

Cross Phase Evolution and Its Role in ELM Bursts

P.H. Diamond

**WCI Center for Fusion Theory, NFRI, Korea
CMTFO and CASS, UCSD, USA**

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Collaborators

- Pengwei Xi – PKU, LLNL

X.-Q. Xu – LLNL

See: P.W. Xi, Xu, P.D. – PRL, PoP 2014

- Related work:

T. Rhee – NFRI

J.M. Kwon – NFRI

W.W. Xiao - UCSD, NFRI, SWIP

See: T. Rhee, et. al. PoP 2012

Acknowledgements

- Xiao-Lan Zou, R. Singh, Todd Evans, Rajesh Mainji, Sang-Hee Hahn, Michael Leconte, Guilhem Dif-Pradalier

Caveat Emptor

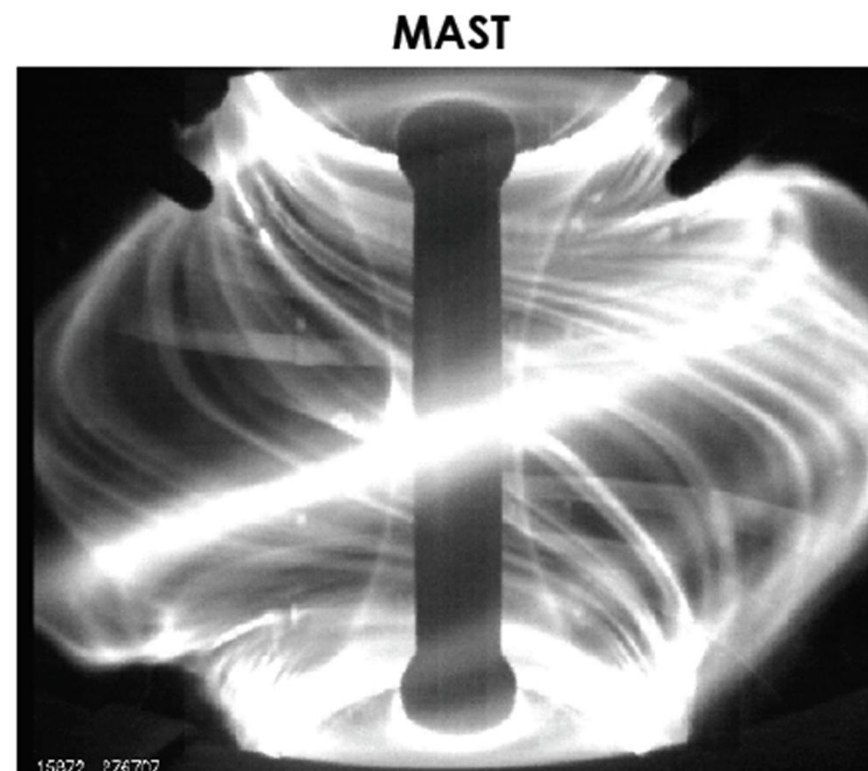
- Not a professional ELM-ologist
- Perspective is theoretical, and focus is on issues in **understanding dynamics**
- Perspective is that of a transport theorist
- Aim is to distill elements critical to model building
- Unresolved issues are discussed

Outline

- ELMs
 - Conventional wisdom: A Quick Look
 - Some physics questions
- Recent Progress:
 - cross phase coherence and the origin of bursts
 - phase coherence as leverage for ELM mitigation
 - a deeper – but incomplete – look at phase dynamics
- Conclusions and Discussion

Terra Firma: Conventional Wisdom of ELMs

- ELMs are ~ quasi-periodic relaxation events occurring at edge pedestal in H-mode plasma
- ELMs
 - Limit edge pedestal –
 - Expel impurities +
 - Damage PFC –
- ELMs → a serious concern for ITER
 - $\Delta W_{ELM} \sim 20\% W_{ped} \sim 20 \text{ MJ}$
 - $W_{ELM} / A \sim 10 \times \text{limit for damage}$
 - $\tau_{rise} \sim 200 \mu\text{sec}$



Terra Firma: Conventional Wisdom of ELMs

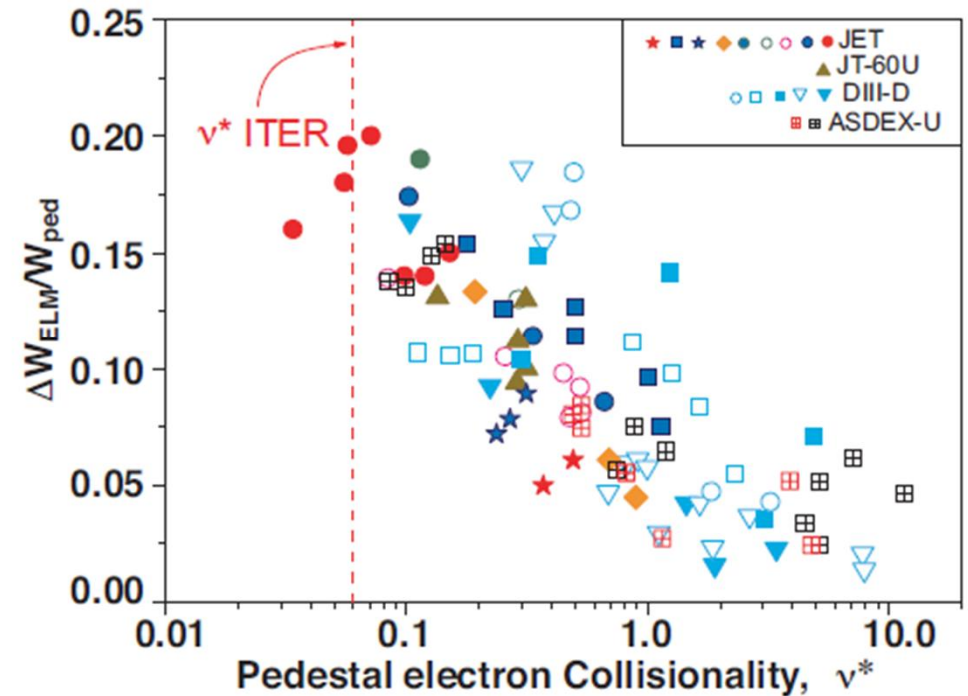
- ELM Types

- I, II: $\omega_{ELM} \uparrow$ as $P \uparrow$, greatest concern, related to ideal stability
- III: $\omega_{ELM} \downarrow$ as $P \uparrow$, closer to P_{Th} , unknown \rightarrow resistive ??

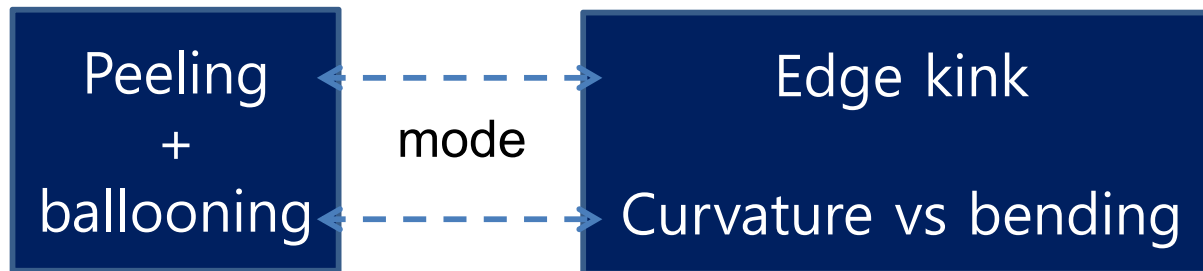
- Physics

- Type I, II ELM onset \rightarrow ideal stability limit
- i.e. peeling + ballooning

Relative ELM Energy Loss



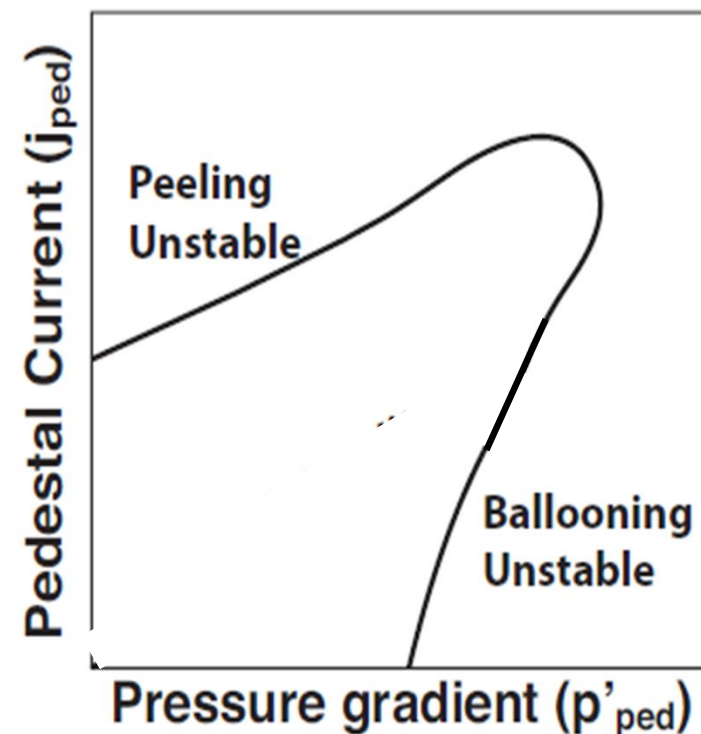
[A. Loarte, *Plasma Phys. Control. Fusion* 45 1549 (2003)]



δW
+
Pedestal, geometry

Terra Firma: Conventional Wisdom of ELMs

- Some relation of ELM character to collisionality is observed
 - Low collisionality \rightarrow peeling \sim “more conductive”
 - High collisionality \rightarrow ballooning \sim “more convective”
- Many basic features of ELMs consistent with ideal MHD peeling-ballooning theory
- Pedestal perturbation structure resembles P-B eigen-function structure (?!)



Some Physics Questions

- What IS the ELM? Why is the ELM?
 - ELMs single helicity or multi-helicity phenomena?
Relaxation event \leftrightarrow pedestal avalanche?, turbulence spreading?
 - How and why do actual **bursts** occur?
Why doesn't turbulence force $\nabla P \sim \nabla P_{crit}$ oscillations?
 - Pedestal turbulence develops during ELM. Thus, how do P-B modes interact with turbulence? – either ambient or as part of MH interaction?
 - Does, or even should, the linear instability boundary define the actual ELM threshold?

