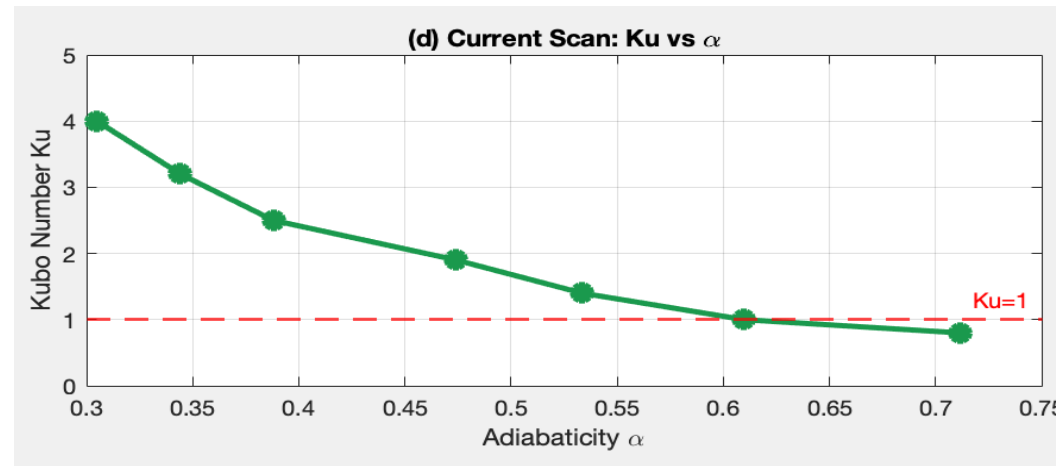
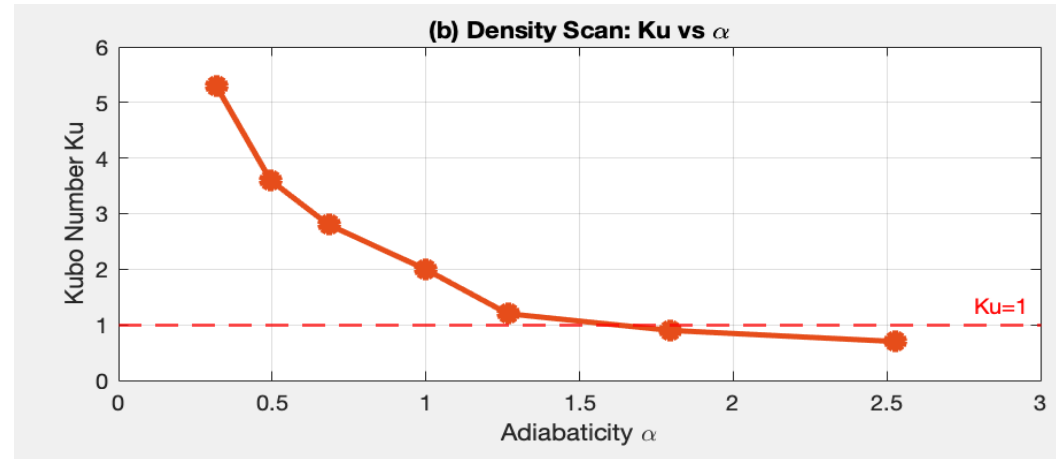
# Results



→  $Ku \gg 1$  at 'large' Greenwald fraction !

# Results, cont'd

- Physics parameter adiabaticity  $\alpha = k_{\parallel}^2 V_{th}^2 / \omega \nu$



- $Ku \gg 1$  in hydrodynamic regime  $\rightarrow$  coherent cells

## Results, cont'd

- What is the anomalous physics at  $Ku \gg 1$  ?



- Blobs ... (see previous)

- Interesting to define  $Ku_{Blob}$

$$Ku_B \sim u_b \tau_{L,b} / \Delta_b$$

~ analogous to Knudsen #

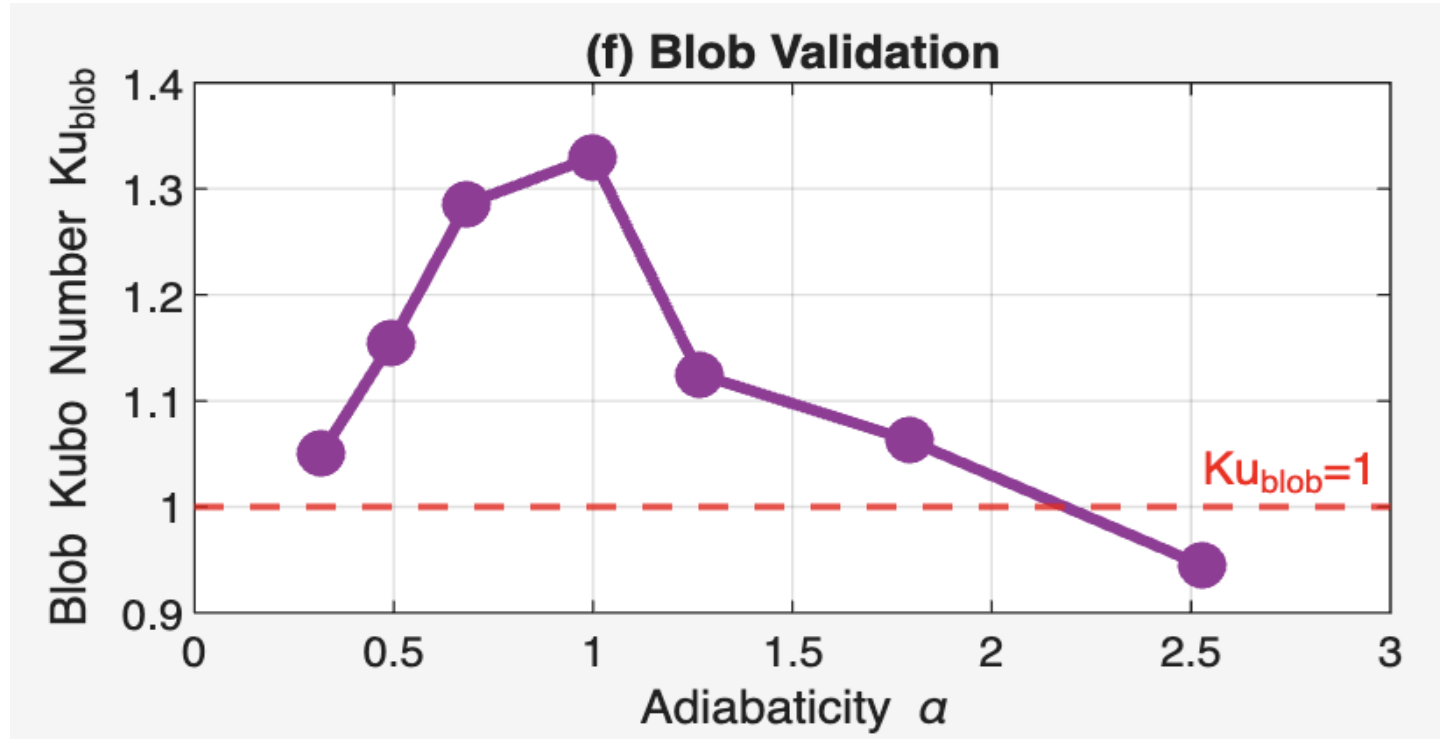
$u_b$  ~ blob speed

$\tau_{L,b}$  ~ blob life time

$\Delta_b$  ~ blob scale

- If turbulence → ensemble of blobs, expect  $Ku_B \rightarrow 1$

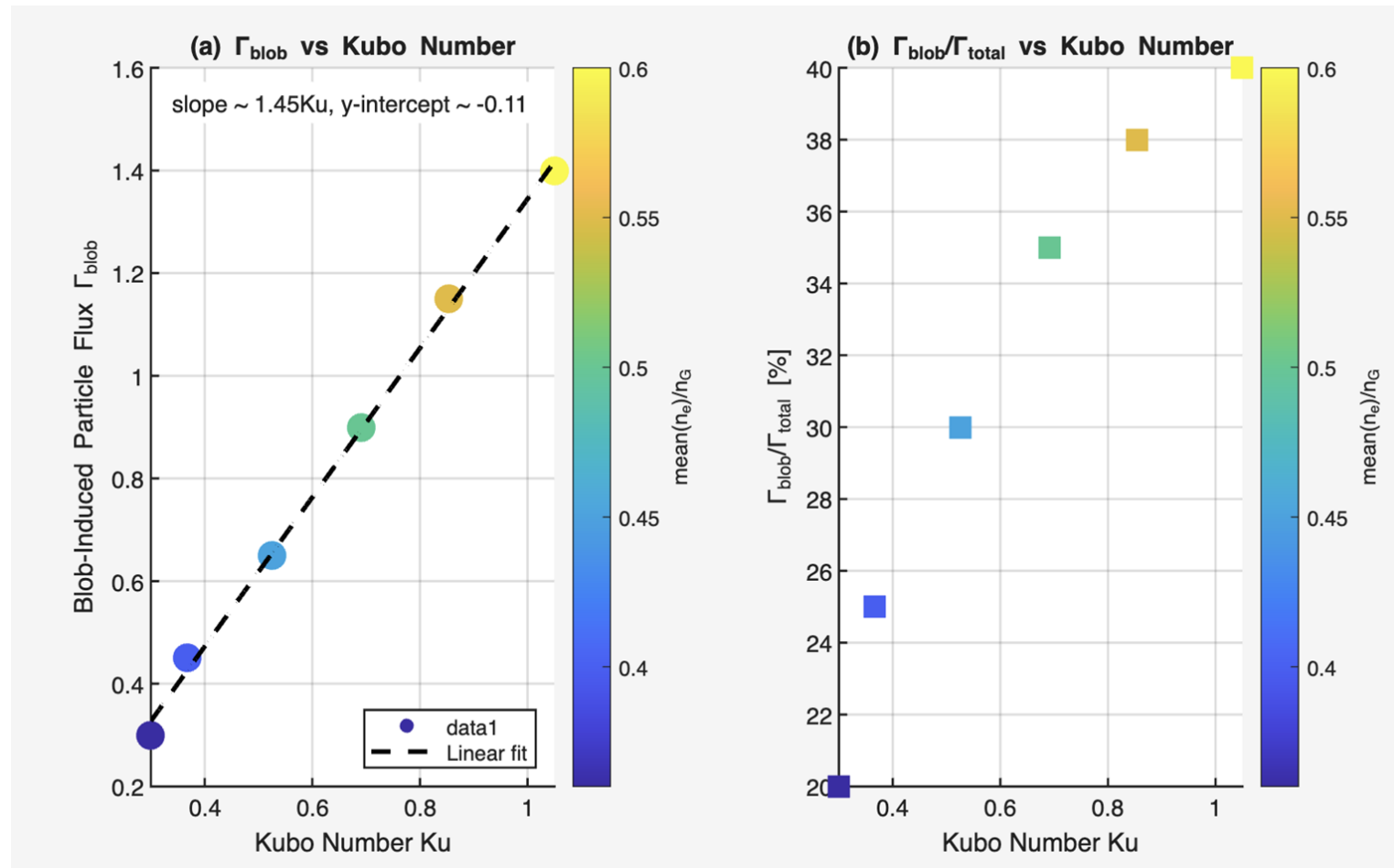
# Results, cont'd



$$Ku_{blob} \sim 1$$

- At high  $nn_G$ , turbulence best modeled as a gas of blobs
- Message for edge modellers in relevant BP regimes.