I) Basic Notions of ELMs:

ELM Bursts and Thresholds as

Consequence of Stochastic Phase Dynamics

Simulation model and equilibrium in BOUT++

- 3-field model for nonlinear ELM simulations
 - ✓ Including essential physics for the onset of ELMs

- Peeling-ballooning instability
- Resistivity
- Hyper-resistivity
- Ion diamagnetic effect

$$\frac{d\varpi}{dt} = B \nabla_{\parallel} J_{\parallel} + 2\mathbf{b}_0 \times \boldsymbol{\kappa} \cdot \nabla \tilde{P} + \mu_{i,\parallel} \partial_{\parallel}^2 \varpi$$

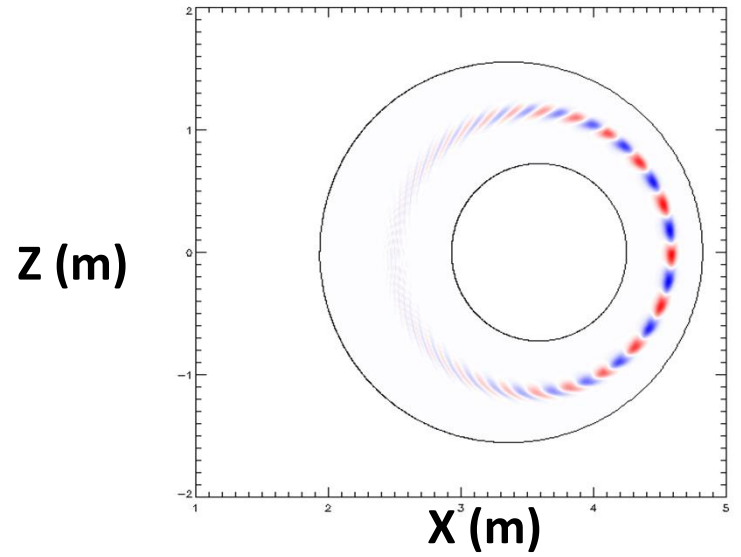
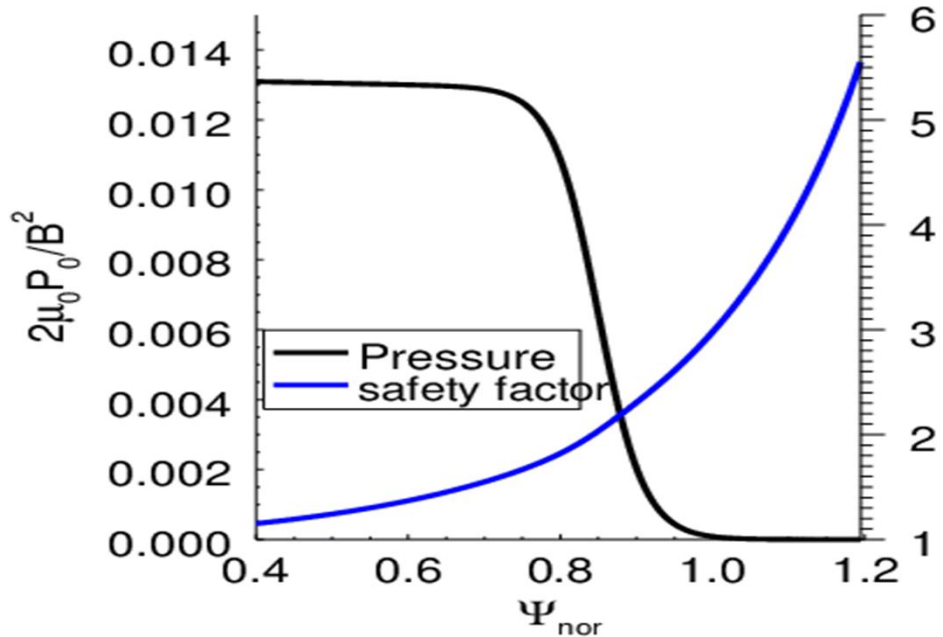
$$\frac{d\tilde{P}}{dt} + \mathbf{V}_E \cdot \nabla P_0 = 0$$

$$\frac{\partial A_{\parallel}}{\partial t} + \partial_{\parallel} \phi_T = \frac{\eta}{\mu_0} \nabla_{\perp}^2 A_{\parallel} + \frac{\eta_H}{\mu_0} \nabla_{\perp}^4 A_{\parallel}$$

$$\varpi = \frac{m_i n_0}{B} \left(\nabla_{\perp}^2 \phi + \frac{1}{en_0} \nabla_{\perp}^2 \tilde{P}_i \right)$$

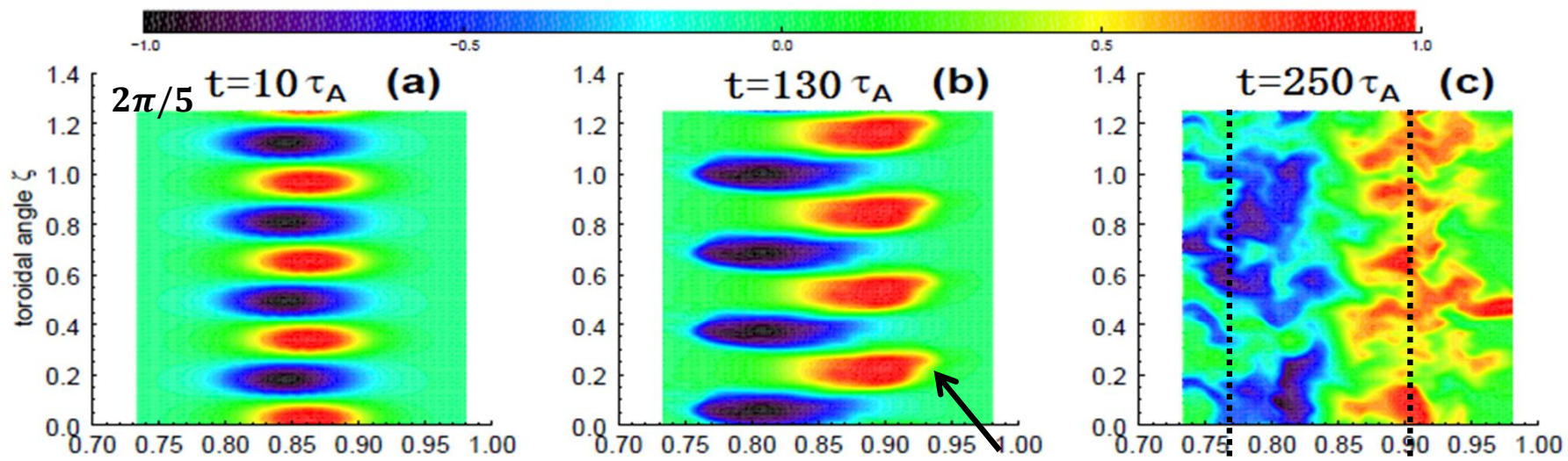
hyper resistivity

$$d/dt = \partial/\partial t + \mathbf{V}_{ET} \cdot \nabla, \mathbf{V}_{ET} = \frac{1}{R} \mathbf{b}_0 \times \nabla \phi_T, \phi_T = \phi_0 + \phi, \nabla_{\parallel} f = B \partial_{\parallel} \frac{f}{R}, \partial_{\parallel} = \partial_{\parallel}^0 + \delta \mathbf{b} \cdot \nabla, \delta \mathbf{b} = \frac{1}{B} \nabla A_{\parallel} \times \mathbf{b}_0, J_{\parallel} = J_{\parallel 0} + \tilde{J}_{\parallel}, \tilde{J}_{\parallel} = -\nabla_{\perp}^2 A_{\parallel} / \mu_0$$



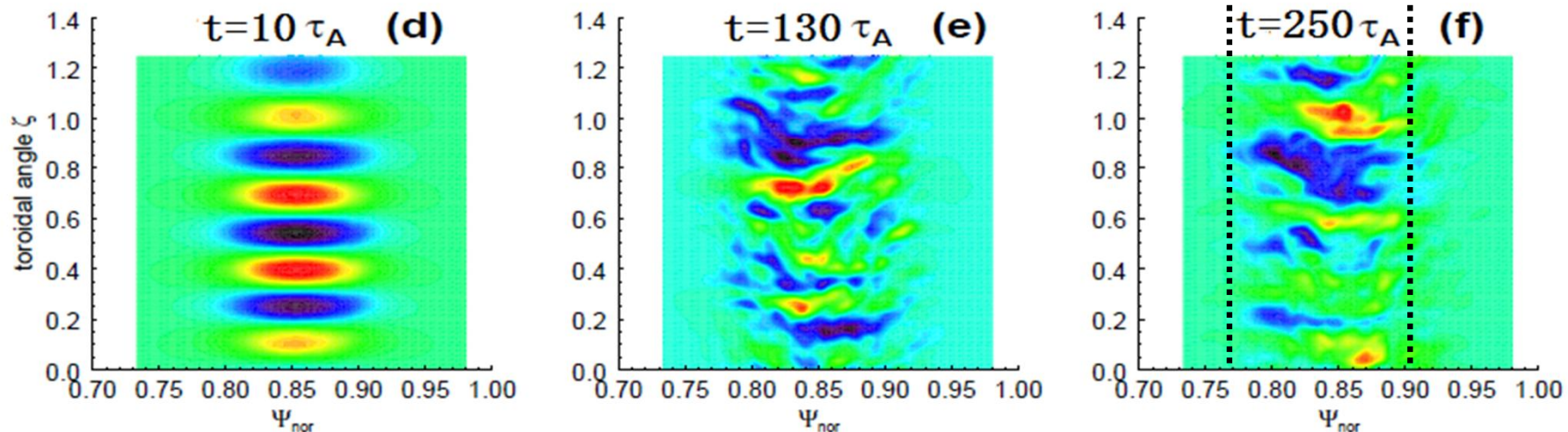
Contrast of perturbation evolution (1/5 of the torus)