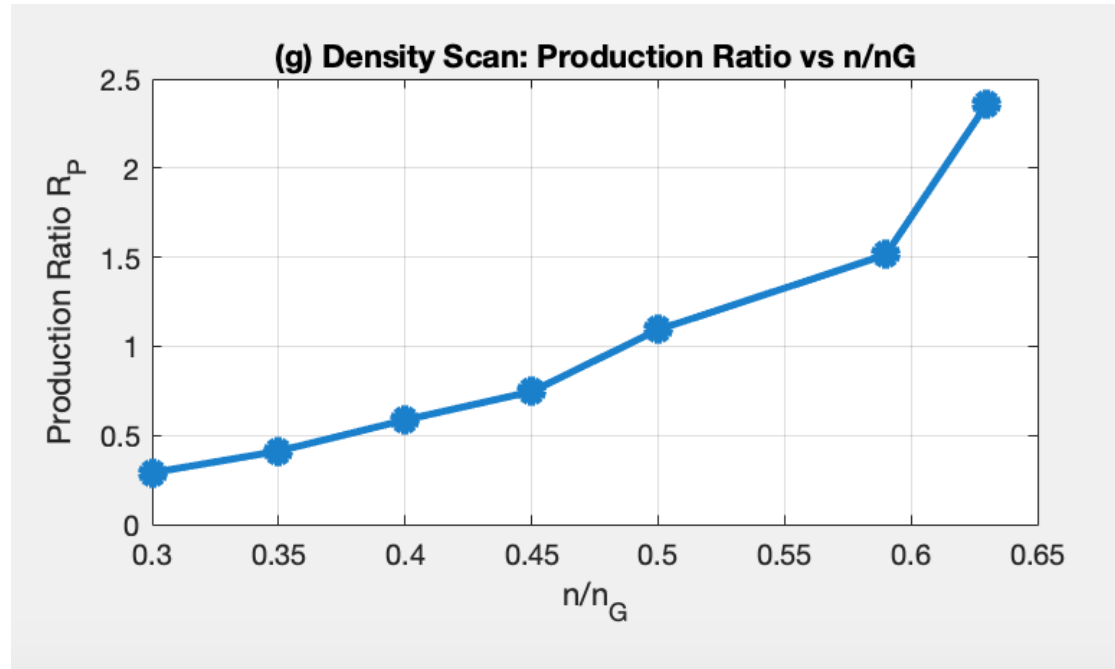
# Results, cont'd



- Fraction of particle flux carried by blobs increases with  $Ku$
- $\Gamma_{blob} \equiv$  particle flux windowed for blob contribution  $\rightarrow$  conditional average

# Results, cont'd

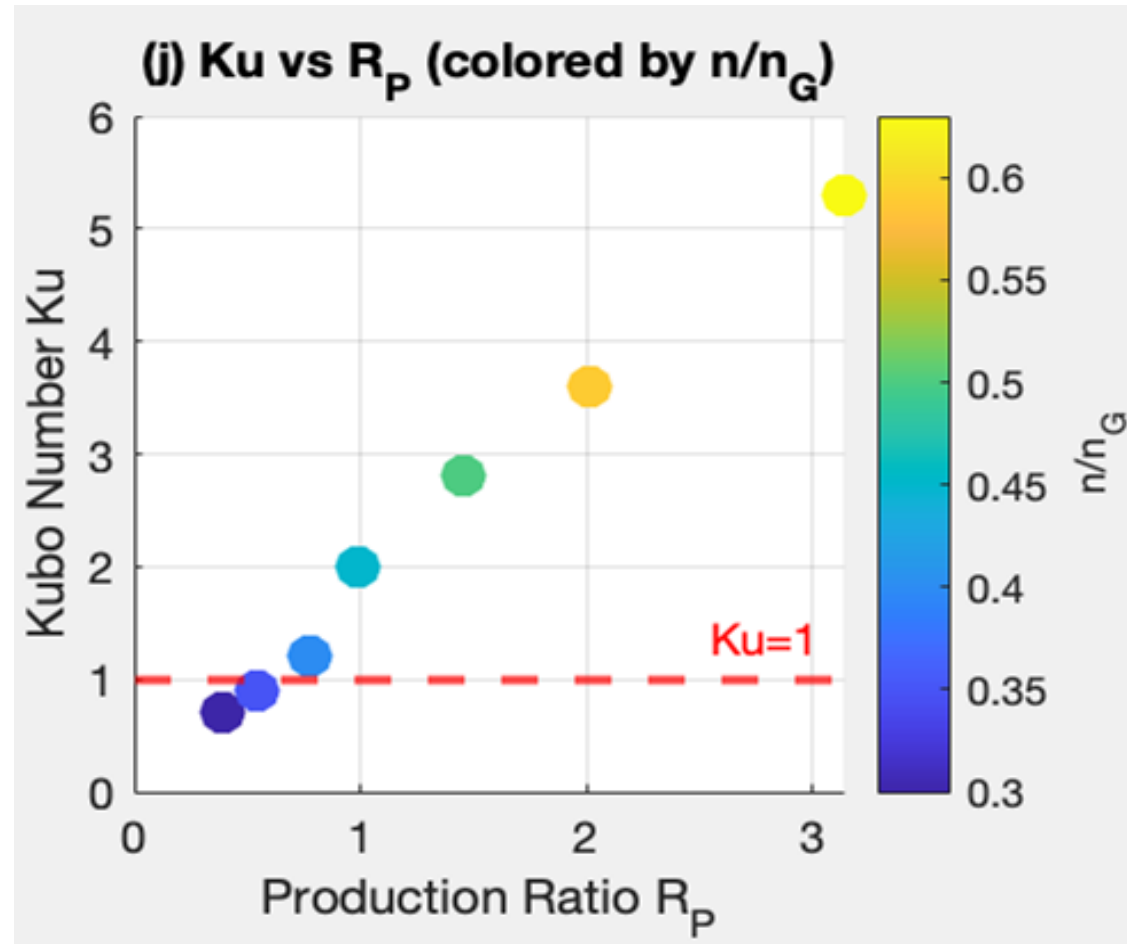
- What are consequences for relaxation?



- Production ratio  $R_p$  exceeds unity for large  $n/n_G$

→ Blob turbulence pulse 'contamination' controls drive

# Results, cont'd



- $Ku$  closely tracks  $R_p$
- $R_p > 1 \rightarrow Ku > 1$

Long  $\tau_{ac}$  phenomena linked to (generalized) turbulence spreading pulse

# Results – General Lessons

- $Ku > 1$  happens in relevant, macroscopically stable edge regimes – i.e. DL
- $Ku \sim P_R \rightarrow$  Spreading, contamination – due blobs – underlies  $Ku > 1$  phenomena
- General question: Long  $\tau_{ac}$  phenomena  $\leftrightarrow$  contamination connection (DIII-D upcoming – Kin, P.D., Zeyu Li)

# Results – General Lessons

- $Ku > 1$  regimes – mode ensemble description is pointless
- Better: inhomogeneous turbulence closure model  
Output:  $\langle \tilde{n}^2 \rangle, \langle \tilde{v}^2 \rangle$  etc.

# Results – next

- “But what of gyrokinetics ?”
- Diagnosing/quantifying turbulence spreading is challenging in GK
  - local relaxation rarely vanishes...
- Kinetic Production Ratio (ITG)

$$P_R = d \int d^3 v \langle \tilde{v}_\perp \tilde{g} \tilde{g} \rangle / R$$

$$R = \int_{r-\delta}^{r+\delta} dr \int d^3 v \frac{|e| \langle f \rangle}{T} \left[ \left\langle \frac{\partial \phi}{\partial t} \tilde{g} \right\rangle + V(E) \left\langle \frac{\partial \phi}{\partial y} \tilde{g} \right\rangle \right]$$

$$d(F) = F(r + \delta) - F(r - \delta)$$

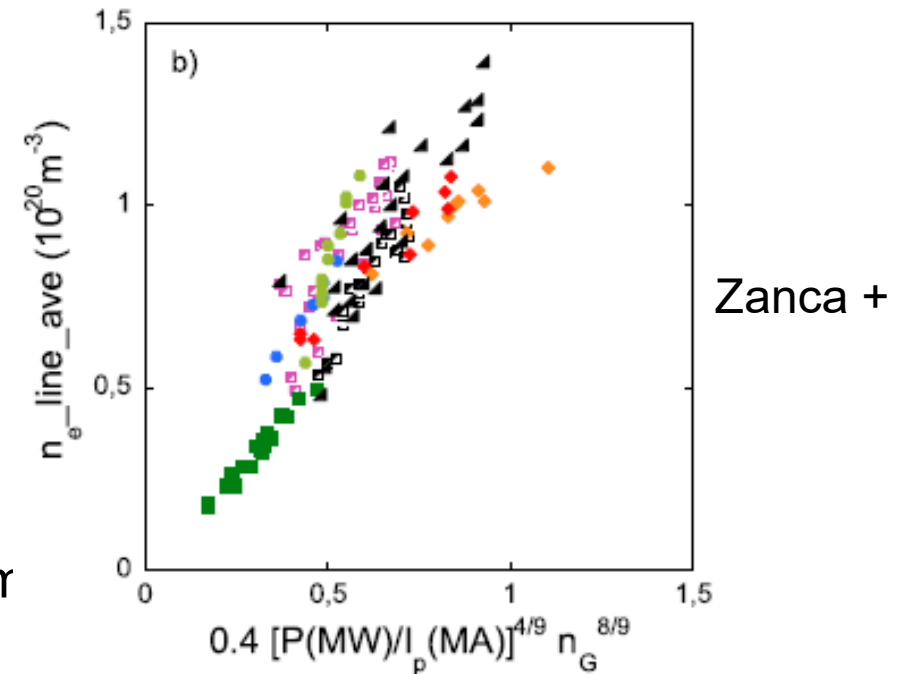
- Extensions possible...  $P_R$  for moments, i.e.  $\langle \tilde{v}_r \tilde{T} \tilde{T} \rangle$
- Formulate from 2pt theory  $\langle \tilde{g}(1) \tilde{g}(2) \rangle$