Single Mode



Filaments

Multiple Mode



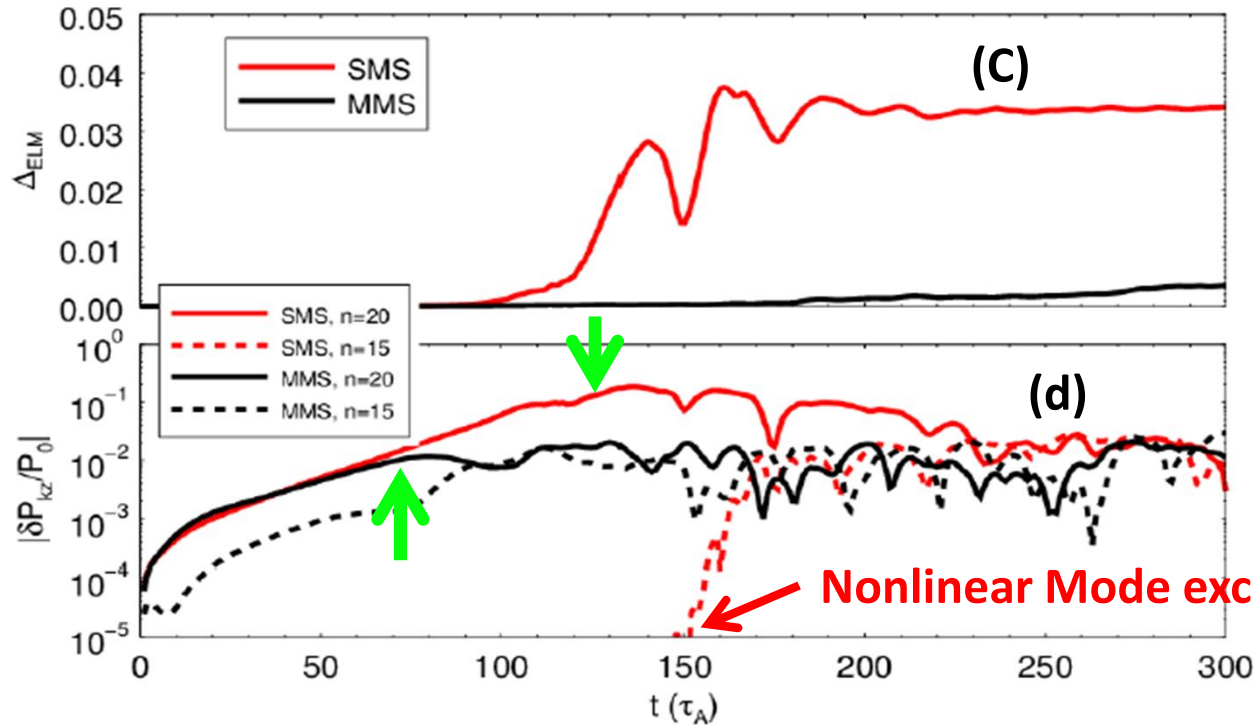
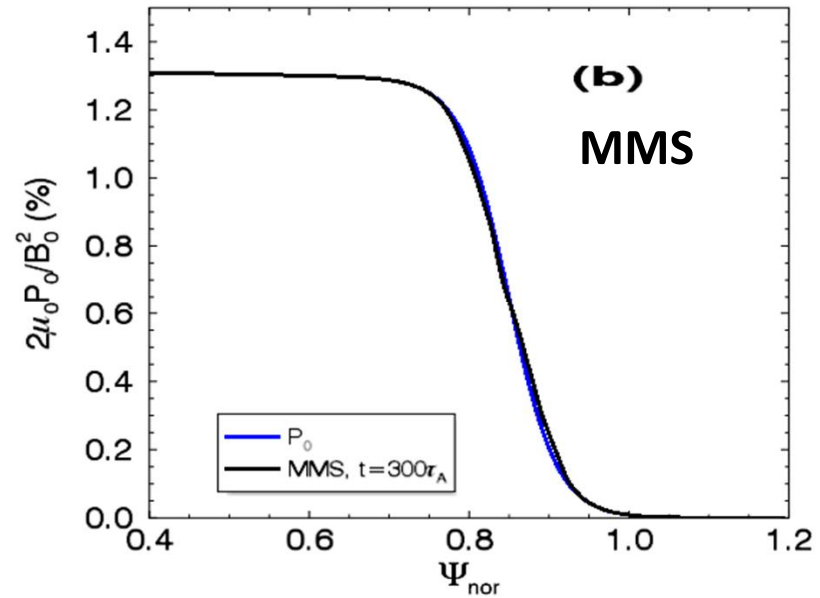
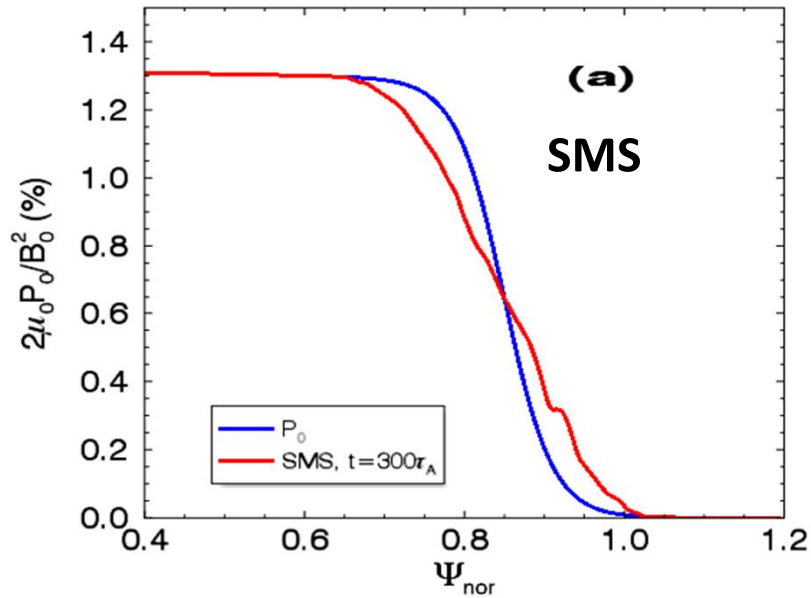
Linear phase

Early nonlinear phase

Late nonlinear phase

- Single mode: Filamentary structure is generated by linear instability;
- Multiple modes: Linear mode structure is disrupted by nonlinear mode interaction and no filamentary structure appears

Single mode: ELM crash || Multiple modes: P-B turbulence



- ELM size larger for SMS

$$\Delta_{ELM} = \frac{\Delta W_{ped}}{W_{ped}} = \frac{\int dx^3 (P_0 - \langle P \rangle_\zeta)}{\int dx^3 P_0}$$

- SMS has longer duration linear phase than MMS


Relative **Phase** (Cross Phase) **Dynamics** and Peeling-Ballooning Amplification

Peeling-Ballooning Perturbation Amplification is set by Coherence of Cross-Phase

i.e. schematic P.B. energy equation:

$$\frac{\partial}{\partial t} E_k = \langle \tilde{\phi} 2\hat{b}_0 \times \vec{k} \cdot \nabla \tilde{P} \rangle_{\vec{k}} \longleftarrow \sim \langle \tilde{v}_r \tilde{P} \rangle \rightarrow \begin{array}{l} \text{energy release from } \nabla \langle P \rangle \\ \rightarrow \text{quadratic} \end{array}$$

$$+ \sum_{\vec{k}', \vec{k}''} \tau_{c\vec{k}} C(\vec{k}', \vec{k}'') E_{\vec{k}'} E_{\vec{k}''} - \sum_{\vec{k}'} \tau_{c\vec{k}+\vec{k}'} C(\vec{k}', \vec{k}) E_{\vec{k}'} E_{\vec{k}} - \text{dissipation}$$

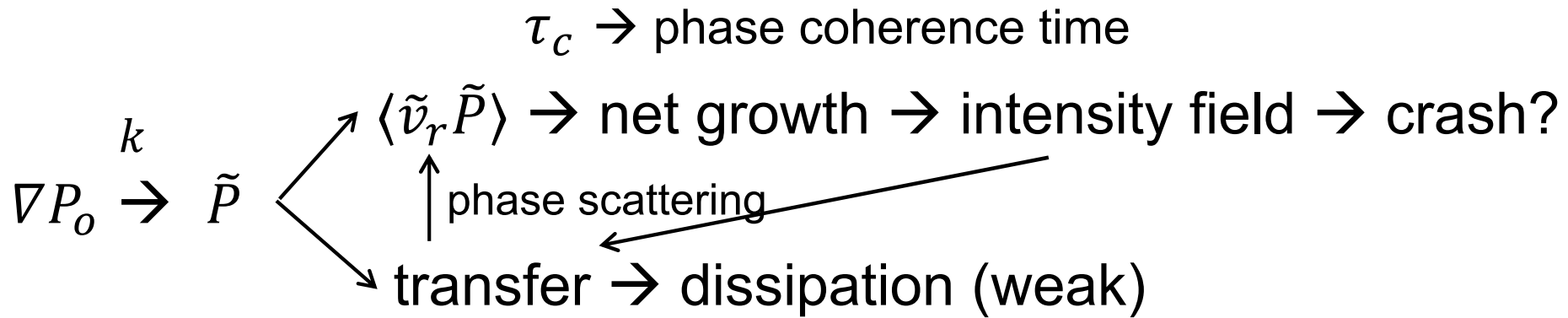


 nonlinear mode-mode coupling → quartic

NL effects

- energy couplings to transfer energy (weak)
- response scattering to de-correlate $\tilde{\phi}$, $\tilde{P} \rightarrow$ regulate drive