# Density Limits beyond Greenwald ...

→ Power (Heat Flux) Scaling

# Power Scaling and Physics of L-mode Density Limit (Singh, P.D. PPCF 2022)

- Power Scaling is an old story, keeps returning
  - Zanca+ (2019) fits  $\rightarrow \bar{n} \sim P^{1/4}$
- ↑
- Giacomini+: Simulations recover power scaling
  - Observe:  $Q_i|_{\text{bndry}}$  will drive shear layer  $\rightarrow$  LH mechanism
  - So:  $P_{\text{scaling}} \leftrightarrow$  shear layer physics: a natural connection
- also support vs radiation
- $Q_i, Q_e$  at boundary as physical quantities



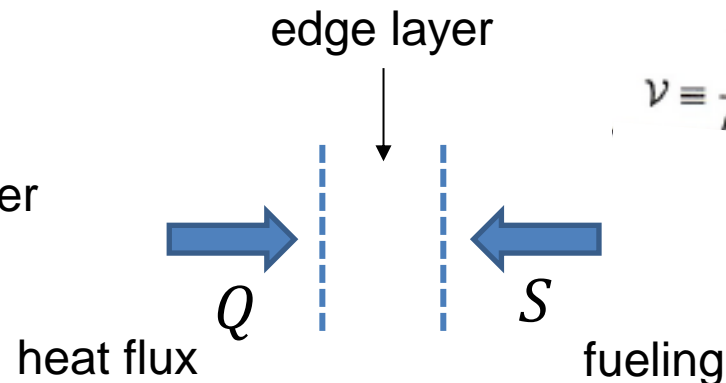
# Expanded Kim-Diamond Model

- KD '03 – useful model of L→H dynamics (0D)
- See also Miki, P.D. et al '12, et. seq. (1D)
- Evolve  $\varepsilon, V_{ZF}, n, T_i, V'_E$

↔

- Treats mean and zonal shearing
- Separates density and temperature contributions to  $P_i$
- Heat and particle sources  $Q, S$

N.B. i) ZeroD → interpret as edge layer  
 ii) Does not determine profiles  
 iii) Coeffs for ITG



$$\frac{\partial \varepsilon}{\partial t} = \frac{a_1 \gamma(N, T) \varepsilon}{1 + a_3 \nu^2} - a_2 \varepsilon^2 - \frac{a_4 v_z^2 \varepsilon}{1 + b_2 \nu^2} \quad \text{Fluctuation Intensity}$$

$$\frac{\partial v_z^2}{\partial t} = \frac{b_1 \varepsilon v_z^2}{1 + b_2 \nu^2} - b_3 n v_z^2 + b_4 \varepsilon^2 \quad \text{Zonal Intensity}$$

$$\frac{\partial T}{\partial t} = -c_1 \frac{\varepsilon T}{1 + c_2 \nu^2} - c_3 T + Q \quad \left\{ \begin{array}{l} T_i \\ Q \rightarrow \text{power} \end{array} \right.$$

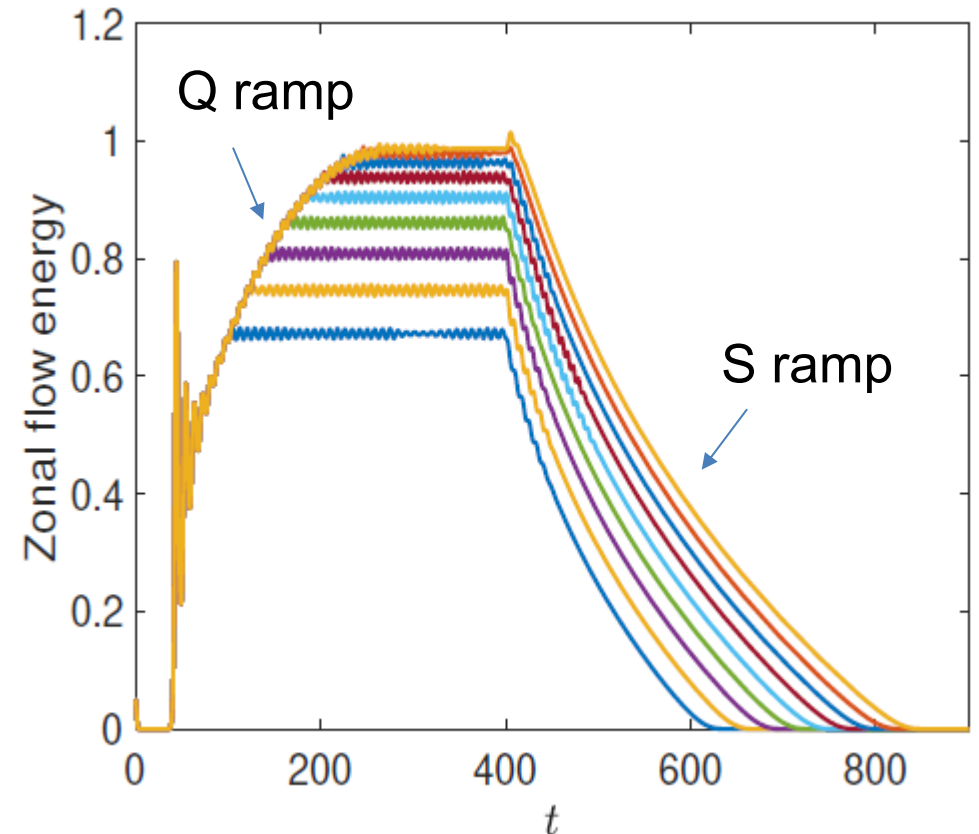
$$\frac{\partial n}{\partial t} = -d_1 \frac{\varepsilon n}{1 + d_2 \nu^2} - d_3 n + S \quad \left\{ \begin{array}{l} n \\ S \rightarrow \text{fueling} \end{array} \right.$$

$$V'_E = -\rho_i v_{thi} L_n^{-1} (L_n^{-1} + L_T^{-1}) \quad \text{Shear (mean)}$$

$$\nu \equiv \frac{V'_E a}{\rho^* v_{thi}} = -\frac{n_0}{n} \mathcal{N} \left( \frac{n_0}{n} \mathcal{N} + \frac{T_0}{T} \mathcal{T} \right)$$

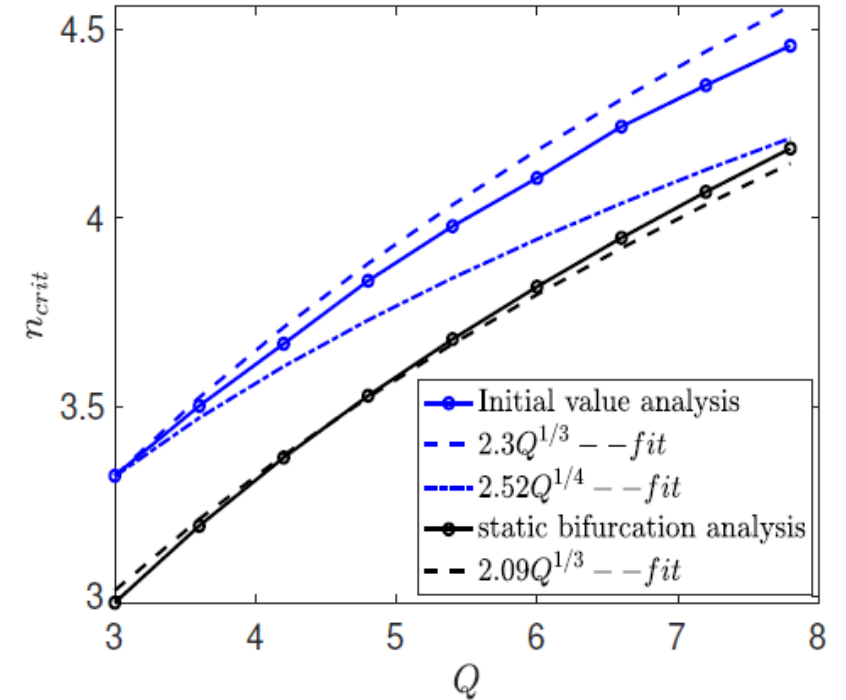
# L → DL Studies: Shear Layer Physics ↔ Power Scaling

- Look for shear layer collapse
  - $Q$  ramp-up to L-mode, followed by  $S$  ramp-up
  - Oscillations → predator-prey cycles
  - $n$  for ZF collapse increases with  $Q$
- scaling of  $n_{\text{crit}}$  emerges



# Power Scaling: LDL

- $n_{crit} \sim Q^{1/3}$
- Distinct from Zanca, but close (model)
- In K-D, with neoclassical screening  $n_{crit} \sim I_p \rightarrow I_p^2$
- Physics is  $\gamma(Q)$  vs ZF damping
- Shear layer drive underpins power scaling



Physics:  $Q_i \rightarrow$  Turbulence  $\rightarrow$  Reynolds Stress  $\rightarrow$  ZF shear

Increased ZF damping  $\rightarrow$  Confinement degradation

NB: Unavoidable model dependence in scalings

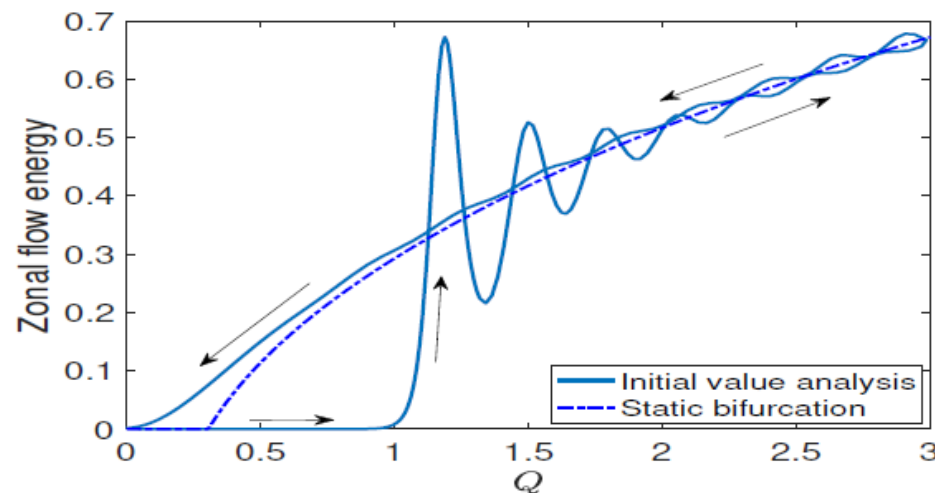
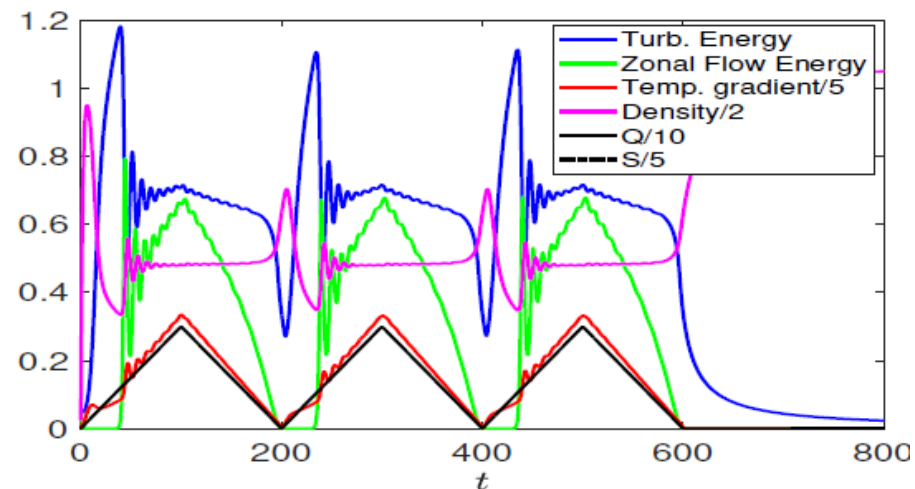
# Beyond Scalings: L→DL 'Transition' Physics

“If it Flux Like a Duck... (M.N. Rosenbluth, after F. Wagner)”

- Hysteresis ! in  $\varepsilon_{ZF}$  vs  $Q$  → Critical slowing down effect
- Expected, given 2 states transport
- Not familiar bistability !
- Physics prediction... beyond scaling

Also:

- Is there torque effect of density limit, i.e.  $\nabla P/n$  vs  $B_\theta V_\phi$  ?
- Torque  $\leftrightarrow V'_E$  → Mean field  
→ Reyn. stress coherence



# Partial Conclusions: $V'_E$ as Ubiquitous Edge Order Parameter

- Density limits as “back-transition” phenomena;  $V'_E$  physics crucial
- L-DL mechanism:
  - Shear layer degradation
  - Strong turbulence spreading → Blob emission
- $\alpha$  is key parameter, but not only
- Scalings of L-DL emerge from zonal flow physics
  - $I_p$  scaling → neo dielectric
  - $P$  scaling → Reynolds stress, radial force balance
- Novel hysteresis evident in L-DL dynamics
- Back Transition is in state of edge plasma.

# Recent L-DL Experiments in DIII-D with NT

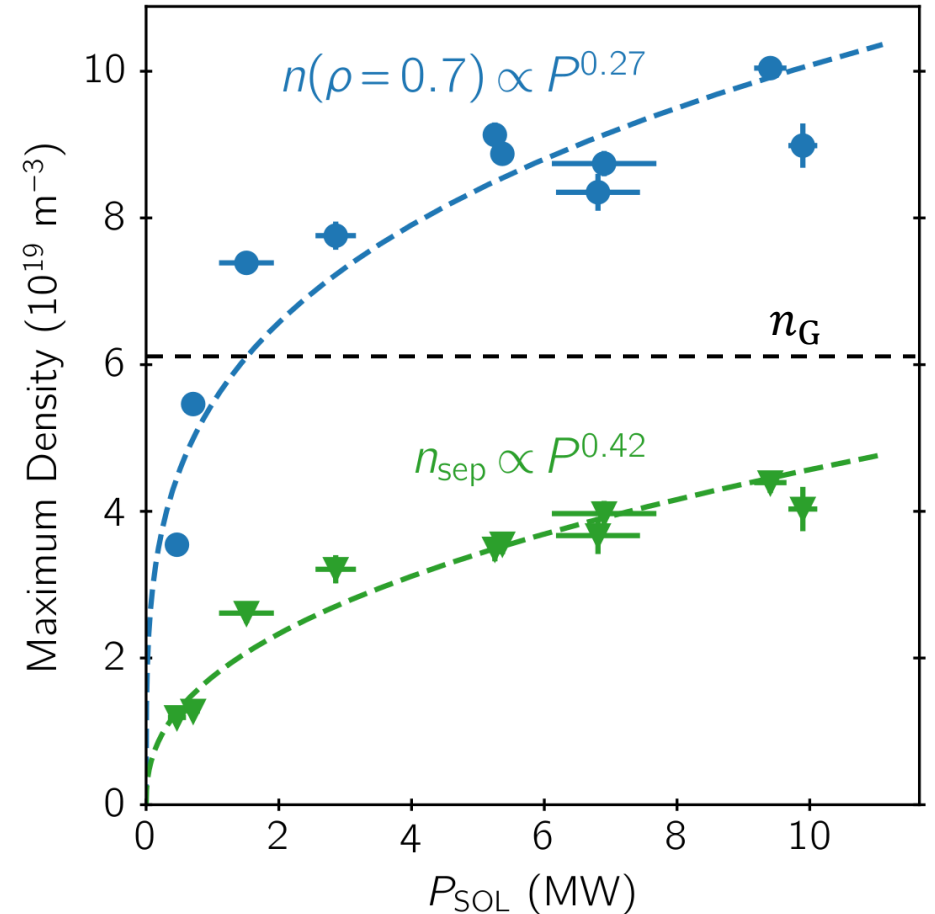
R. Hong, P.D., O. Sauter + → NF 2026

O. Sauter, R. Hong + → N.F. 2025

- N.B. :
- NT suppresses L→H transition, even at high power
  - Extends dynamic range of  $P$  for L-DL studies

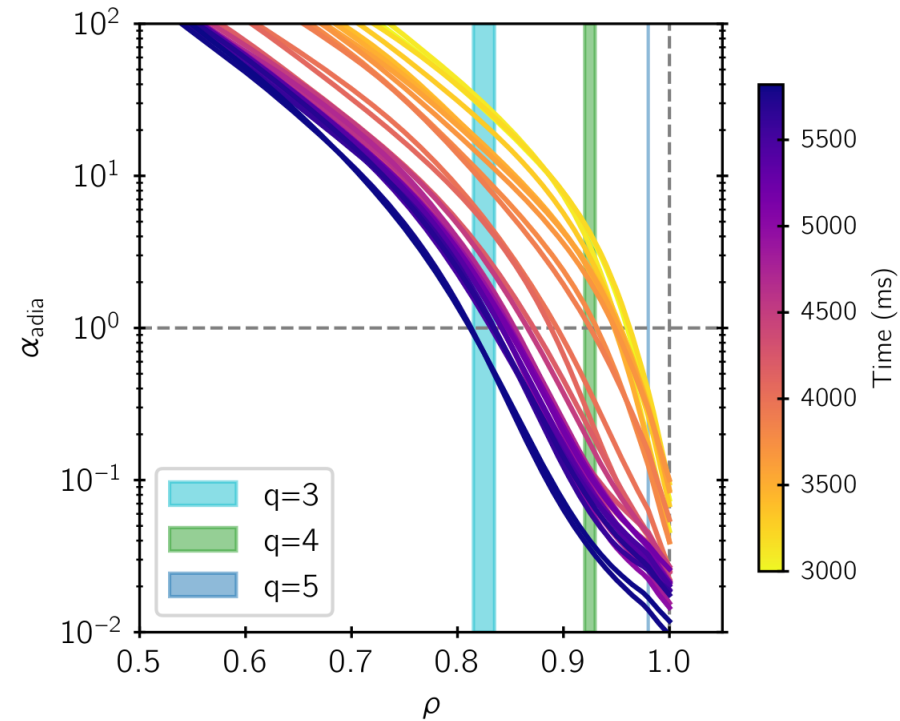
# Power Scalings

- Distinct power scalings of sep. and core density, over wide range
- No unique “Density Limit”
- $n_{sep} < n_G$  but steadily increasing with power
- Most cases don't terminate in disruption
- Radial departure from conventional wisdom



# Adiabaticity vs Time

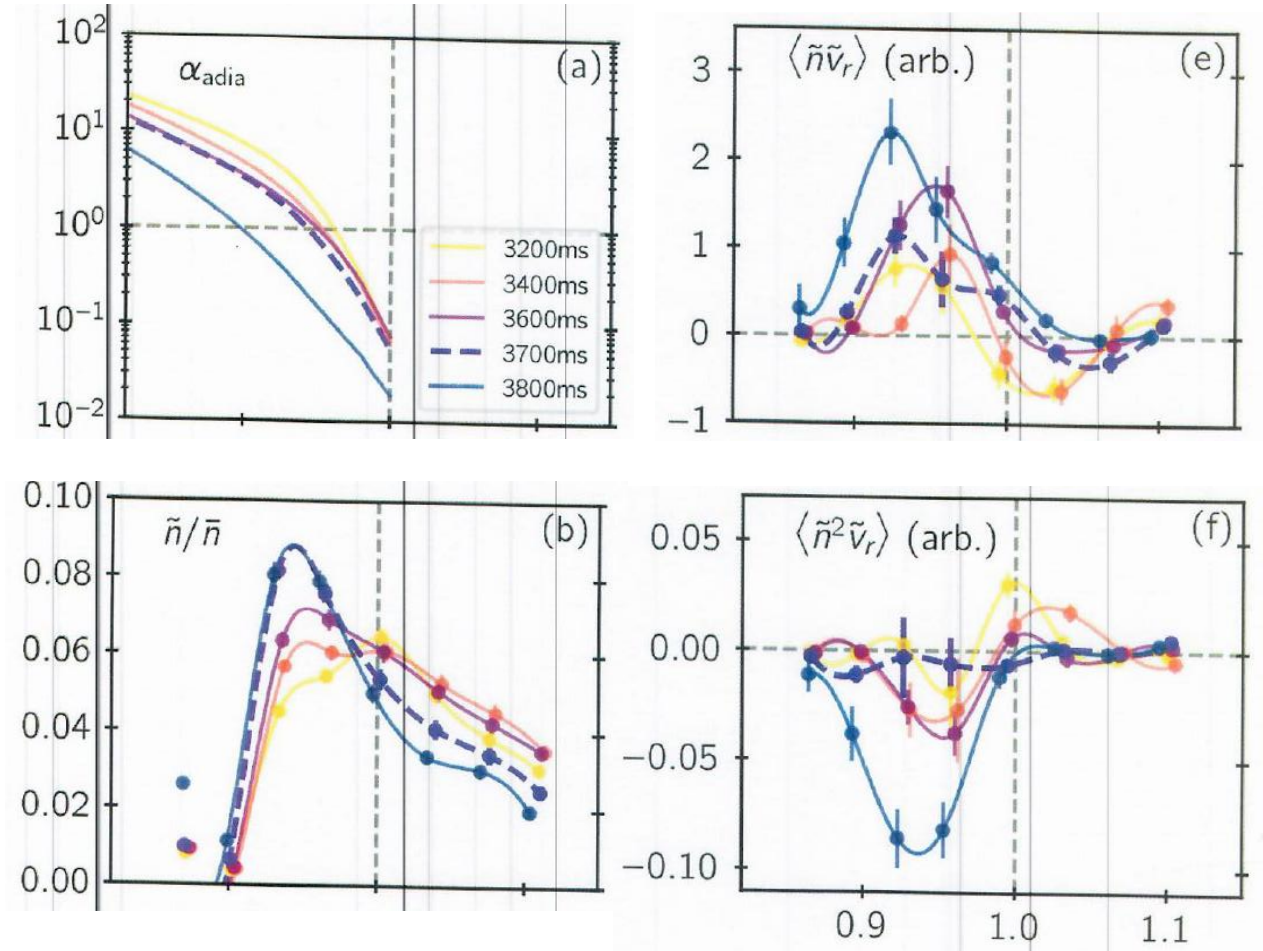
- $\alpha = k_{\parallel}^2 V_{th}^2 / \omega \nu$  drops, with low values penetrating inward with cooling
- Broad range of hydrodynamic turbulence develops after radiation