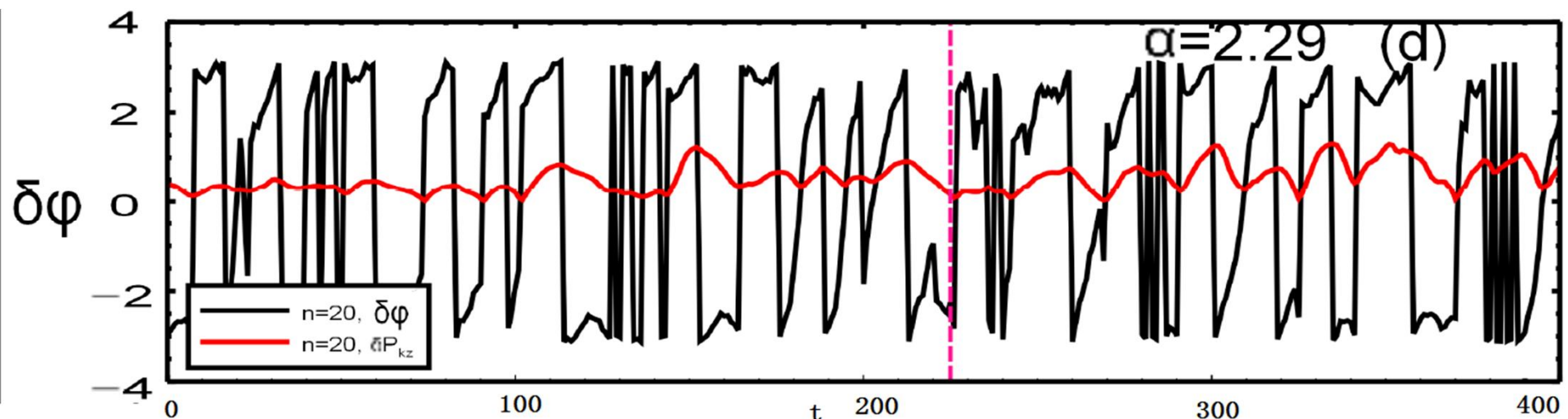
Growth Regulated by Phase Scattering



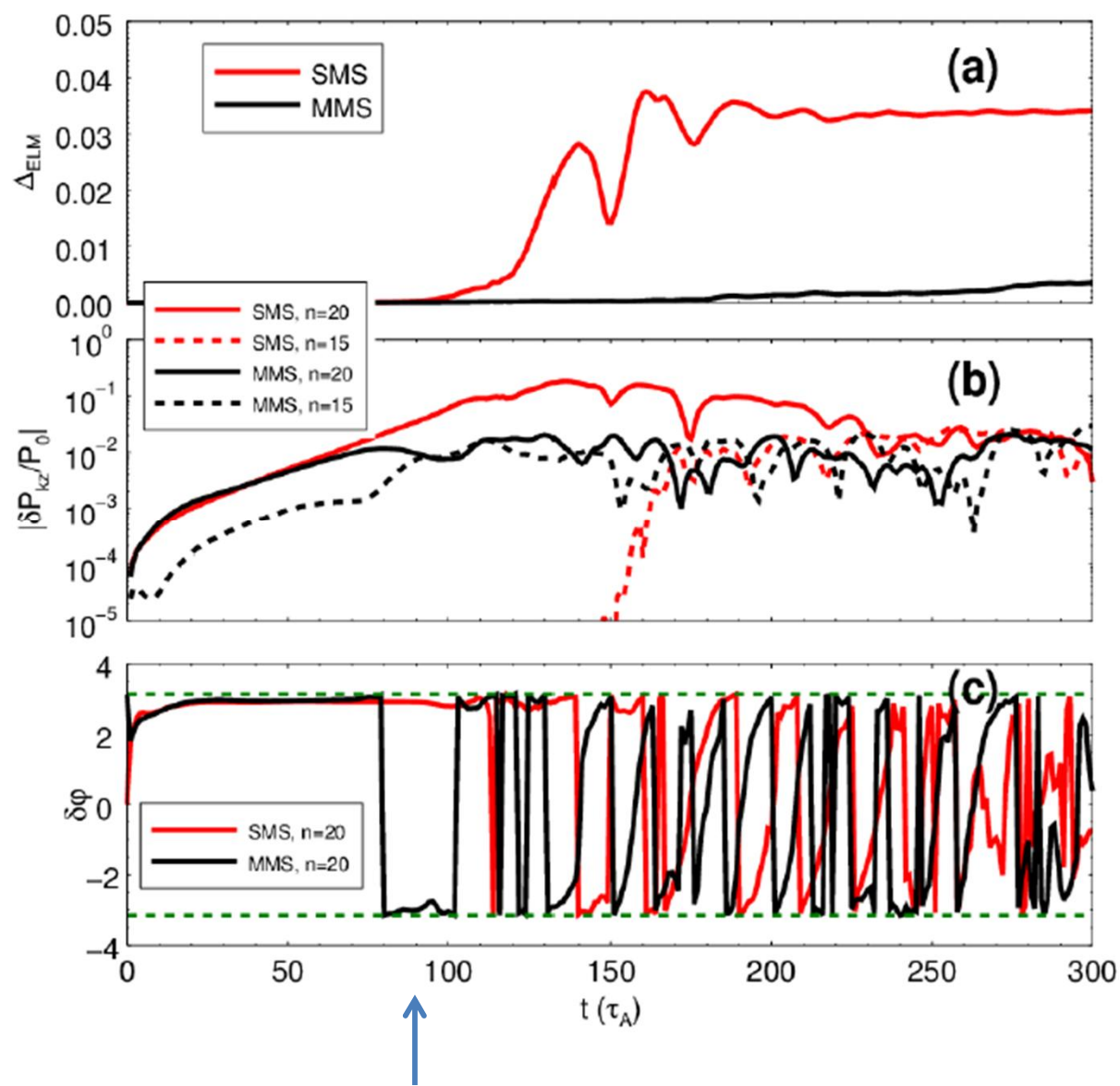
Critical element: relative phase

$$\delta\phi = \text{arg} [\hat{p}_n / \hat{\phi}_n]$$

Phase coherence time sets growth



Cross Phase Exhibits Rapid Variation in Multi-Mode Case



- Single mode case \rightarrow coherent phase set by linear growth \rightarrow rapid growth to 'burst'
- Multi-mode case \rightarrow phase de-correlated by mode-mode scattering \rightarrow slow growth to turbulent state

Key Quantity: Phase Correlation Time

- Ala' resonance broadening (Dupree '66):

$$\frac{\partial}{\partial t} \hat{P} + \tilde{v} \cdot \nabla \hat{P} + \langle v \rangle \cdot \nabla \hat{P} - D \nabla^2 \hat{P} = -\tilde{v}_r \frac{d}{dr} \langle P \rangle$$

Nonlinear
scattering

Linear streaming
(i.e. shear flow)

Ambient
diffusion

$$\hat{P} = A e^{i\phi}$$

Relative phase \leftrightarrow cross-phase

Amplitude

$$\hat{v} = B$$

Velocity amplitude

$$\rightarrow \partial_t \tilde{\phi} + \tilde{v} \cdot \nabla \tilde{\phi} + \langle v(r) \rangle \cdot \nabla \tilde{\phi} - D \nabla^2 \tilde{\phi} - \frac{2D}{A} \nabla A \cdot \nabla \tilde{\phi} = 0$$

NL scattering shearing

$$\partial_t A + \tilde{v} \cdot \nabla A + \langle v(r) \rangle \cdot \nabla A + D (\nabla \tilde{\phi})^2 A - D \nabla^2 A = -B \frac{d}{dr} \langle P \rangle$$

Damping by phase fluctuations

Phase Correlation Time

- Stochastic advection:

$$\frac{1}{\tau_{ck}} = \vec{k} \cdot D_\phi \cdot \vec{k} + k^2 D$$

$$D_\phi = \sum_{k'} \tau_{ck'} |\tilde{v}'_{\perp k}|^2$$

- Stochastic advection + sheared flow:

$$\frac{1}{\tau_{ck}} \approx \left(k_\perp^2 (D_\phi + D) \langle v_\perp \rangle'^2 \right)^{1/3} \quad \rightarrow \text{Coupling of radial scattering and Shearing shortens phase correlation}$$

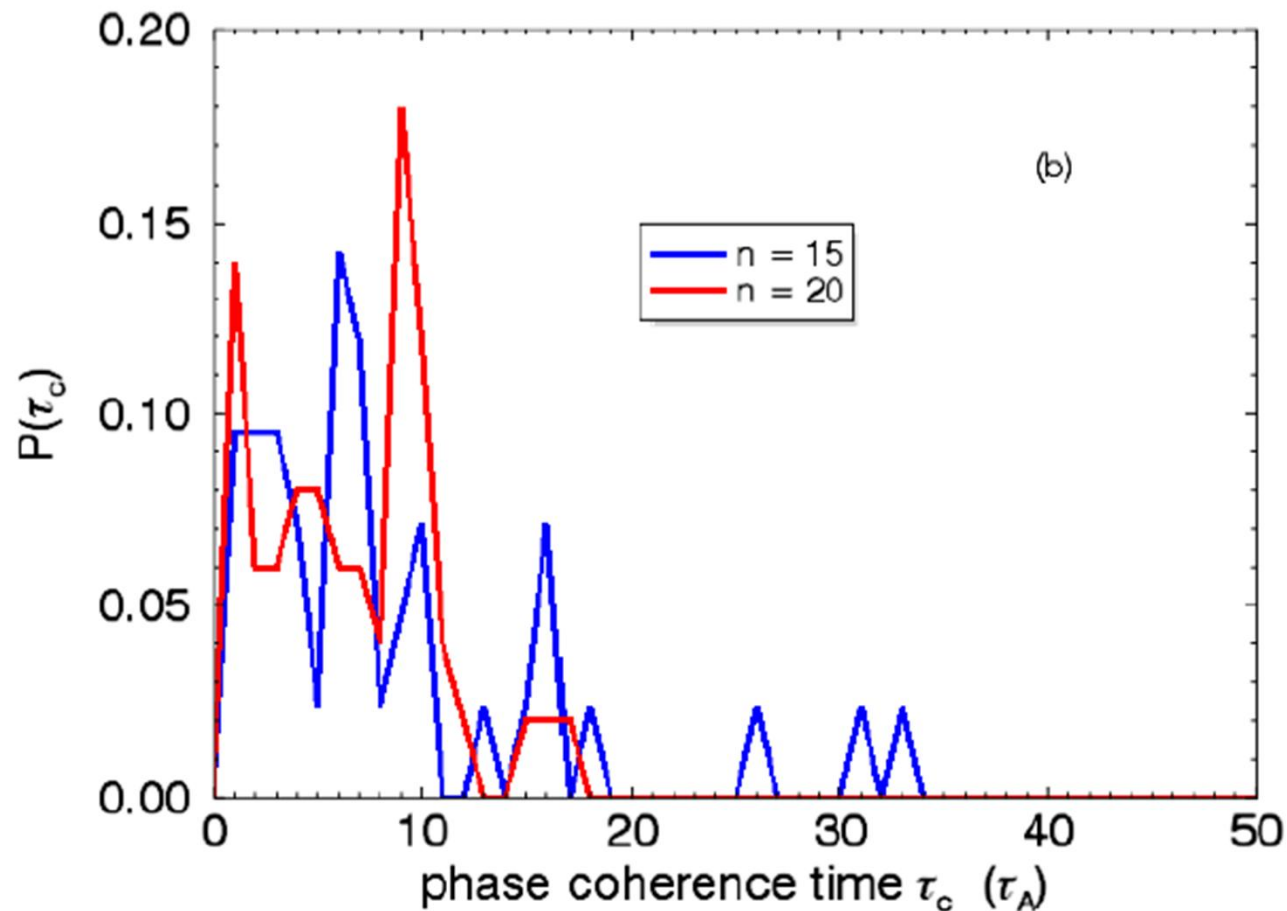
- Parallel conduction + diffusion:

$$\frac{1}{\tau_{ck}} \approx \left[\frac{\hat{s}^2 k_\perp^2}{(Rq)^2} \chi_\parallel (D_\phi + D) \right]^{1/2} \quad \rightarrow \text{Coupling of radial diffusion and conduction shortens phase correlation}$$

What is actually known about fluctuations in relative phase?

- For case of P.-B. turbulence, a broad PDF of phase correlation times is observed