- $\alpha \leq 1$  at low  $q$  linked to disruption  
→ proximity  $q_{res}$  to  $q_{95}$  relevant

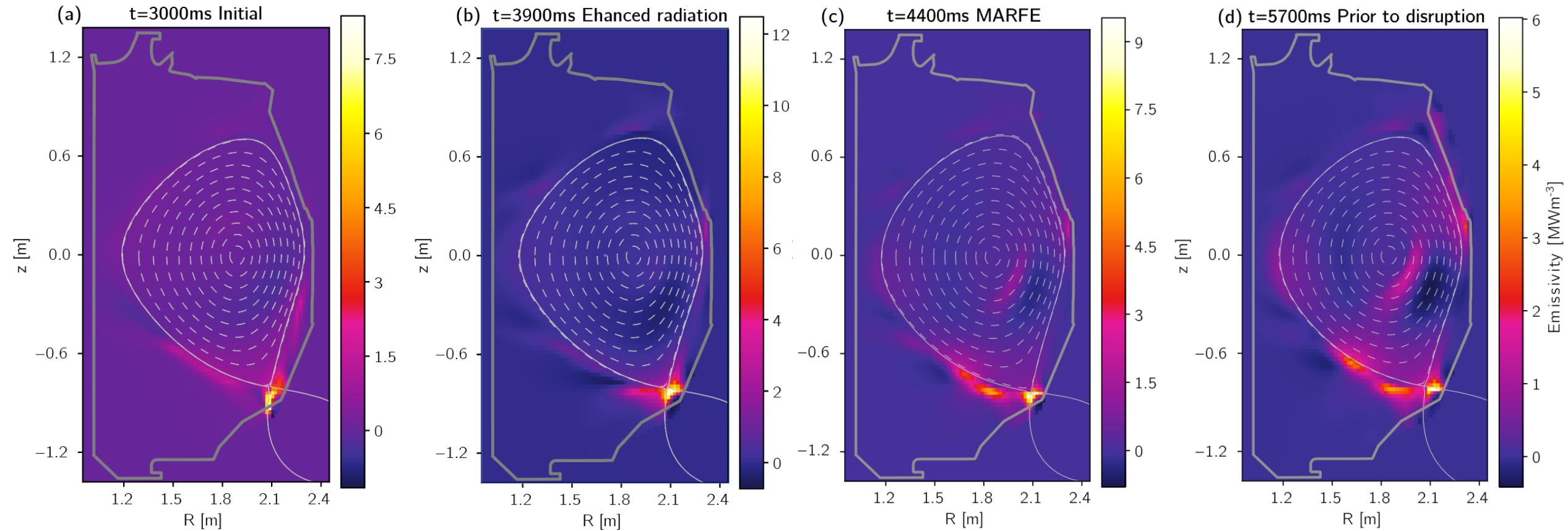
# $\alpha$ , Transport and Spreading

- With strong radiation:
  - $\alpha$  drops
  - $\tilde{n}/n$  increases
  - $\langle \tilde{v}_r \tilde{n} \rangle$  increases
  - spreading flux  $\langle \tilde{v}_r \tilde{n}^2 \rangle$  increases, spreading inward



- Suggests DL as Radiation  $\leftrightarrow$  Condensation  $\leftrightarrow$  Turbulent Transport Synergy

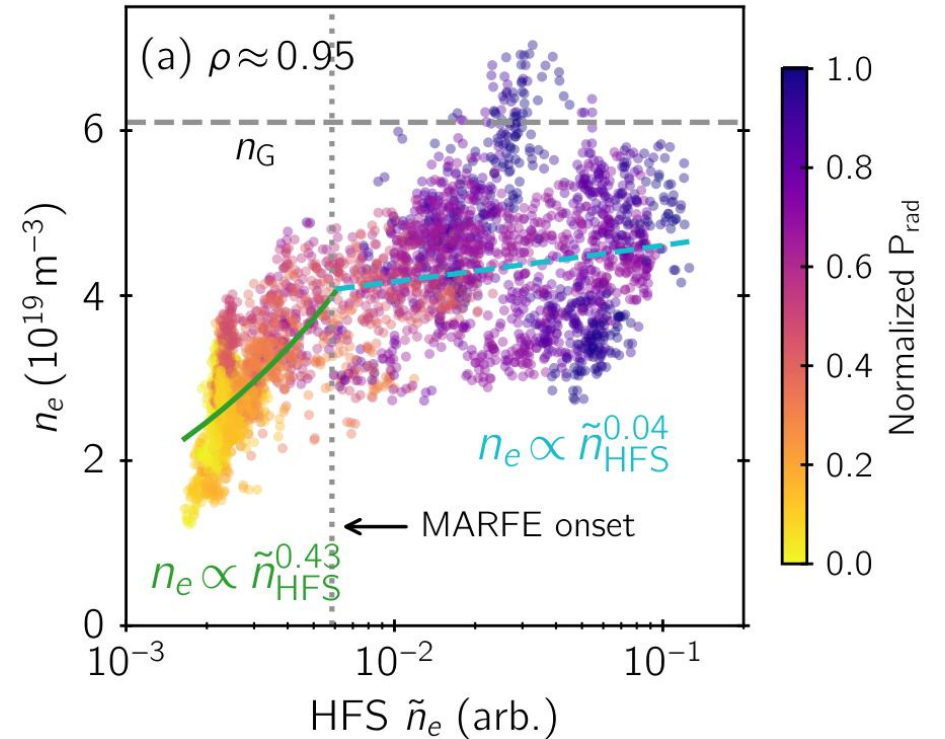
# Evolution of Radiation and MARFE



- Edge density limit linked to radiation
- Disruption follows strong MARFE

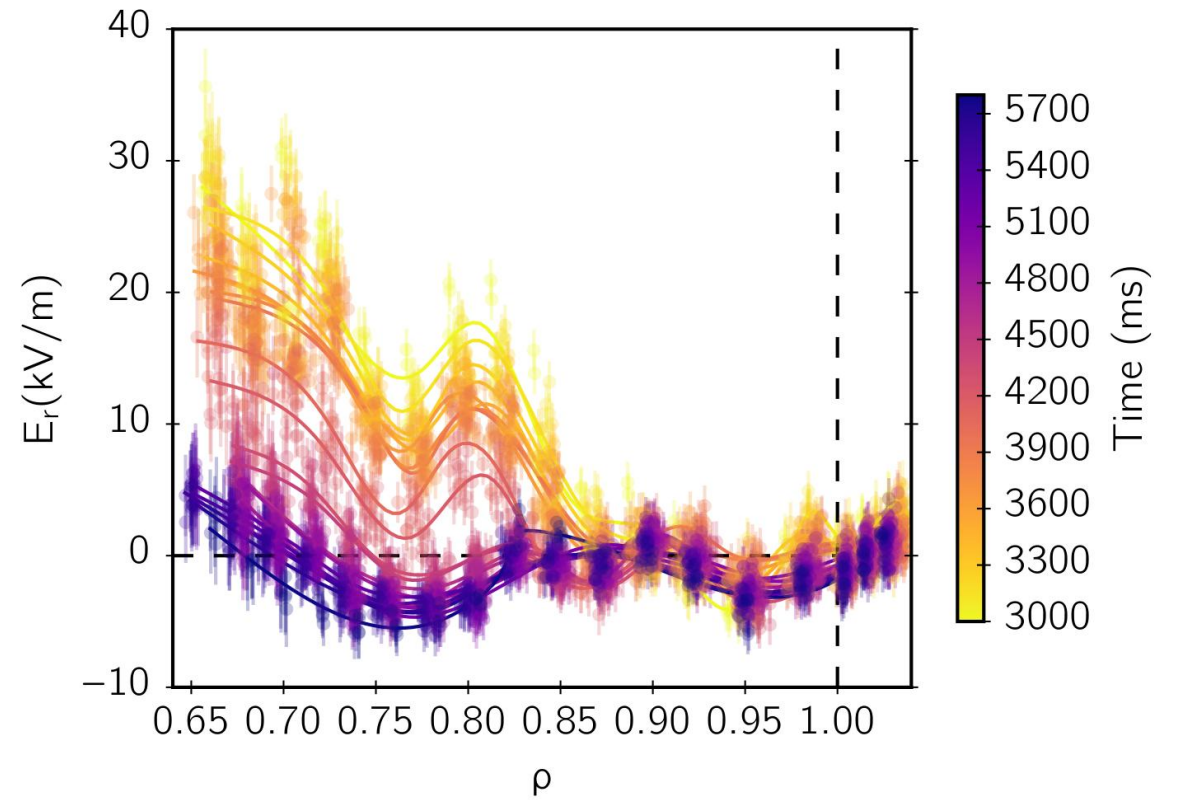
# Fluctuations and Density

- Edge density rises with  $\tilde{n}_{HFS}$  pre-MARFE
- Post MARFE edge density  $\sim$  “clamps”, manifesting a ‘limit’, as  $\tilde{n}_{HFS}$  increases.
- Broad range edge  $n_e$ , some exceeding  $n_G$ .
- Larger fluctuations and density saturation follow radiation onset.



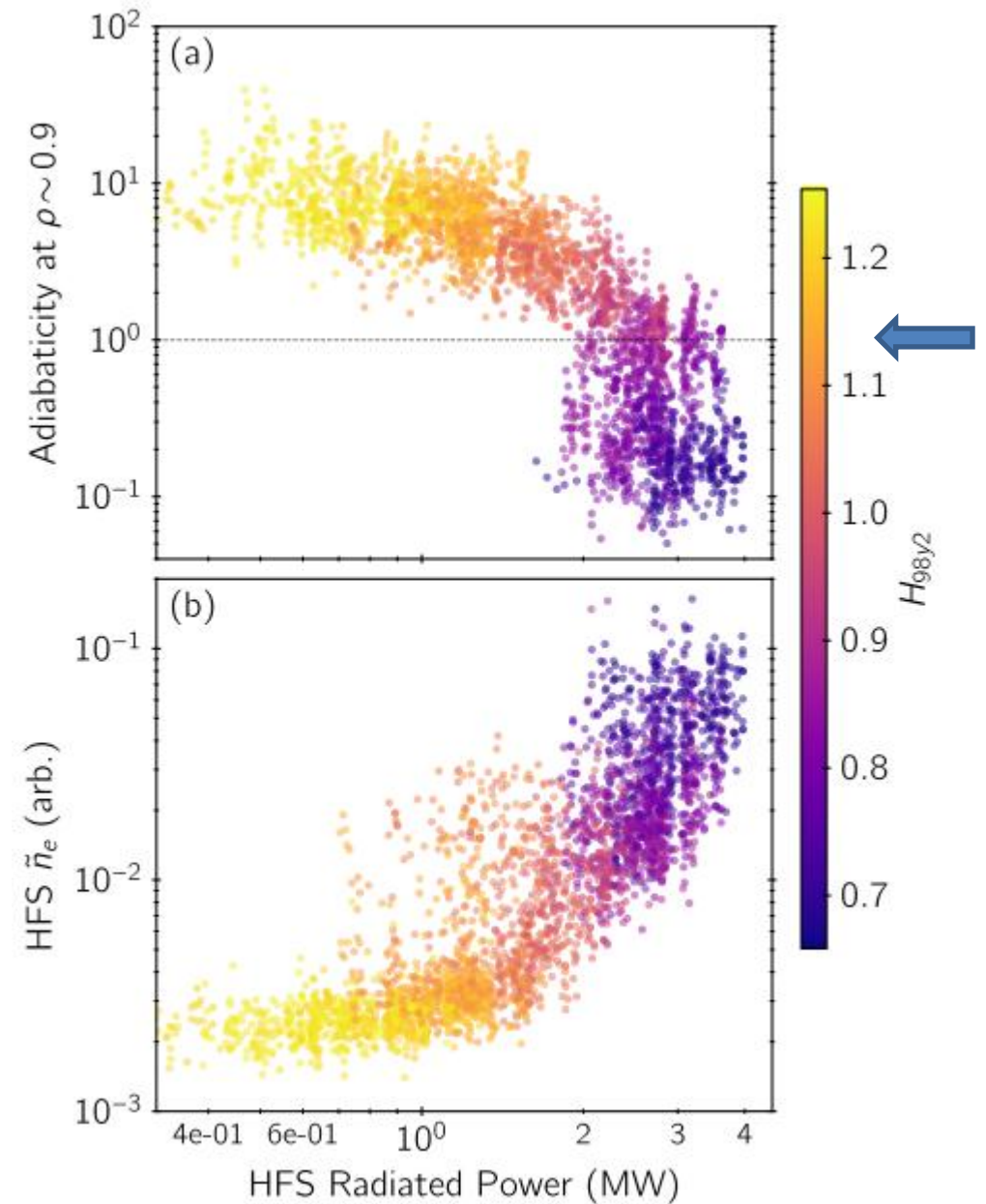
# Electric Field Shear Decays

Electric field shear for  $\rho > 0.85$   
decays with time



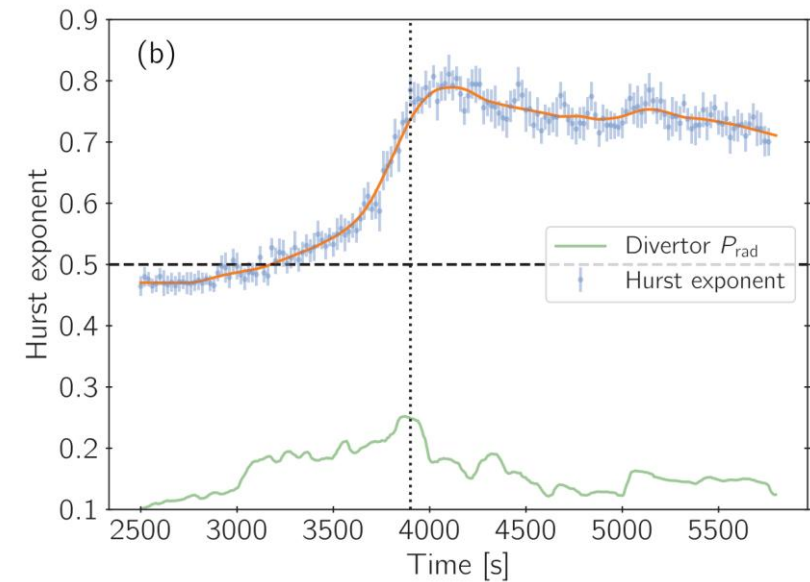
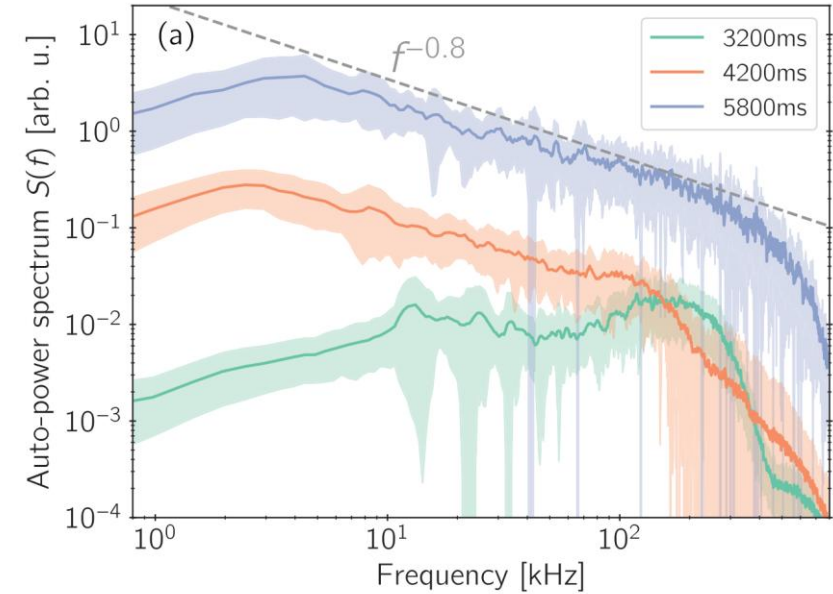
# $\alpha$ and turbulence level versus radiated power

- Rise in fluctuations tracks drop in  $\alpha < 1$  as HFS  $P_{rad}$  increases
- Radiation induces drop in adiabaticity to hydrodynamic regime:  $\omega\tau > 1$



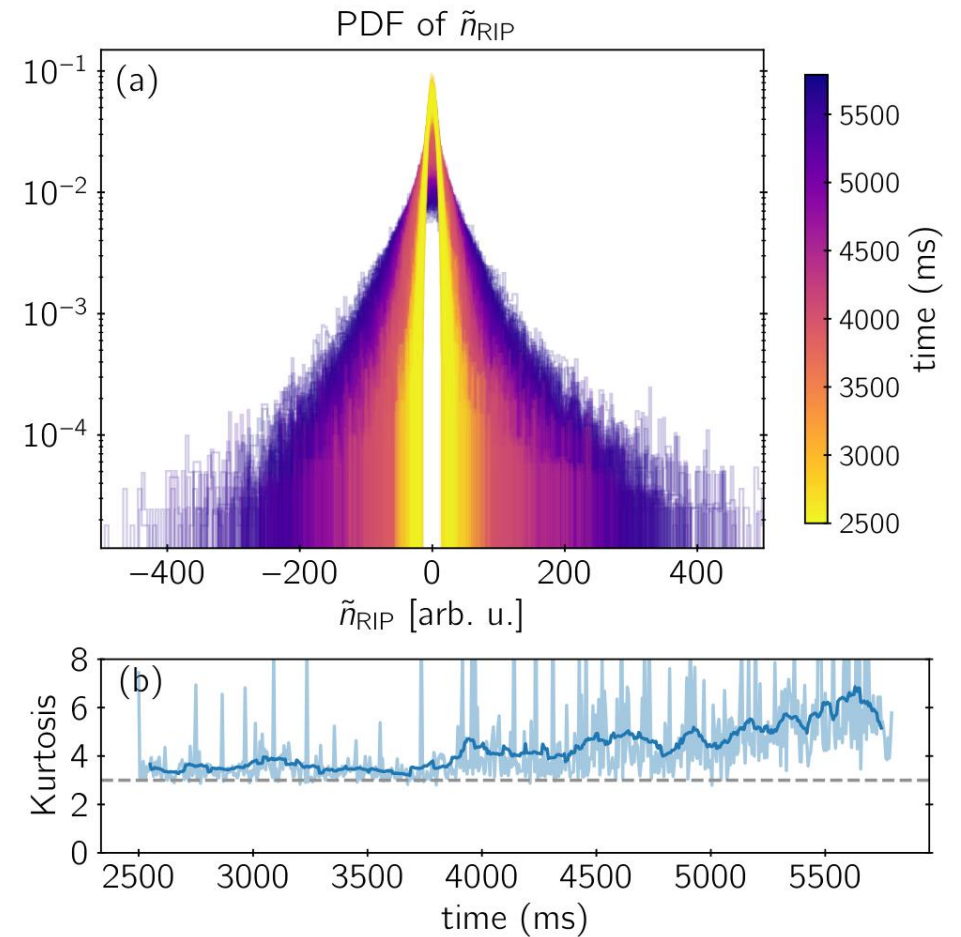
# Core Fluctuations

- Recall  $\sim$  independent core density limit
  - as  $n$  rises,  $S(f) \sim f^{-0.8}$
  - $P_{rad}$  Hurst exponent  $\rightarrow \sim 0.7$
  - Core DL  $\leftrightarrow$  Avalanching ?!



# Fluctuation PDF

- PDF of  $\tilde{n}$  exhibits 'fat tails' at longer time
  - Kurtosis rises
- Core density limit enforced by density  
avalanching ?!