pdf
of τ_c



Implications for: i) Bursts vs Turbulence
ii) Threshold

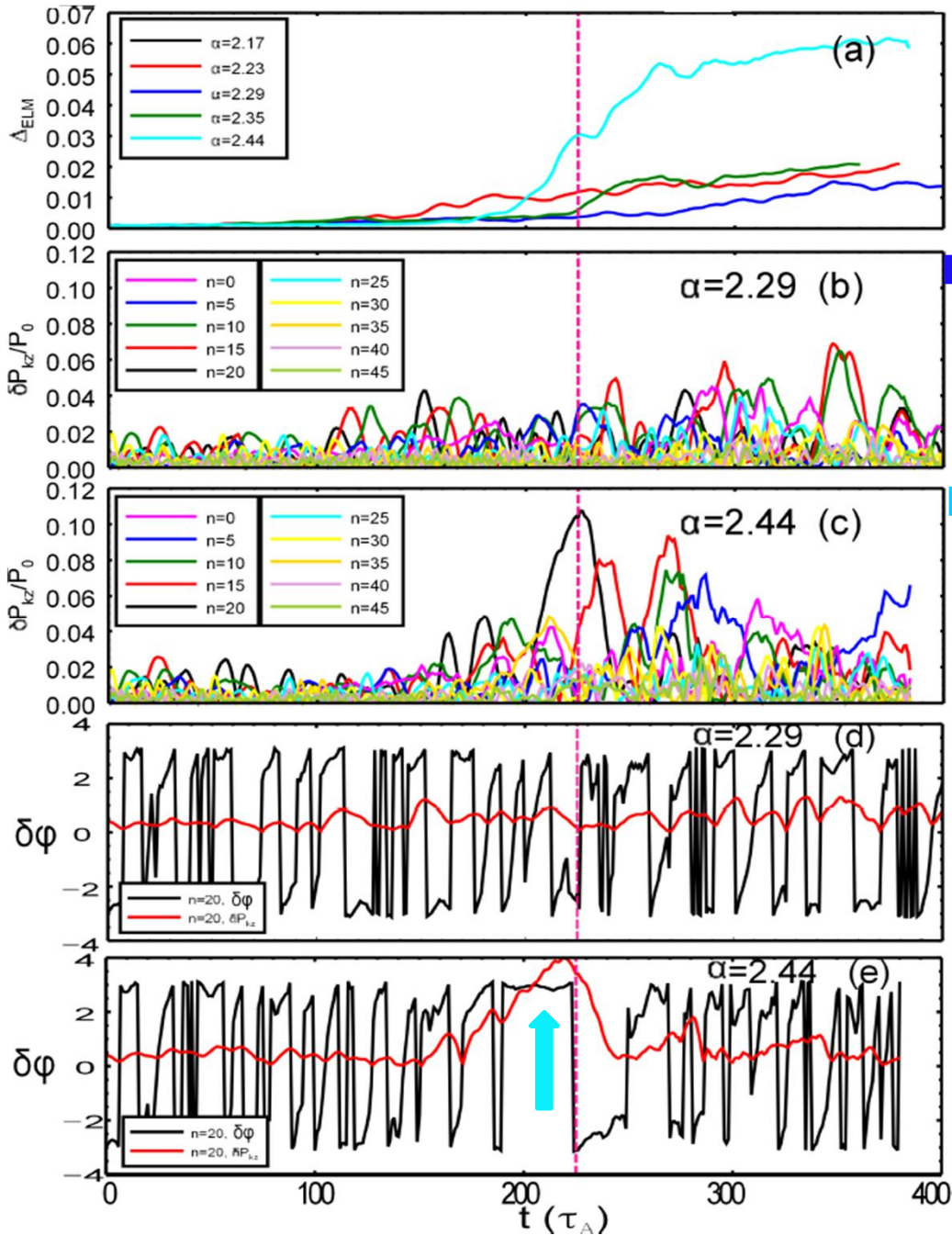
Bursts, Thresholds

- P.-B. turbulence can scatter relative phase and so reduce/limit growth of P.-B. mode to large amplitude
- Relevant comparison may be:

$$\gamma_k^L \text{ (linear growth) vs } \frac{1}{\tau_{ck}} \text{ (phase de-correlation rate)}$$

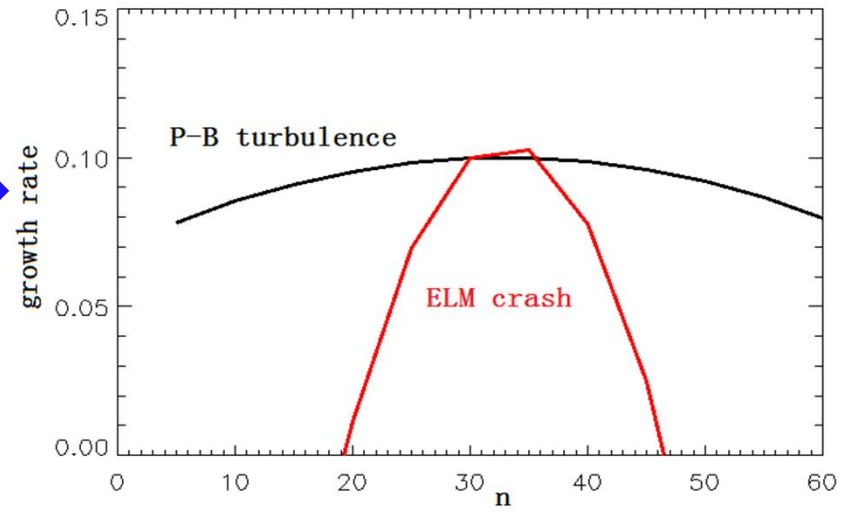
- Key point: Phase scattering for mode \vec{k} set by ‘background modes \vec{k}' ’ i.e. other P.-B.’s or micro-turbulence
→ is the background strong enough??

The shape of growth rate spectrum determines burst or turbulence



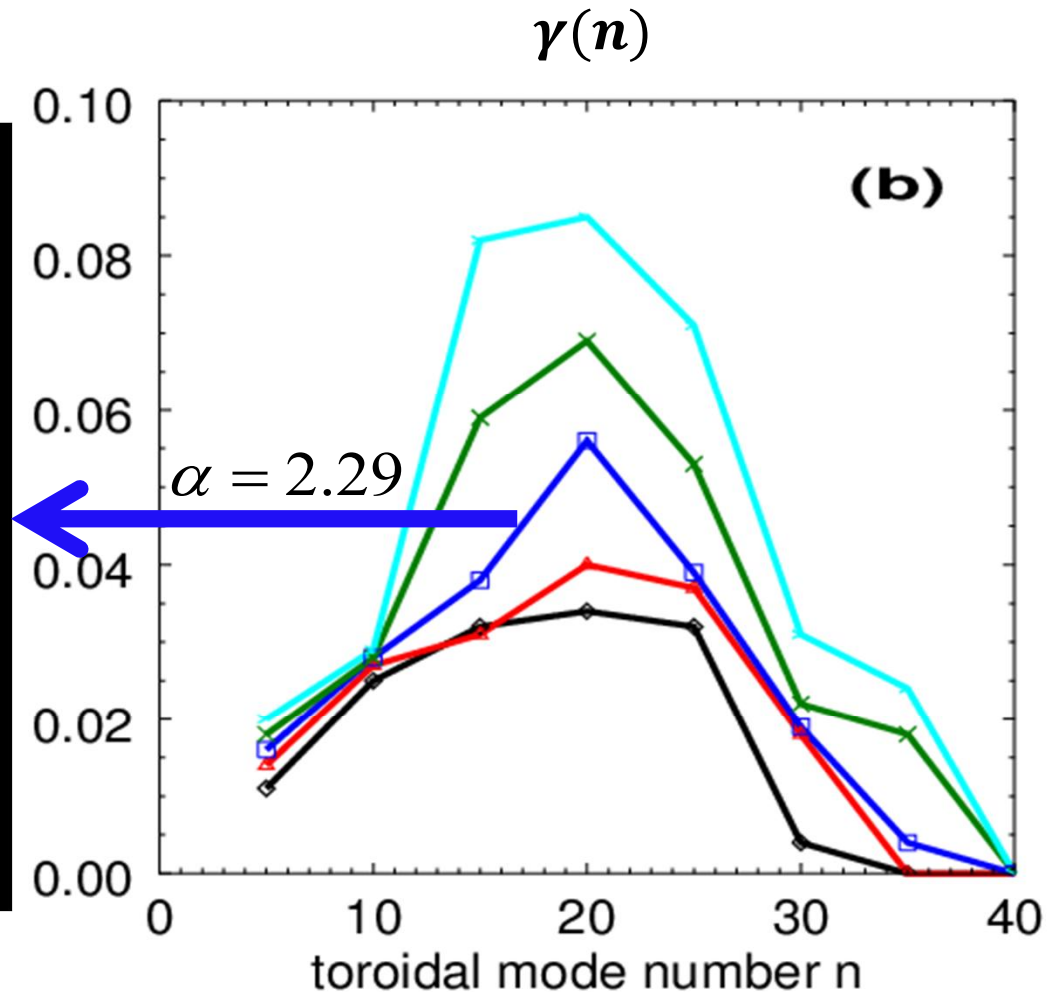
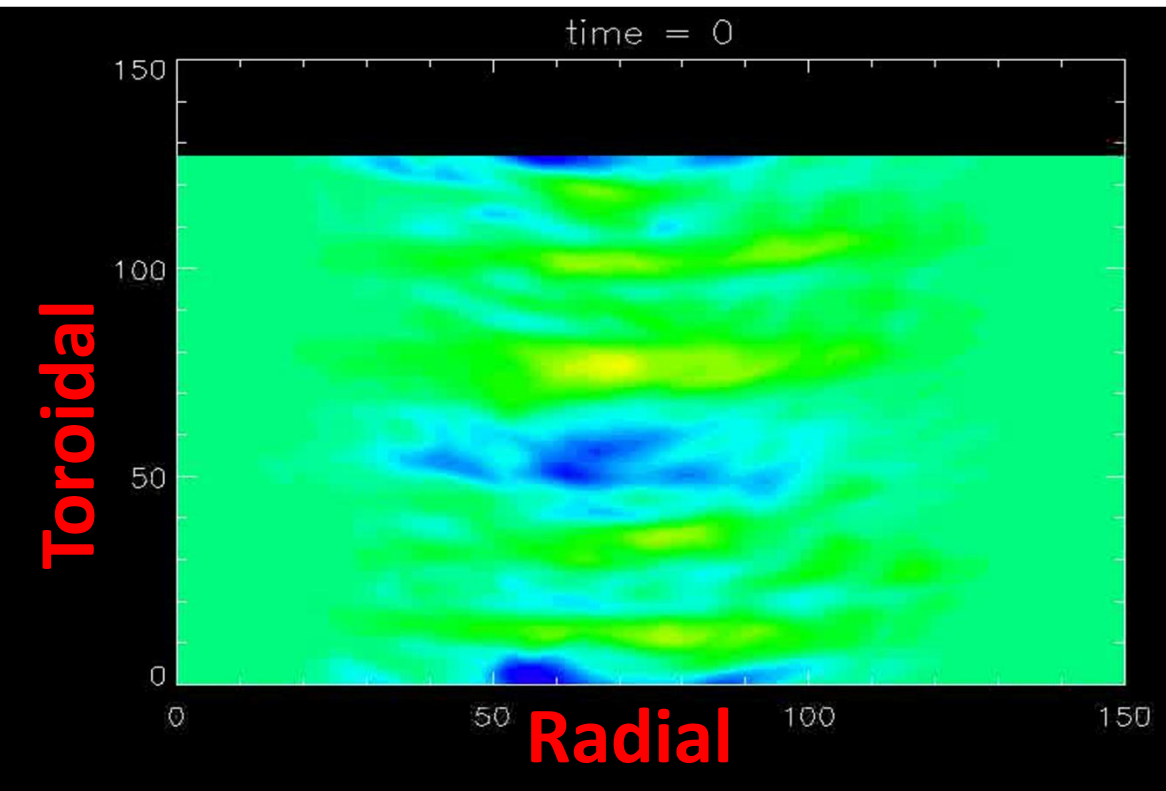
P-B turbulence
 $\gamma(n)\tau_c(n) < \ln 10$

Isolated ELM crash
 $\begin{cases} \gamma(n)\tau_c(n) > \ln 10, n = n_{dom} \\ \gamma(n)\tau_c(n) < \ln 10, n \neq n_{dom} \end{cases}$



So When Does it Crash?