# What have we learned?

- No single “DENSITY LIMIT”, but rather an edge and core  $n$  saturation, with different mechanisms
- Power dependence unambiguous
- Soft limit  $\rightarrow$  most plasmas don't disrupt
- Suggests evolution:

Radiative condensation MARFE  $\rightarrow$  edge cooling  $\rightarrow \alpha < 1$  hydro-regime  $\rightarrow$

enhanced transport with shear layer collapse  $\rightarrow$  DL via particle outflow

i.e. Radiative cooling releases strong particle transport  $\rightarrow$  ‘density clamp’

# More Theory

- A Work in Progress

# What is Needed?

- Model of radiative condensation/MARFE in turbulent medium
- Radiative condensation: Thermal (cooling) instability s/t  $\omega < k_{\parallel} c_s$  so  $\delta P = 0$

$$\gamma = \frac{2}{5n} \left( \frac{2L}{T} - \frac{\partial L}{\partial T} \right) - \chi_{\perp} k_{\perp}^2 - \chi_{\parallel} k_{\parallel}^2 \quad L \sim n n_z f(\tau_c)$$

- c.f. G. Field .. 1965 → Drake 1987 (linear theory in cylinder)



infinity of papers on ISM c.f. Balbus, 1995 for tutorial

- R.C. plus turbulence intensively studied in ISM cf Max Gronke+ MNRAS 2021

# Strategy

- Incorporate radiative cooling into reduced model
- Defining competition for power scaling will be Heat Flux vs Turbulent Transport + Cooling
  - $\alpha$  sets branching ratio
  - Coupling: cooling  $\uparrow \rightarrow \alpha \downarrow \rightarrow$  transport  $\uparrow$   
Ratio [R.C. / Transport] of interest
- Minimal Model:
  - Fluctuation energy
  - Zonal energy
  - $T_e \rightarrow$  high  $n$ ,  $T_e \sim T_i$ ;  $D \sim \chi_e$  for electrostatics

# Proto-Model

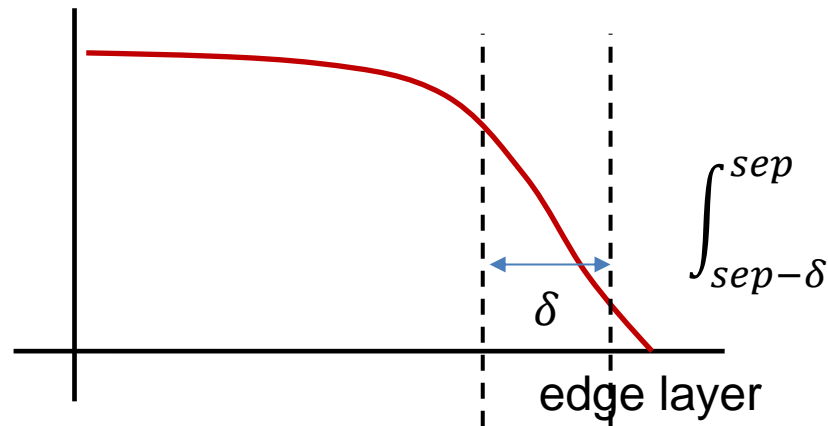
- Temperature Equation – Mean field,  $k_{\parallel} = 0$

$$n \frac{\partial T}{\partial t} = \nabla_r n \chi(\varepsilon) \nabla_r T - L$$

radiative loss ( $L > 0$ )  
 $L = L(n, T) \approx n^2 \Lambda$

turbulent transport  
 $\chi(\varepsilon)$  - fluctuation dependent

- So



$$\langle n \frac{\partial T}{\partial t} \rangle = n \chi(\varepsilon) \nabla_r T \Big|_{sep} - n \chi(\varepsilon) \nabla_r T \Big|_{sep-\delta} - \int_{sep-\delta}^{sep} L$$

# Proto-Model, cont'd

- But:  $-n \chi(\varepsilon) \nabla_r T|_{sep-\delta} = Q \rightarrow$  heat flux from core

So for edge layer:

$$\left\langle n \frac{\partial T}{\partial t} \right\rangle = Q_{core} - (\delta)L + n \chi(\varepsilon) \nabla_r T \Big|_{sep}$$

radiative losses

transport loss at sep.  $\nabla_r T < 0$

and can simplify to:

$$\frac{\partial T}{\partial t} = -\frac{\chi(\varepsilon)T}{\Delta^2} + q - \frac{L(n, T)}{n}$$

becomes a simple mean field temperature equation

$\Delta \equiv$  layer width

# Key Ratio

- Physics of Power Scaling of Great Interest
- From  $T$  eqn, radiation and transport compete for  $q$ , so

$$\frac{L/n}{\chi(\varepsilon)T/\Delta^2} \sim \frac{\gamma_{RC}}{\chi(\varepsilon)/\Delta^2} \sim D_{a,R} (Q, n, \dots)$$

$D_a \sim \tau_{turb}/\tau_{react} \rightarrow$  Damkohler #, from combustion

$D_a \gg 1 \rightarrow$  reaction time short  $\rightarrow$  flame sheets

$D_a \ll 1 \rightarrow$  mixing time short  $\rightarrow$  pre-mixed flame

Radiative Damkohler #  
(after Gronke)

- Expect  $D_{a,R}$  first
  - rises to  $\geq 1$  as R.C. grows
  - drops to  $\leq 1$  as edge plasma cools
  - $\therefore$  time history of interest  $\rightarrow$  drop  $D_{a,R} \rightarrow$  density clamp
  - Is  $D_{a,R} \sim 1$  indicative of edge turbulence levels at high density ?

# Minimal Model → ala' Singh + P.D.


$$\frac{\partial T}{\partial t} = -\frac{\chi \varepsilon T}{\Delta^2} + q - \frac{L(n, T)}{n} \quad \rightarrow \text{temp}$$

$$\frac{\partial \varepsilon}{\partial t} = a_1 \gamma(n, T) \varepsilon - a_4 V_z^2 \varepsilon f(\alpha) - a_2 \varepsilon^2 \quad \rightarrow \text{fluctuation energy}$$

$$\frac{\partial V_z^2}{\partial t} = a_4 V_z^2 \varepsilon f(\alpha) - b_3 n V_z^2 \quad \rightarrow \text{zonal flow}$$

→  $f(\alpha) \equiv$  adiabaticity “switch” -  $T$  evolves,  $\alpha$  follows

$\alpha \sim k_{\parallel}^2 v_{th}^2 / \omega \nu \sim T^2$  accounts for Z.F. decay,  
↔ production drop


$$f(\alpha) \sim 1, \alpha > 1$$
$$f(\alpha) \ll 1, \alpha \ll 1$$

$f(\alpha) \ll 1 \rightarrow$  Z.F. collapse

# Next Steps

- Exercise Model → Prediction!

Edge DL defined by  $\gamma_{rad}(n)$  sufficient to induce  $\alpha < 1$  !?  
 $n$  limit to induce strong particle transport

- 1D version → cooling fronts !?

→ cooling + transport fronts ?!

- Profile structure, magnetic geometry?

N.B.  $\chi_{\parallel} k_{\parallel}^2 \sim (D_T \chi_{\parallel})^{1/2} k_{\theta} / L_s$  → higher m modes damped by  
conduction + turbulence interaction  
 $D_r / \Delta^2 \sim \chi_{\parallel} k_{\parallel}'^2 \Delta^2$  → supports mean field theory approach

- Can radiative condensation couple to turbulence dynamics directly??

# Partial Conclusion

- “Causality” counter-intuitive

Radiation/MARFE  $\rightarrow$  cooling  $\rightarrow$  strong transport

$$\alpha > 1 \rightarrow \alpha < 1$$

Opposite to conventional wisdom!

- Non-disruptive termination
- Two channels for power scaling, strongly coupled.  $D_{a,R}$  is useful  $\rightarrow$  experimental analysis !?

$$D_{a,R} \sim \gamma_{rad} / \frac{\chi(\varepsilon)}{\Delta^2}$$

- Model extended to encompass radiative condensation

# From then to now of DL

- Greenwald (1988) scaling → Incomplete → Power Scaling, Multiple Limits

$$\bar{n} \sim I_p / \pi a^2$$

- Sudo (1990) Scaling (stellarator) → entering mainstream → G + S unification

$$\bar{n} \sim P^{1/2} \quad \text{Transport Physics Central}$$

- DL as MHD + Disruption phenomenon → frequently, even usually, not the case  
Rebut, Gates, White ?? (revision needed)
- DL as a 'Back-Transition' → from heresy to convention
- Radiation triggers MHD → Rad. triggers transport...
- Better Density 'Saturation' than 'Limit' !

# Remaining Issue: "Causality"

- Demonstrated link: Radiation  $\leftrightarrow$  Transport
- Which is trigger?
- RI mode – reanimated for SPARC (via Temasek fund)

How avoid MARFE's  $\rightarrow$  shift plasma off inner wall

$\therefore$

$n \leq n_{DL}$  and shift plasma toward inner wall

- $\rightarrow$  induce MARFE and increased particle transport, i.e. hit DL?

**Back-Up**

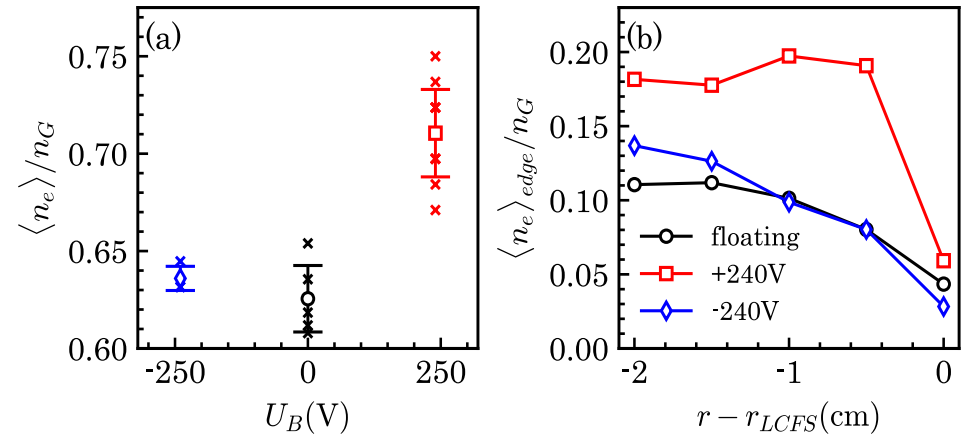
# The Obvious Question

- Can driving the shear layer sustain high densities, where  $L \rightarrow DL$ , otherwise ?
- “Driving”  $\longrightarrow$  bias electrode – here (J-TEXT). Not a conventional H-mode
- Long history of bias-driven shear layers in  $L \rightarrow H$  saga – R.J. Taylor, et. seq.
- Recent: Shesterikov, Xu et. al. 2013 - Textor
- Electrode  $\rightarrow J_r \rightarrow V_\theta \rightarrow V'_E$  etc.
- New Here?
  - High Density
  - Gas Puffing  $\rightarrow$  push on DL
  - Analysis

c.f. Rui Ke, P.D. + NF 2022

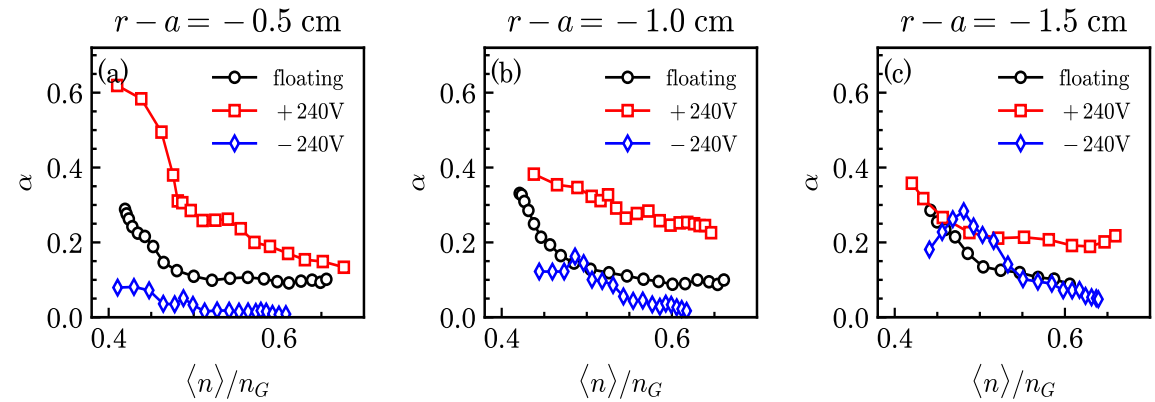
# The Answer – Looks Promising!

- Edge density doubled for +240V bias
- $\bar{n}_{\text{max,bias}} > \bar{n}_{\text{max,float}}$
- Note:  $\bar{n}_{\text{max,float}} \sim 0.7n_G$



Experiment limited by graphite probe sputtering

- Key parameter?
  - $\alpha$  systematically higher with +bias
  - $\alpha \sim T^2/n$   $\leftarrow$  Reduced transport  $\rightarrow$  higher T



- Turbulence spreading quenched by positive bias

# The Physics

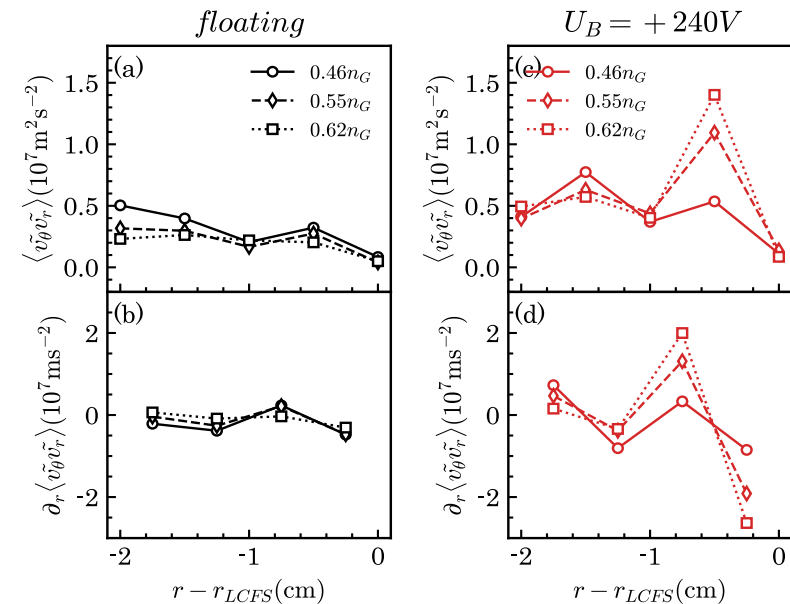
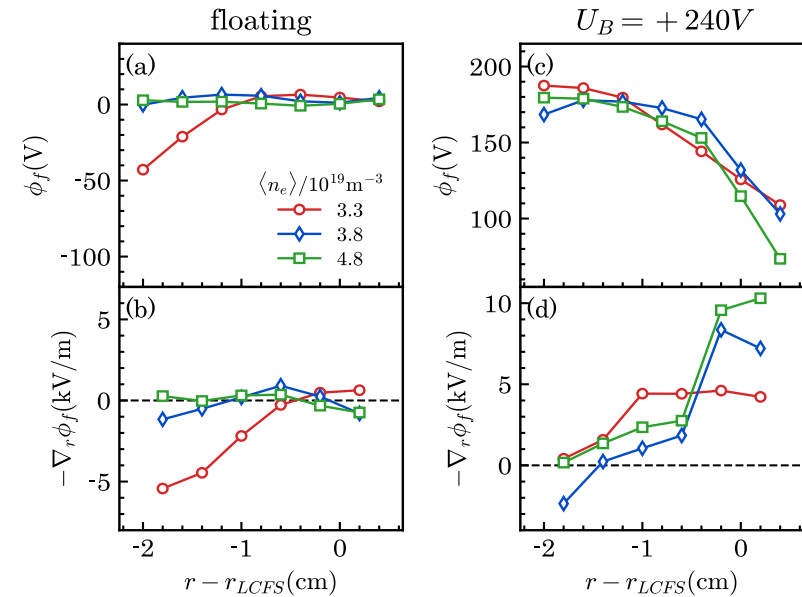
- Edge Shear Layer produced for +bias

N.B. Not an  $E_r$  well

- Reynolds stress, force increase for +bias

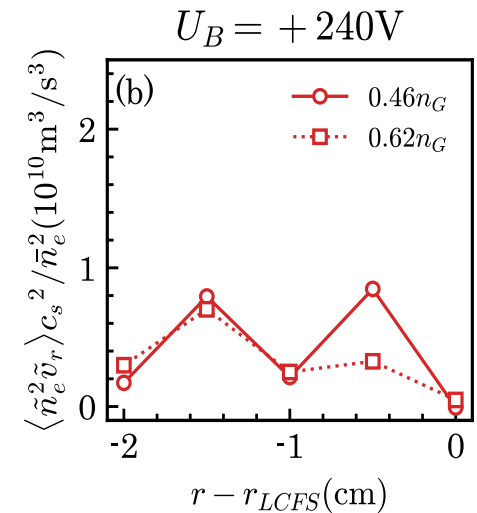
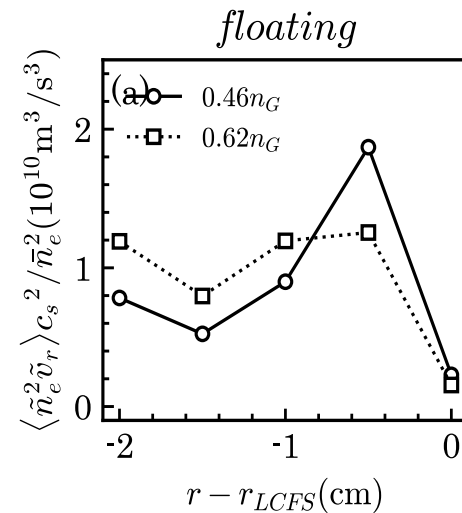
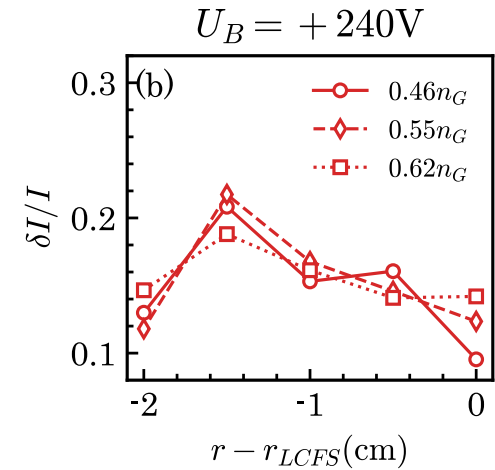
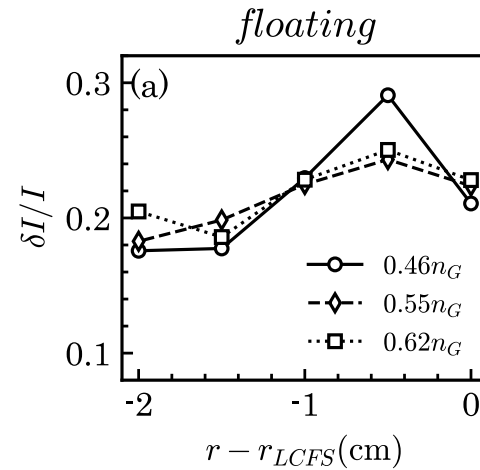
$\leftrightarrow$  bias effect on eddy alignment

“Shearing”  $\leftrightarrow$  interplay of bias and Reynolds stress



# The Physics

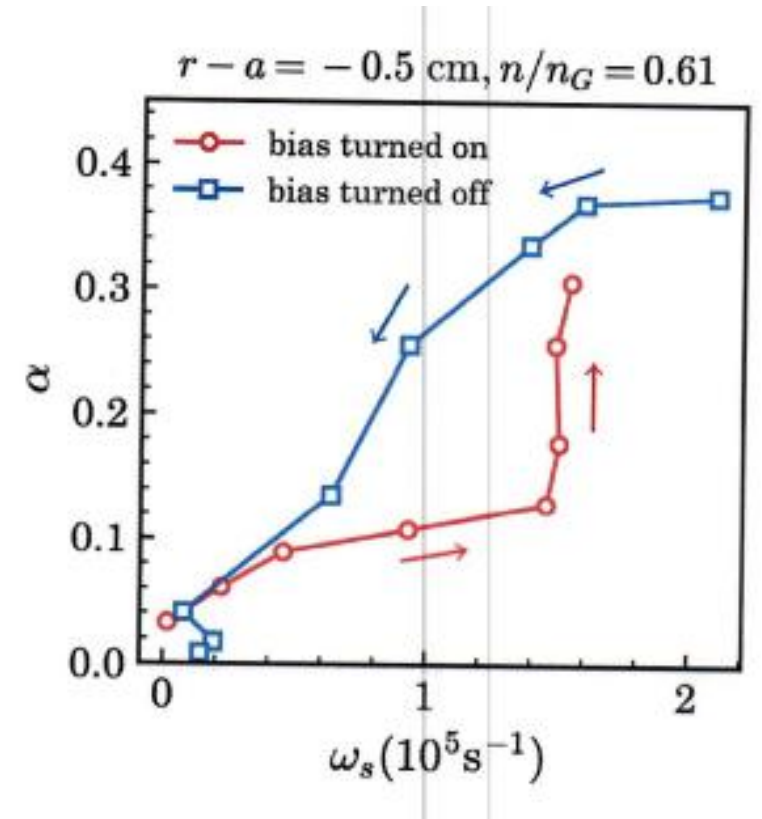
- $\delta I/I$  ( $\rightarrow \tilde{n}/n$ ) fluctuations sharply reduced by +bias



- Turbulence spreading quenched by +bias

# Key Parameter vs Control Parameters

- $\alpha$  vs  $\omega_{shear}$  exhibits hysteresis loop
- Cntr clockwise rotation  $\rightarrow \omega_{shear}$  'leads'  $\alpha$
- Is  $\alpha$  unique 'key parameter'?
- For drift waves,  $\alpha \sim T^2/n$ 
  - $\rightarrow$  shear  $\uparrow \rightarrow$  turbulence  $\downarrow \rightarrow$  heat transport  $\downarrow$
  - $\rightarrow \alpha$  increases
- Is  $\omega_{shear}$  the control parameter?



# Ongoing and Future Work

- Bias experiment with improved probe
- Ip scan vs  $n/n_G$  scan ? – obvious ‘Greenwald test’ (Long+ 2024, N.F.):

Ip ramp down explained via  $\omega_{shear} \tau_{cor}$

- Physics of spreading (Long, PD+ 2024)
  - Spreading  $\leftrightarrow$  Blob emission
  - Broken symmetry: “Spreading” dominated by large blobs
  - No apparent correlation of spreading and particle flux

**Thank You !**

**Supported by U.S. D.O.E.**