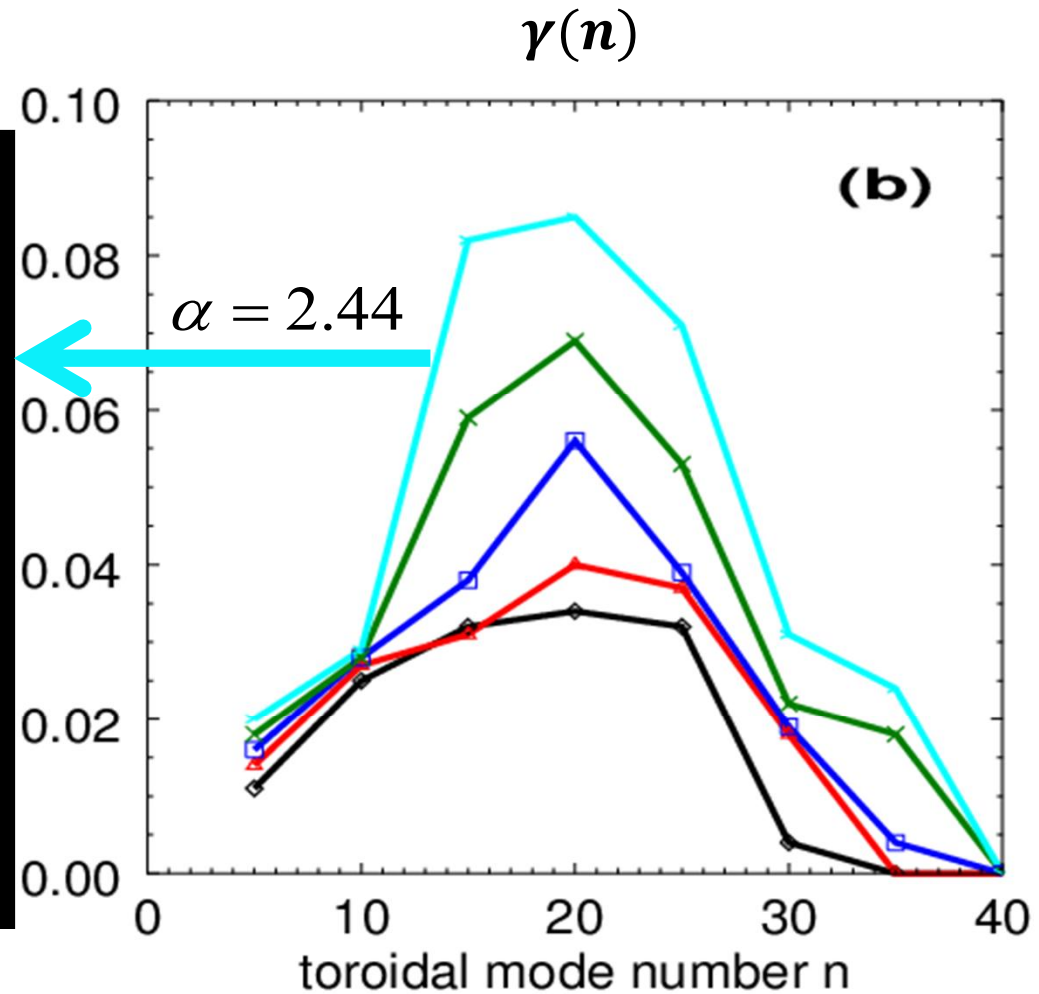
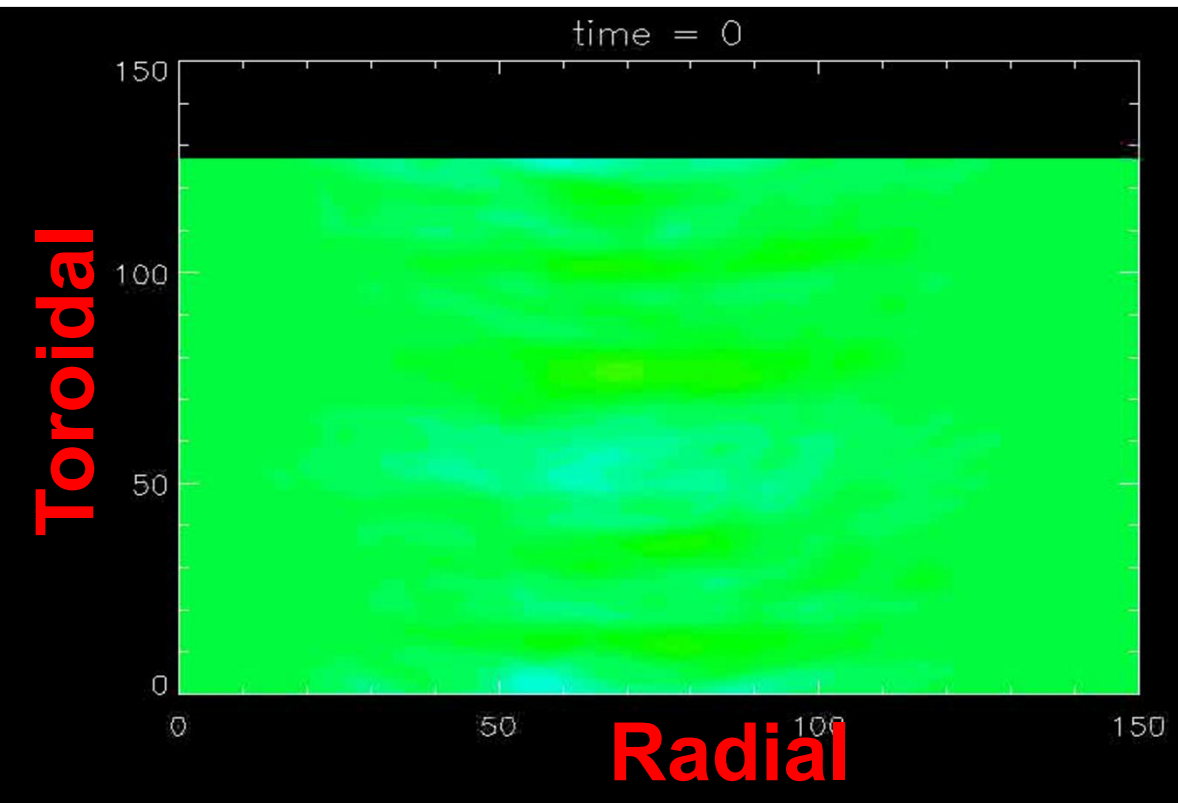
Modest $\gamma(n)$ Peaking \rightarrow P.-B. turbulence



- Evolution of P-B turbulence
 - No filaments
 - Weak radial extent

$$\alpha = -2\mu_0 R P_0' q^2 / B^2$$

Stronger Peaking $\gamma(n) \rightarrow$ ELM Crash

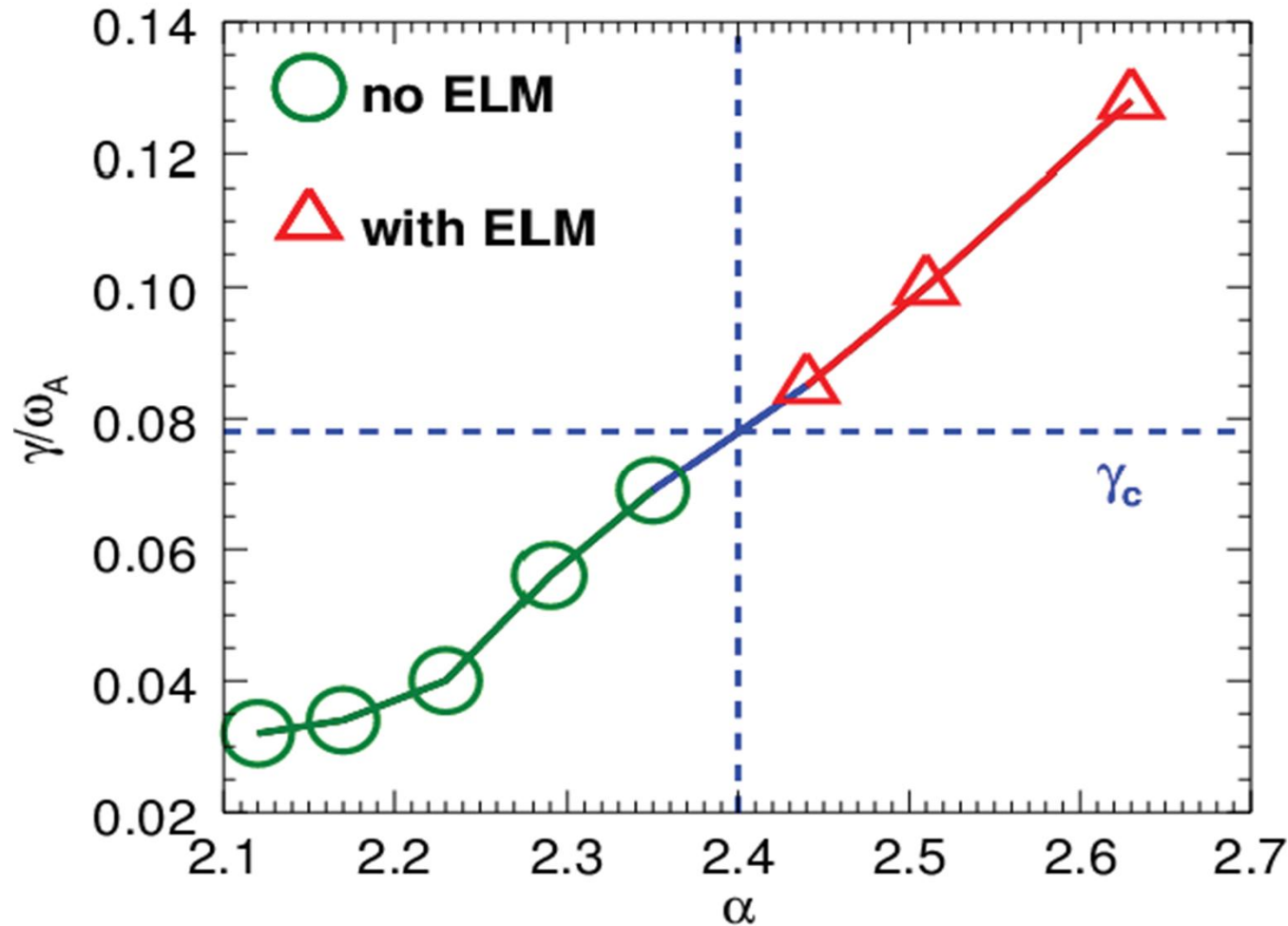


- ELM crash is triggered
- Wide radial extension

$$\alpha = -2\mu_0 R P_0' q^2 / B^2$$

Linear criterion for the onset of ELMs $\gamma > 0$ is replaced by the nonlinear criterion

$$\gamma > \gamma_c \sim 1 / \tau_c$$



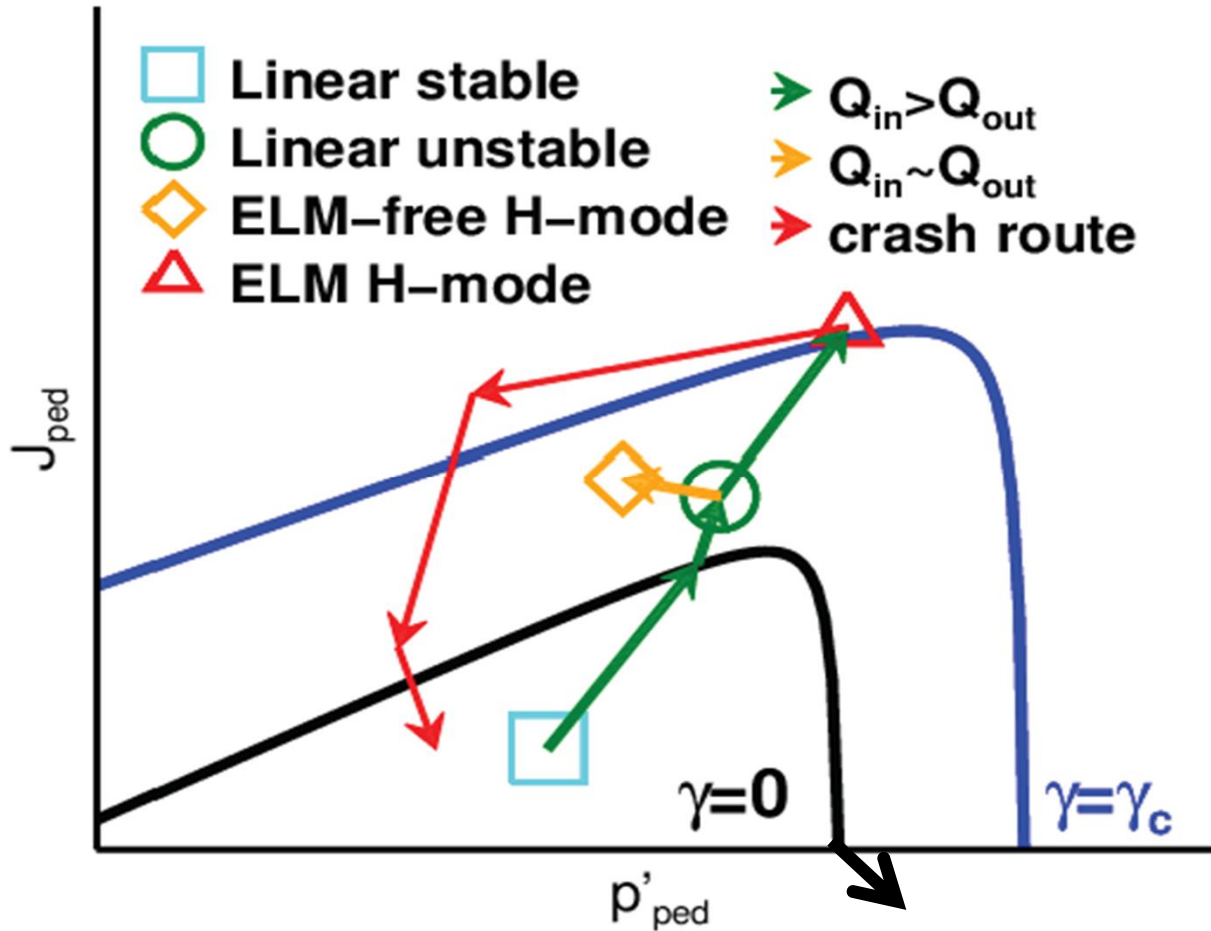
- Criterion for the onset of ELMs
 $\gamma\tau_c > \ln 10 \Rightarrow \gamma > \frac{\ln 10}{\tau_c} \equiv \gamma_c$

- Linear limit

$$\lim_{\tau_c \rightarrow \infty} \Rightarrow \gamma > 0$$

- γ_c is the critical growth rate which is determined by nonlinear interaction in the background turbulence
- N.B. $1 / \tau_c$ - and thus γ_{crit} - are functionals of $\gamma_L(n)$ peakedness

Nonlinear Peeling-ballooning model for ELM:



- $\gamma < 0$:
Linear stable region
- $0 < \gamma < \gamma_c$: Turbulent region
Possible ELM-free regime →
Special state: EHO, QCM (?!)
- $\gamma > \gamma_c$:
ELMy region

✓ Different regimes depend on **both linear instability and the turbulence** in the pedestal.

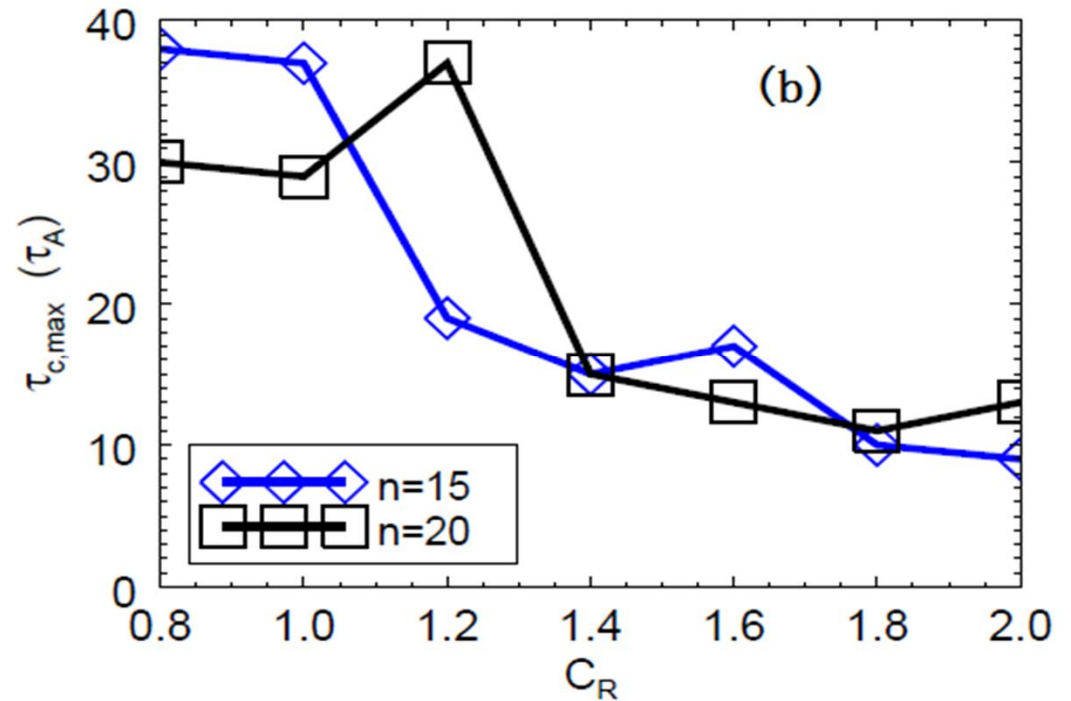
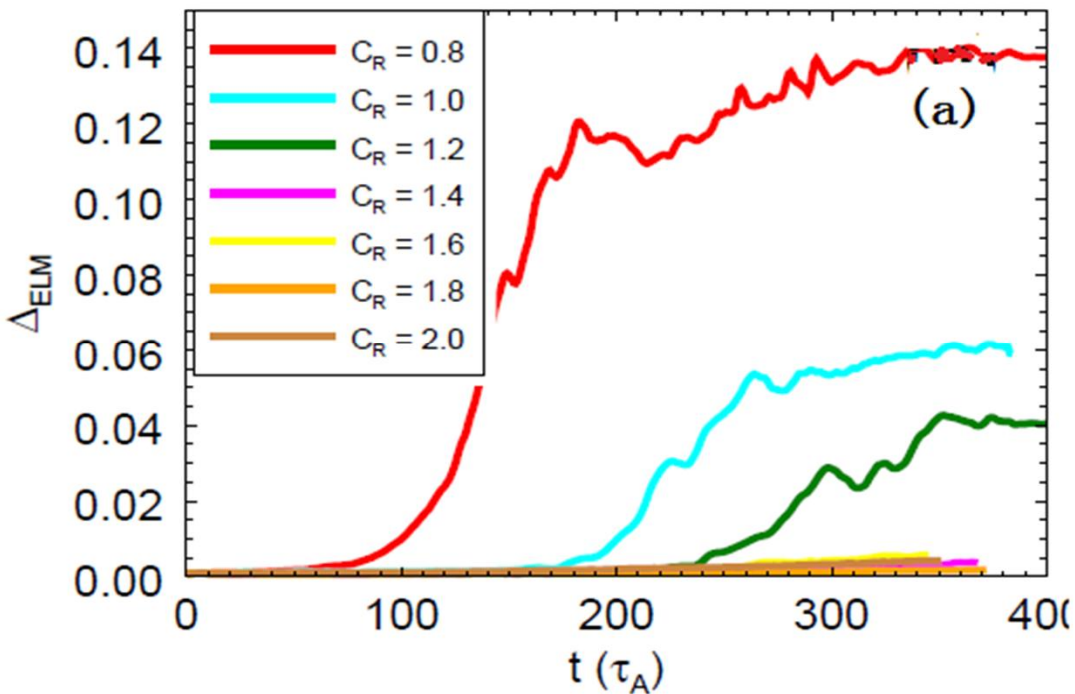
Including relevant linear physics
 (not only ideal P-B with ω_*)
 Resistivity / Electron inertia / ...

➔ **Turbulence can maintain ELM-free states**

How can these ideas be exploited
for ELM mitigation and control?

ELMs can be controlled by reducing phase coherence time

$$\frac{\partial \varpi}{\partial t} + C_R \frac{\mathbf{b} \times \nabla \phi}{B} \cdot \nabla \varpi = RHS \quad \text{i.e. scan } C_R \text{ for fixed profiles}$$



- ELMs are determined by the product $\gamma(n)\tau_c(n)$;
- Reducing the phase coherence time can limit the growth of instability;
- **Different turbulence states lead to different phase coherence times and, thus different ELM outcomes**

Keys to τ_c

- Scattering field
- ‘differential rotation’ in \hat{P} response to \hat{v}_r
→ enhanced phase de-correlation

Knobs:

- ExB shear
- Shaping
- Ambient diffusion
- Collisionality

Mitigation States:

- QH mode, EHO
- RMP
- SMBI
- ...

Scenarios

- QH-mode

- enhanced ExB shear $\rightarrow \frac{1}{\tau_c} \rightarrow (k_{\perp}^2 D \langle V_E \rangle'^2)^{1/3}$
- Triangularity strengthens shear via flux compression (Hahm, KHB)
- Enhanced de-correlation restricts growth time

Also:

- Is EHO peeling/kink + reduced τ_c ? How maintained?
- $\langle V_E \rangle'$ works via γ_L and τ_c

N.B. See Bin Gui, Xu; for more on shearing effects

Scenarios

- RMP

- $\frac{1}{\tau_c} = \left(\frac{k_{\perp}^2 \hat{s}^2}{(Rq)^2} \chi_{\parallel} D \right)^{1/2} \quad D = D_{\phi} + D_{amb}$

- RMP $\rightarrow D_{amb} \uparrow \rightarrow$ enhanced de-correlation

or

- Enhanced flow damping \rightarrow enhanced turbulence \rightarrow increased D_{ϕ}

(Leconte, P.D., Y. Xu)

- SMBI

- enhanced $D_{\phi} \rightarrow$ reduced τ_c ?

and/or

- Disruption of pedestal avalanches?

Phase Dynamic → A Deeper Look

Phase Dynamics in P.-B. Turbulence is INTERESTING

- i.e. usually $\tau_{ac} \sim \left| \left(\frac{\omega}{k} - \frac{d\omega}{dk} \right) \Delta k \right|^{-1}$ sets phase coherence time by \sim linear processes \rightarrow wave propagation and dispersion
- P.-B. turbulence in strong coupling regime $\omega \rightarrow 0$,
if insist on ω_* , non-dispersive
 $\therefore \tau_c$ set by nonlinear dynamics

Phase Dynamics in P.-B. Turbulence is HARD

- Recall:

$$\partial_t \tilde{\phi} + \tilde{v} \cdot \nabla \tilde{\phi} + \langle v(r) \rangle \cdot \nabla \tilde{\phi} - D \nabla^2 \tilde{\phi} - \frac{2D}{A} \nabla A \cdot \nabla \tilde{\phi} = 0$$

$$\partial_t A + \tilde{v} \cdot \nabla A + \langle v(r) \rangle \cdot \nabla A + D (\nabla \phi)^2 A - D \nabla^2 A = -B \frac{d\langle P \rangle}{dr}$$

- Turbulent \tilde{v} , self-consistency?
- ϕ, A coupling?
- Vorticity equation, Ohm's law?

Phase Dynamics in P-B Turbulence is HARD

- For hard problems recall advice of G. Polya in “How to Solve It”
 - ~ “If you didn’t know how to solve a problem, convert it to an easier problem you **do** understand.”
- What familiar paradigm(s) does the phase dynamics problem resemble?

Paradigms

- There are at least 2; both involving phase dynamics:

a) if ignore $2D\nabla A \cdot \nabla\phi / A$:

$$\partial_t \tilde{\phi} + \tilde{v} \cdot \nabla \tilde{\phi} + \langle v(r) \rangle \cdot \tilde{\phi} - D\nabla^2 \tilde{\phi} = \text{Noise}$$

→ scalar evolution with noise

→ if ignore feedback on \tilde{v} → **passive** scalar

∴ considerable body of insight into pdf[$\tilde{\phi}$].

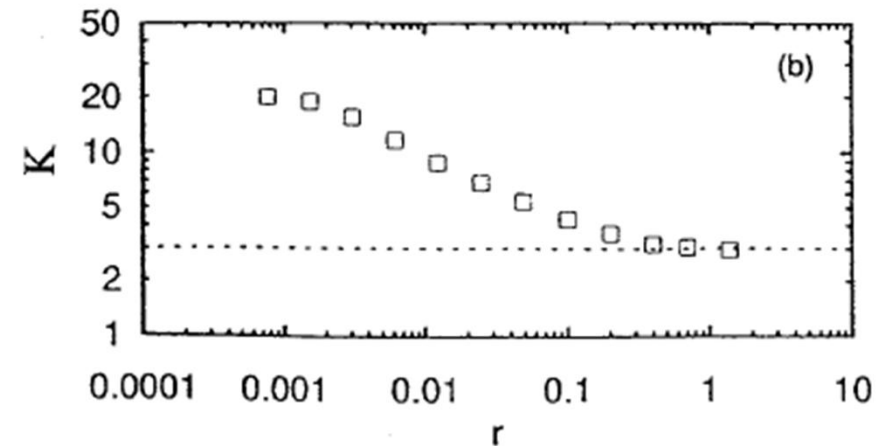
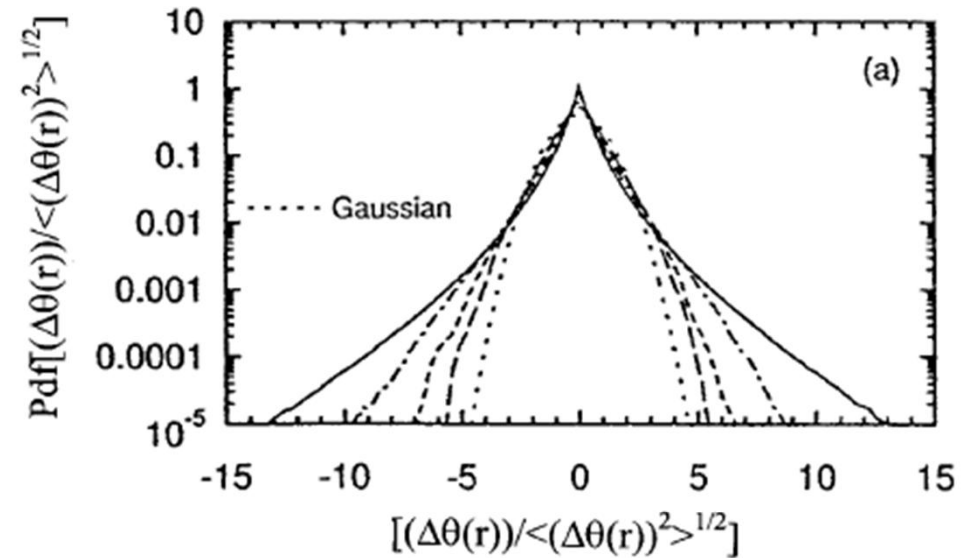
N.B. Little or no “physics” analysis of passive scalar statistics
with shear flow

A) Scalar

If $\langle v \rangle' \rightarrow 0$

$$\partial_t \tilde{\phi} + \tilde{v} \cdot \nabla \tilde{\phi} - D \nabla^2 \tilde{\phi} = 0$$

- PDF $\Delta\theta$ as function r
- Significant deviation from Gaussian for smaller scales
- Approaches Gaussian at large scales
- See also large kurtosis



Expect strongly non-Gaussian phase statistics

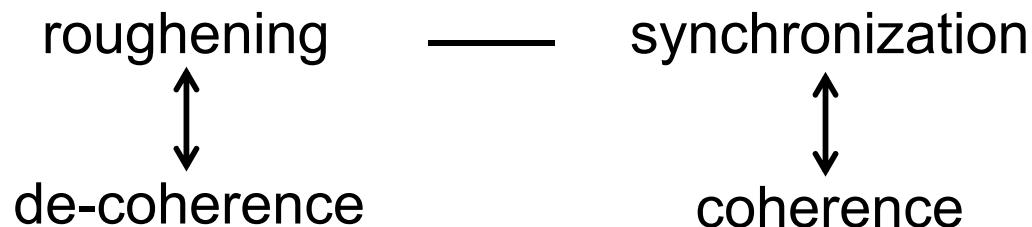
b) If ignore or simplify spatial structure, have synchronization problem \leftrightarrow Kuramoto

i.e.

i) Single oscillator

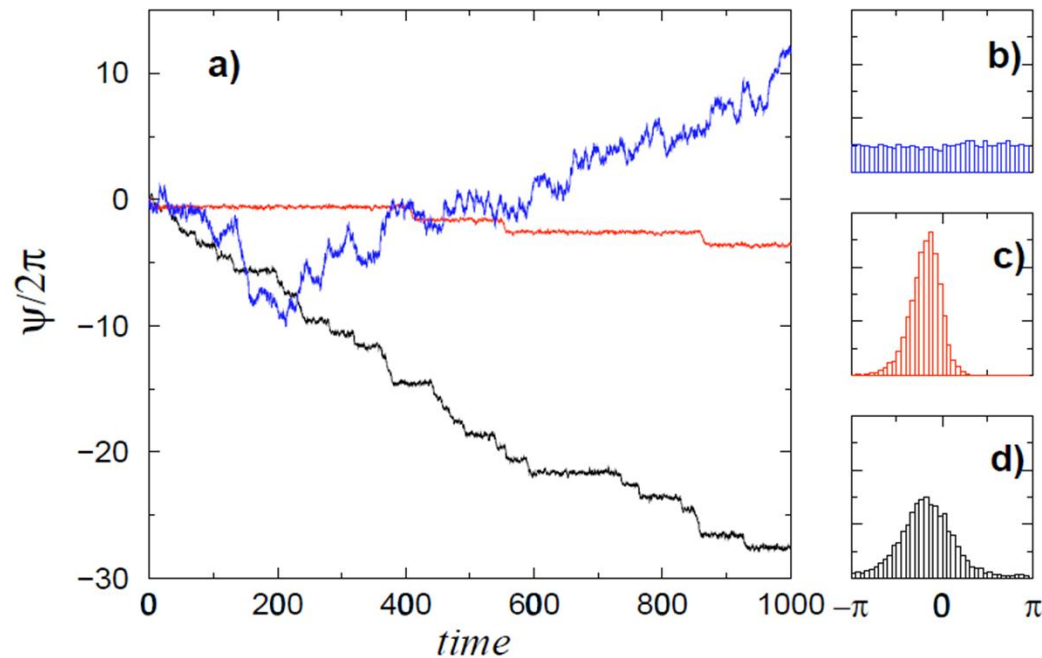
$$\frac{d\psi}{dt} = -\gamma + \epsilon q(\psi) + \epsilon(t)$$

ii) Oscillator lattice, continuum \rightarrow



$$\frac{\partial \phi(x,t)}{\partial t} = \omega(x) + \alpha \nabla^2 \phi(x,t) + \beta (\nabla \phi)^2 + \epsilon \quad (\text{extended KPZ})$$

B) Noisy Phase Dynamics \leftrightarrow Single Oscillator



- **Phase slips** occur
- Phase slips resemble cross phase jumps

- **Observe**

- No noise (blue) \rightarrow phase diffusive, distribution flat
- Weak noise (red) \rightarrow phase slips occur, distribution narrow
- Strong noise (black) \rightarrow more phase slips, broader distribution

Note:

→ Qualitative consistency of noisy oscillator phase and P.-B. turbulence cross phase

i.e. slips/jumps occurrence **increase** with
noise/turbulence level

→ in P.-B. turbulence,

- noise multiplicative
- profile evolution introduces another interaction channel.

→ Phase Dynamics Problem

Promising but Challenging

Conclusions – Coarse Grained

Conclusions

- ELM phenomena are intrinsically multi-mode and involve turbulence
- P.-B. growth regulated by phase correlation
 - determines crash + filament vs turbulence
- Phase coherence can be exploited for ELM mitigation

Where to Next?

- Simulations **MUST** move away from IVP – even if motivated by experiment – and to dynamic profile evolution, with:
 - sources, sinks i.e. flux drive and particle source essential
 - pedestal transport model
 - anomalous electron dissipationi.e. → - what profiles are actually achieved?
 - how evolve near P.-B. marginality?

- Should characterize:
 - pdf of phase fluctuations, correlation time
 - Dependence on τ_c control parameters
 - Threshold for burst
- Need understand feedback of P.-B. growth on turbulent hyper-resistivity
- Continue to develop and extend reduced models.

Some Inflammatory Questions

- Why bother with RMP?
 - Focus on ‘self-kinked state’ of QH with EHO.
Is this relevant and reproducible?
- Is SMBI ELM mitigation due:
 - phase de-correlation?
 - avalanche fragmentation?
 - stability of SMBI-mitigated EAST profiles?